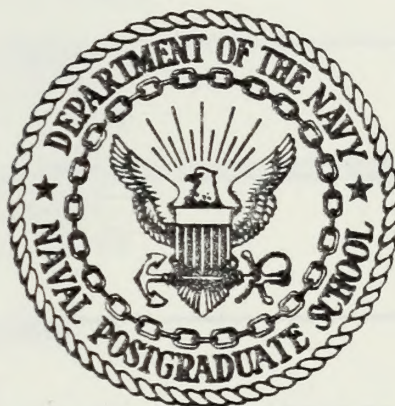


THEORETICAL ANALYSIS OF
TRANSONIC FLOW PAST
UNSTAGGERED OSCILLATING CASCADES

Peter Carlton Olsen

NAVAL POSTGRADUATE SCHOOL

Monterey, California



THESIS

THEORETICAL ANALYSIS OF
TRANSONIC FLOW PAST
UNSTAGGERED OSCILLATING CASCADES

by

Peter Carlton Olsen

September 1978

Thesis Advisor:

M.F. Platzer

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Unstaggered Oscillating Cascades

by

Peter Carlton Olsen
Lieutenant, United States Coast Guard
B.S., United States Coast Guard Academy, 1970
M.S., University of West Florida, 1975
M.S.A.E., Naval Postgraduate School, 1977
M.S.O.R., Naval Postgraduate School, 1978

Submitted in partial fulfillment of the
requirements for the degree of

AERONAUTICAL ENGINEER

from the

NAVAL POSTGRADUATE SCHOOL

September 1978

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LIST OF SYMBOLS

a	= local speed of sound	II
a _o	= speed of sound in the uniform flow	III
c	= blade semichord	III
f _j	= elementary function used in collocation solution	V
F	= specific energy ("head")	II
G	= function describing the surface of the airfoil as a function of time	III
H	= function specifying location of blade surface in the vertical axis	III
i	= $\sqrt{-1}$	II, III, IV, V
k	= Strouhal number, nondimensional frequency	III, IV, V
m	= $\sqrt{(\gamma+1)w}$	IV, V
M	= Mach number = $\frac{ \vec{V} }{a}$	III
n	= number of collocation points - 1, order of highest spanning function	V
p	= pressure	II
	= nondimensional interblade distance	V
R	= universal gas constant	II
R.P.	= "real part of"	III, IV, V
T	= temperature	II
t	= time, nondimensional time	II, III, IV, V
U _o	= uniform velocity from infinity	II, III, VI
u	= x-component of velocity	II, III
u ^o , u ¹	= interference vertical velocities due to reference and adjacent blades respectively, solved so as to satisfy the tangential flow conditions	V, VI

u'	= small disturbance velocity	III
v	= y-component of velocity	II,III,IV,V
\vec{V}	= general velocity vector	II
v'	= small disturbance velocity	III
v^0, v^1	= vertical velocities due to the reference and adjacent blades respectively, determined from the tangential flow condition	V
w	= $\tilde{\phi}_x$, a constant used in Gorelov's approximation of the transonic flow potential	IV,V
x	= horizontal coordinate, may be non-dimensional	
x_*	= mp (transformed interblade distance in Gorelov's approximation)	V
x_l	= blade leading edge	IV
x_0	= center of pitch of the unstaggered cascade	IV,V
y	= vertical coordinate, may be non-dimensional	
y, y_1	= vertical coordinates attached to the reference and adjacent blades respectively, may be non-dimensional	IV,V
z, z_1	= transformed vertical coordinates used in Gorelov's approximation, attached to the reference and adjacent blades respectively. $z = my, z_1 = my_1$	IV,V
$\frac{D}{Dt}$	= substantial derivative w.r.t. time $= \frac{\partial}{\partial t} + \frac{\partial x}{\partial t} \frac{\partial}{\partial x} + \frac{\partial y}{\partial t} \frac{\partial}{\partial y}$ $= \frac{\partial}{\partial t} + u \frac{\partial}{\partial x} + v \frac{\partial}{\partial y}$	II,III
$O(\omega^2)$	= "of the order of magnitude of ω^2 "	V
α	= angle of attack	II,III,IV,V
α_0	= maximum amplitude of pitch oscillations	IV

γ	= ratio of specific heats, c_p/c_v	
δ_{io}	= Dirac Delta function = 1 when $i = 0$ = 0 when $i \neq 0$	V
η	= $\cos^{-1}(-x)$	V
η_*	= $\cos^{-1}(1-x_*)$	V
$\bar{\eta}$	= $\cos^{-1}(-s)$	V
$\hat{\eta}$	= $\cos^{-1}(x_*-x)$	V
θ_j^0, θ_j^1	= interference potential coefficients for reference and adjacent blades respectively	V
λ	= k/m^2	IV,V
μ_j^0, μ_j^1	= Fourier coefficients describing the motion of the reference and adjacent blades respectively	V
ν	= angular frequency of oscillation	IV,V
ρ	= density	II
σ	= phase angle	V
τ	= cascade solidity, $\frac{2}{p}$	V
Φ	= general velocity potential	II,III,VI
Φ_0	= uniform flow velocity potential	
$\tilde{\Phi}$	= steady flow perturbation potential	III,IV
ϕ^0, ϕ^1	= perturbation potential in collocation solution due to reference and adjacent blades respectively	V,VI
Φ^0, Φ^1	= transformed potentials	V
ψ	= oscillatory flow potential	III,IV,V
ϕ	= transformed oscillatory potential in Gorelov's coordinates corresponding to ψ	V
ψ^0, ψ^1	= interference potentials due to reference and adjacent blades respectively	V,VI

ψ^0, ψ^1 = transformed potentials in Gorelov's coordinates, corresponding to ψ^0 and ψ^1 V

$$\omega = \frac{k(1-m)^2}{m^4}$$

∇ = $\vec{i} \frac{\partial}{\partial x} + \vec{j} \frac{\partial}{\partial y}$, Gradient operator, \vec{i} and \vec{j} are unit vectors in the x and y direction respectively

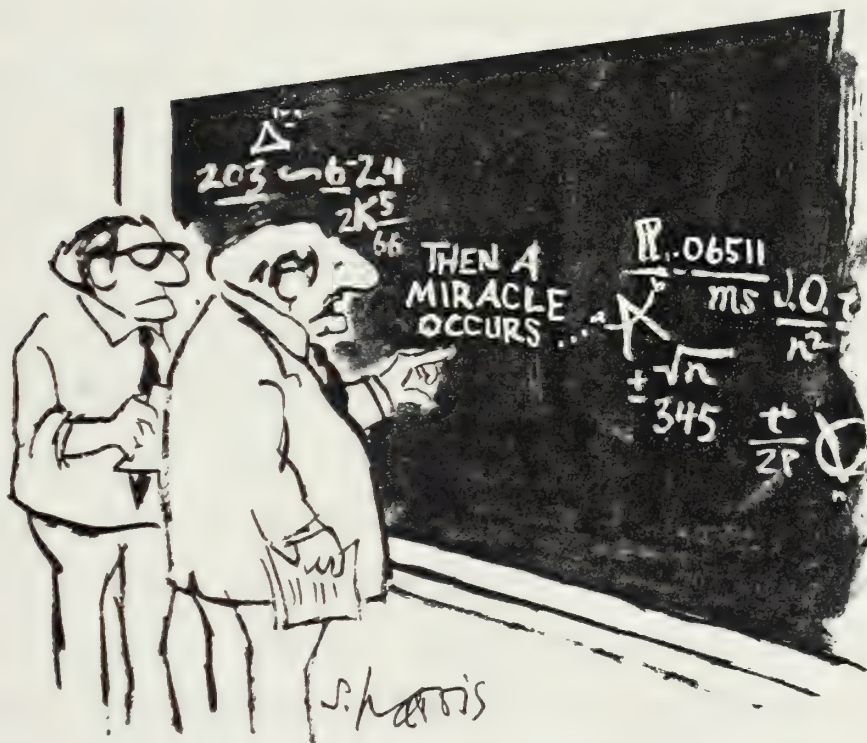
Computer Variables

DK	=	Reduced Frequency, k
DLAMDA	=	λ
DM2	=	m^2
DR	=	$mp = x_*$
ETA	=	η
ETASTR	=	η_*
IPT	=	Print Parameter
N	=	n
NF	=	not used
OFFSET	=	r
OMEGA	=	ω
QALPHA	=	$0 + i(\lambda - k)$
QCONST	=	$e^{i\sigma}$
QDCL	=	C_{l_α}
QDCM	=	C_{m_α}
QDK	=	$0 + ik$
QEXP	=	$0 + i\lambda$
QINTAP, QINTRP	=	Variables used to transmit boundary condition integrals
Q1ABCF, Q1RBCF	=	Interference coefficients for adjacent and reference blades
Q1COF	=	Right hand side vector in collocation solution
Q1INT	=	Known matrix of integrals in collocation solution
Q1RBP, Q1ABP	=	Not used
Q2CP	=	Not used
Q2EXP	=	$e^{-i\lambda x}$

Q2PT	=	Not used
RHO	=	vertical displacement
RHP	=	local input variable for rho
SIGMA	=	σ
TAU	=	τ
XASTN	=	Current x station in adjacent blade coordinates
XSTN	=	Current x station in reference blade coordinates

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"I think you should be more explicit here in step two."

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I. INTRODUCTION

The analysis of unsteady transonic flows in aircraft turbopropulsion is an area of intense current interest. Rising fuel prices and increasing thrust requirements both point toward the need of turbomachinery capable of performing well with transonic or supersonic internal flow. But, increased flow has increased both the costs and uncertainties of engine designs. Flutter problems have already become a major consideration in engine development. Problems unforeseen in earlier days of turbine engine production have caused long development delays, or forced acceptance of engines producing less than their initial design thrust. These uncertainties cannot be avoided when an attempt is made to extend the state of the art, but they can be reduced by extending the range of analytical modeling.

Such extension must now be done piecemeal. The three-dimensional flows in turbomachinery, including the simultaneous effects of boundary layers, rotation, finite blade thickness, spanwise Mach distributions, and shocks, are well beyond present capability. Perhaps one day complete analysis will be practical, but it is not today. The best that can be done now is to approach the problem from one aspect at a time. Flow through a two dimensional cascade has been a useful tool in this partial analysis.

This thesis was originally to have been an extension of the work of Elder [1] and Schlein [2] to the case of a staggered cascade. Their work, based on Teipel's [3] linearization of the unsteady transonic small perturbation equation, analyzed transonic flow through oscillating unstaggered cascades by use of the collocation method. While the problem was easy to state, it was difficult to solve. Both Elder and Schlein had encountered difficulty in employing the collocation method. Therefore, it was decided that verification of the basic collocation solution presented by Gorelov [4] using a different linearization would be a worthwhile goal in itself.

The following investigation presents a verification of the development in [4], along with numerical results and suggestions for further work.

II. UNSTEADY TRANSONIC FLOW THEORY

Considering inviscid flow only, the following four equations govern the aerodynamic flow problem at hand:

The equation of state

$$p = \rho RT \quad (\text{II-1})$$

and the equations for the conservation of

$$1. \text{ Mass: } \operatorname{div}(\rho \vec{v}) + \frac{\partial \rho}{\partial t} = 0 \quad (\text{II-2})$$

$$2. \text{ Momentum: } \frac{D\vec{v}}{Dt} + \frac{1}{\rho} \nabla p = 0 \quad (\text{II-3})$$

$$3. \text{ Energy: } \frac{DS}{Dt} = 0 \quad (\text{II-4})$$

where

\vec{v} = velocity

p = pressure

S = entropy

R = universal gas constant

T = temperature

t = time

ρ = density

$$\frac{D}{Dt} = \frac{\partial}{\partial t} + \frac{\partial}{\partial x} \frac{\partial x}{\partial t} + \frac{\partial}{\partial y} \frac{\partial y}{\partial t}$$

The analysis starts with a uniform flow from infinity.

This flow has velocity U_0 parallel to the x-axis. This

formulation can be simplified by working with the total velocity potential, Φ , where

$$u = \frac{\partial \Phi}{\partial x} = \Phi_x = x \text{ component of velocity} = \frac{\partial x}{\partial t} \quad (\text{II-5})$$

$$v = \frac{\partial \Phi}{\partial y} = \Phi_y = y \text{ component of velocity} = \frac{\partial y}{\partial t} \quad (\text{II-6})$$

Thus, the initial uniform flow is represented by the uniform flow potential

$$\Phi_0 = U_0 x \quad (\text{II-7})$$

This notation may be applied to the conservation equations for mass and momentum. The equation for conservation of mass

$$\text{div}(\rho \vec{v}) + \frac{\partial \rho}{\partial t} = 0 \quad (\text{II-8})$$

becomes for two-dimensional unsteady flow

$$\begin{aligned} \frac{\partial(\rho u)}{\partial x} + \frac{\partial(\rho v)}{\partial y} + \frac{\partial \rho}{\partial t} &= 0 \\ \left[\frac{\partial \rho}{\partial t} + u \frac{\partial \rho}{\partial x} + v \frac{\partial \rho}{\partial y} \right] + \rho \left(\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} \right) &= 0 \end{aligned} \quad (\text{II-9})$$

but

$$\frac{\partial \rho}{\partial t} + u \frac{\partial \rho}{\partial x} + v \frac{\partial \rho}{\partial y} = \frac{D\rho}{Dt}$$

and

$$u = \phi_x \quad \text{and} \quad v = \phi_y$$

Thus

$$\frac{D\rho}{Dt} + \rho(\phi_{xx} + \phi_{yy}) = 0$$

and

$$\phi_{xx} + \phi_{yy} = -\frac{1}{\rho} \frac{D\rho}{Dt} \quad (\text{II-10})$$

The speed of sound is given by

$$a^2 = \frac{dp}{d\rho}$$

Thus

$$\frac{D\rho}{Dt} = \frac{d\rho}{dp} \cdot \frac{Dp}{Dt} = \frac{1}{a^2} \frac{Dp}{Dt}$$

Applying this to equation (II-10) yields

$$\phi_{xx} + \phi_{yy} = -\frac{1}{\rho a^2} \frac{Dp}{Dt} \quad (\text{II-11a})$$

$$\phi_{xx} + \phi_{yy} = -\frac{1}{\rho a^2} (u p_x + v p_y + p_t) \quad (\text{II-11b})$$

$$= -\frac{1}{\rho a^2} [(\nabla\phi) \cdot (\nabla p) + p_t] \quad (\text{II-11c})$$

where ∇ is the gradient operator, $P_x = \frac{\partial P}{\partial x}$, $P_y = \frac{\partial P}{\partial y}$, $P_t = \frac{\partial P}{\partial t}$.

Laying this aside for the moment, consider the momentum equation (II-3)

$$\frac{D\vec{v}}{Dt} + \frac{1}{\rho} \nabla p = 0$$

$$\frac{D\vec{v}}{Dt} = \frac{\partial \vec{v}}{\partial t} + (\vec{v} \cdot \nabla) \vec{v}$$

Thus

$$\rho \left[\frac{\partial \vec{v}}{\partial t} + (\vec{v} \cdot \nabla) \vec{v} \right] = \nabla p \quad (\text{II-12})$$

$$\vec{v} = \nabla \phi$$

so

$$\frac{\partial \vec{v}}{\partial t} = \frac{\partial}{\partial t} (\nabla \phi) = \nabla \frac{\partial \phi}{\partial t} \quad (\text{II-13})$$

and

$$(\vec{v} \cdot \nabla) \vec{v} = \nabla \frac{v^2}{2} - \vec{v} \times (\nabla \times \vec{v})$$

where

$$v^2 = u^2 + v^2$$

$$= \vec{v} \cdot \vec{v}$$

For irrotational flow

$$\nabla \times \vec{v} = 0$$

thus

$$(\vec{v} \cdot \nabla) \vec{v} = \frac{\nabla v^2}{2} \quad (\text{II-14})$$

Thus

$$\frac{\nabla p}{\rho} + \nabla \left[\phi_t + \frac{v^2}{2} \right] = 0 \quad (\text{II-15})$$

which after integration along a streamline becomes

$$\int \frac{dp}{\rho} + \phi_t + \frac{v^2}{2} = F(t) \quad (\text{II-16})$$

For uniform flow from infinity $F(t) = \frac{1}{2} U_o^2$ and thus the final result is

$$\int \frac{dp}{\rho} + \phi_t + \frac{v^2}{2} = \frac{1}{2} U_o^2 \quad (\text{II-17})$$

Differentiation with respect to t gives

$$p_t = -\rho(\phi_{tt} + \frac{1}{2} \frac{\partial V^2}{\partial t}) \quad (\text{II-18})$$

From (II-3)

$$-\nabla p = \rho \frac{D\vec{v}}{Dt} \quad (\text{II-19})$$

Substitute (II-18) and (II-19) into (II-11c) to obtain

$$\phi_{xx} + \phi_{yy} = \frac{1}{a^2} [(\nabla\phi) \cdot \frac{D\vec{v}}{Dt} + \phi_{tt} + \frac{1}{2} \frac{\partial V^2}{\partial t}] \quad (\text{II-20})$$

This may be further simplified

$$\begin{aligned} \frac{D\vec{v}}{Dt} &= \frac{\partial \vec{v}}{\partial t} + (\vec{v} \cdot \nabla) \vec{v} \\ &= \frac{\partial \vec{v}}{\partial t} + \frac{1}{2} \nabla V^2 \end{aligned}$$

Hence

$$\phi_{xx} + \phi_{yy} = \frac{1}{a^2} [(\nabla\phi) \cdot (\frac{\partial \vec{v}}{\partial t} + \frac{1}{2} \nabla V^2) + \phi_{tt} + \frac{1}{2} \frac{\partial V^2}{\partial t}] \quad (\text{II-21})$$

Expanding terms

$$\begin{aligned} (\nabla\phi) \cdot (\frac{\partial \vec{v}}{\partial t}) &= \nabla\phi \cdot (\frac{\partial}{\partial t} \nabla\phi) = \phi_x \phi_{xt} + \phi_y \phi_{yt} \\ \nabla\phi \cdot \frac{\nabla V^2}{2} &= \frac{\phi_x^2 \phi_{xx}}{2} + \frac{\phi_y^2 \phi_{yy}}{2} + \frac{\phi_x \phi_y \phi_{xy}}{2} + \frac{\phi_y \phi_x \phi_{yx}}{2} \\ &= \frac{\phi_x^2 \phi_{xx}}{2} + \frac{\phi_y^2 \phi_{yy}}{2} + \frac{2\phi_x \phi_y \phi_{xy}}{2} \end{aligned}$$

$$\begin{aligned}
\frac{1}{2} \frac{\partial V^2}{\partial t} &= \frac{1}{2} \frac{\partial}{\partial t} [(\nabla \phi) \cdot (\nabla \phi)] \\
&= \frac{1}{2} [\phi_x \phi_{xt} + \phi_{xt} \phi_x + \phi_y \phi_{yt} + \phi_{yt} \phi_y] \\
&= \phi_x \phi_{xt} + \phi_y \phi_{yt}
\end{aligned}$$

The final result obtained by combining terms is

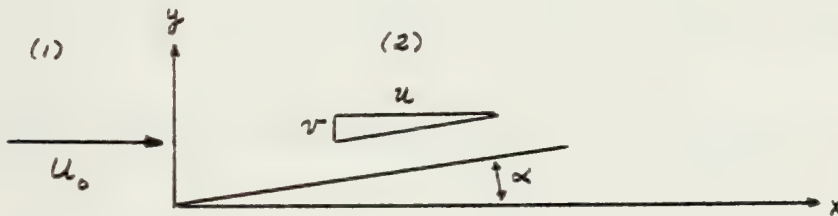
$$\begin{aligned}
(1 - \frac{\phi_x^2}{a^2})_{xx} + (1 - \frac{\phi_y^2}{a^2})_{yy} - \frac{2\phi_x \phi_y \phi_{xy}}{a^2} \\
- \frac{2\phi_x}{a^2} \phi_{xt} - \frac{2\phi_y}{a^2} \phi_{yt} - \frac{1}{a^2} \phi_{tt} = 0
\end{aligned} \tag{II-22}$$

This result is valid for irrotational, inviscid, two-dimensional, unsteady, compressible flows where gravity has been neglected.

III. SMALL PERTURBATION THEORY OF TRANSONIC FLOW

A. GENERAL CASE

A thin body at a small angle of attack will cause only a slight disturbance in the fluid. A flat plate is an example. Consider flow past a flat plate at angle of attack, α .



The flow at (2) must be parallel to the plate. To achieve this, small disturbance velocities u' and v' must be added to the free stream velocity yielding

$$u = U_0 + u'$$

$$v = v'$$

The potential of the disturbed flow may be considered as the sum of the uniform flow potential, $\Phi_0 = U_0 x$, and the disturbance potential, ϕ

$$\Phi = \Phi_0 + \phi \quad (\text{III-1})$$

Thus

$$\phi_x = U_o + u' \quad (\text{III-2a})$$

$$\phi_y = v \quad (\text{III-2b})$$

If ϕ is a function of time, then

$$\phi_t = \phi_t$$

This result may be substituted into (II-22) leading to

$$\begin{aligned} & \left[1 - \frac{(U_o + u')^2}{a^2}\right] \phi_{xx} + \left[1 - \frac{v^2}{a^2}\right] \phi_{yy} - 2 \frac{(U_o + u')v}{a^2} \phi_{xy} \\ & - 2 \frac{(U_o + u')}{a^2} \phi_{xt} - \frac{2v}{a^2} \phi_{yt} - \frac{1}{a^2} \phi_{tt} = 0 \end{aligned} \quad (\text{III-3})$$

This expands to yield

$$\begin{aligned} & \left[1 - \frac{U_o^2 + 2U_o u' + u'^2}{a^2}\right] \phi_{xx} + \left[1 - \frac{v^2}{a^2}\right] \phi_{yy} - 2 \frac{(U_o v + u'v)}{a^2} \phi_{xy} \\ & - 2 \frac{(U_o + u')}{a^2} \phi_{xt} - \frac{2v}{a^2} \phi_{yt} - \frac{1}{a^2} \phi_{tt} = 0 \end{aligned} \quad (\text{III-4})$$

Equation (III-4) may be further simplified as shown by Landahl [3]. All non-linear terms except the $\phi_x \phi_{xx}$ product

term can be neglected yielding the following transonic small disturbance equation

$$\begin{aligned} [(M^2-1) + (\gamma+1)M^2 \frac{\phi_x}{U_0}] \phi_{xx} \\ - \phi_{yy} + \frac{2M_0}{a_0} \phi_{xt} + \frac{1}{a_0} \phi_{tt} = 0 \end{aligned} \quad (\text{III-5})$$

where a_0 is the velocity of sound in the free-stream, and γ is the ratio of specific heats, and M_0 = Mach number.

B. BOUNDARY CONDITION

The tangential flow condition requires that the flow be tangent to the airfoil surface at each instant of time. This means that no fluid may flow through the surface of the airfoil and is expressed by the condition

$$\frac{DG}{Dt} = 0 \quad \text{on } G(x,y,t) \quad (\text{III-6})$$

where

$G(x,y,t)$ describes the surface of the body as a function of time.

For a thin airfoil restricted to small oscillations, this may be written as

$$G = y - H(x,t) \quad (\text{III-7})$$

where

$H(x,t)$ is the function describing the position of the airfoil.

$H(x,t)$ can be written for harmonic pitch oscillations as

$$H(x,t) = \text{R.P.}[\alpha_0(x-x_0) e^{i\nu t}] \quad (\text{III-8})$$

where the time-varying angle of attack $\alpha(t)$ is given by

$$\alpha(t) = \text{R.P.}[\alpha_0 e^{i\nu t}]$$

and α_0 = maximum amplitude of pitch oscillation

x_0 is the pitch axis

ν is the angular frequency of oscillation

$$i = \sqrt{-1}$$

R.P. = "real part of"

Inserting (III-8) into the flow tangency condition (III-6) gives, after linearization,

$$\phi_y(x,0) = v(x,0) = \alpha_0[U_0 + i\nu(x - x_0)] e^{i\nu t} \quad (\text{III-9})$$

$$\text{on } y = 0$$

This is a condition for the normal velocity to be prescribed at the airfoil's mean position $y = 0$.

C. NONDIMENSIONALIZATION

The terms in equations (III-5) and (III-9) are dimensional. For the following calculations it is convenient to use non-dimensional quantities. Define non-dimensional time and length to be

$$\bar{x} = \frac{x}{c}$$

$$\bar{y} = \frac{y}{c} \quad (\text{III-10})$$

$$\bar{t} = \frac{tU_o}{c}$$

where

U_o = uniform velocity from infinity

c = reference length (blade semichord).

The velocity potential in equation (III-5) may be non-dimensionalized as follows. Let

$$\bar{\phi} = \frac{\phi}{U_o c}$$

Hence:

$$\phi = U_o c \bar{\phi}$$

$$\phi_x = U_o c \bar{\phi}_x \left(\frac{1}{c}\right)$$

$$= U_o \bar{\phi}_x \quad (\text{III-11})$$

and similarly for the other derivatives in (III-5), yielding

$$[(M^2-1) + (\gamma+1)M^2\bar{\phi}_x] \phi_{\bar{x}\bar{x}} - \bar{\phi}_{\bar{y}\bar{y}} + 2M^2\bar{\phi}_{\bar{x}\bar{t}} + M^2\bar{\phi}_{\bar{t}\bar{t}} = 0$$

(III-12)

This equation is non-dimensional.

The boundary condition given in equation (III-9) may be non-dimensionalized in a similar fashion

$$\phi_y(x,0) = \alpha_o [U_o + i\nu(x - x_o)] e^{i\nu t} \quad (\text{III-9})$$

Thus

$$U_o c \bar{\phi}_y = \alpha_o [U_o + ik \frac{U_o}{c} (c\bar{x} - c\bar{x}_o)] e^{i\nu \bar{t} \frac{c}{U_o}} \quad (\text{III-13})$$

where

$$k = \frac{\nu c}{U_o} = \text{Strouhal number or reduced frequency}$$

$$U_o c \bar{\phi}_y \cdot \frac{1}{c} = \alpha_o U_o [1 + ik(\bar{x} - \bar{x}_o)] e^{ik\bar{t}} \quad (\text{III-14})$$

Thus

$$\bar{\phi}_y = \bar{v}(\bar{x},0) = \alpha_o [1 + ik(\bar{x} - \bar{x}_o)] e^{ik\bar{t}} \quad (\text{III-15})$$

Because the final operations are linear in α_o , set $\alpha_o = 1$, yielding

$$\bar{\phi}_y = \bar{v}(\bar{x},0) = [1 + ik(\bar{x} - \bar{x}_o)] e^{ik\bar{t}} \quad (\text{III-16})$$

The overbars denoting nondimensional quantities will be dropped from the remainder of the paper. All further quantities shall be assumed appropriately non-dimensional. This yields the following final equations

$$[(M^2-1)+(\gamma+1)M^2]\phi_{xx} - \phi_{yy} + 2M^2\phi_{xt} + M^2\phi_{tt} = 0 \quad (\text{III-17})$$

and

$$\phi_y(x,0) = v(x,0) = [1 + ik(x - x_0)] e^{ikt} \quad (\text{III-18})$$

where

all quantities are nondimensional and

$$\alpha_0 = 1$$

D. HARMONIC OSCILLATIONS

In the case of harmonic oscillations, equation (III-17) may be simplified still further.

Let

$$\phi = \tilde{\phi} + \text{R.P.}[\psi e^{ikt}]$$

where

$\tilde{\phi}$ = non-dimensional steady flow potential

ψ = non-dimensional oscillatory flow potential

R.P. = "real part of"

Equation (III-17) then becomes

$$\begin{aligned}
 (1-M^2)\psi_{xx} + \psi_{yy} - M^2(\gamma+1)[\psi_x\psi_{xx} + \tilde{\phi}_x\psi_{xx} + \tilde{\phi}_{xx}\psi_x] \\
 + M^2k^2\psi - 2iMk^2\psi_x = 0
 \end{aligned}
 \tag{III-19}$$

For M close to 1, this is a nonlinear mixed elliptic-hyperbolic partial differential equation with variable coefficients, the exact type depending on $\tilde{\phi}_x$ and $\tilde{\phi}_{xx}$. However, because flutter analysis is primarily concerned with the stability of small perturbations about a steady flow, the oscillatory component may be assumed small compared to the steady flow potential and therefore the product term $\psi_x\psi_{xx}$ may be neglected, yielding,

$$\begin{aligned}
 (1-M^2)\psi_{xx} + \psi_{yy} - M^2(\gamma+1)[\tilde{\phi}_x\psi_{xx} + \tilde{\phi}_{xx}\psi_x] \\
 = 2iM^2k\psi_x + M^2k^2\psi = 0
 \end{aligned}
 \tag{III-20}$$

IV. LINEARIZATION OF THE GOVERNING EQUATION

The basic flutter equation, (III-20), is still a non-linear, mixed elliptic-hyperbolic partial differential equation with variable coefficients and difficult to solve. It may yet be further simplified.

A. BASIC SOLUTION

For $M = 1$, equation (III-20) becomes

$$\psi_{yy} - (\gamma+1) [\tilde{\phi}_x \psi_{xx} + \tilde{\phi}_{xx} \psi_x] - 2ik\psi_x + k^2\psi = 0 \quad (\text{IV-1})$$

Now assume

$$\tilde{\phi}_x \approx w = \text{constant} \quad (\text{IV-2})$$

$$\tilde{\phi}_{xx} \approx 0$$

throughout the interblade channel. Setting

$$\tilde{\phi}_x(\gamma+1) = w(\gamma+1) = m^2 \quad (\text{IV-3})$$

yields

$$m^2\psi_{xx} - \psi_{yy} + 2ik\psi_x - k^2\psi = 0 \quad (\text{IV-4})$$

The solution to this equation is found in Garrick and Rubinow [5]

$$\psi(x,y) = -\frac{1}{m} \int_{x_\ell}^{x-my} v(s) J_0[\omega \sqrt{(x-s)^2 - (my)^2}] e^{i\lambda(s-x)} ds \quad \text{for } y > 0 \quad (\text{IV-5a})$$

and

$$\psi(x,y) = \frac{1}{m} \int_{x_\ell}^{x+my} v(s) J_0[\omega \sqrt{(x-s)^2 - (my)^2}] e^{i\lambda(s-x)} dx \quad \text{for } y < 0 \quad (\text{IV-5b})$$

where

$$v(x) = \lim_{y \rightarrow 0} \frac{\partial}{\partial y} \psi(x,y) .$$

$v(x)$ may be obtained directly from the tangential flow boundary condition, and

$$\omega = \frac{k^2(1-m^2)}{m}$$

$$\lambda = \frac{k}{m^2} \sqrt{1+m^2} \approx \frac{k}{m^2} \quad \text{(where this paper employs the approximation used by Gorelov [4])}$$

x_ℓ = blade leading edge,

Gorelov [4], has proposed a further simplification.

Set

$$z = my \quad (\text{IV-6a})$$

$$\Psi(x, z) = \psi(x, y) e^{i\lambda x} \quad (\text{IV-6b})$$

Equation (IV-4) then becomes

$$\Psi_{xx} - \Psi_{zz} + \omega^2 \Psi = 0$$

with solution

$$\Psi(x, z) = -\frac{1}{m} \int_{x_\ell}^{x-z} v(s) J_0[\omega \sqrt{(x-s)^2 - z^2}] e^{i\lambda s} ds \quad (\text{IV-7a})$$

$$z > 0$$

$$\Psi(x, z) = \frac{1}{m} \int_{x_\ell}^{x-z} v(s) J_0[\omega \sqrt{(x-s)^2 - z^2}] e^{i\lambda s} ds \quad (\text{IV-7b})$$

$$z < 0$$

where

$$v(x) = m e^{-i\lambda x} \lim_{z \rightarrow 0} \Psi_z(x, z)$$

$v(x)$ is obtained from the tangential flow boundary condition.

For a thin body immersed in the flow, the solutions for $y > 0$, $z > 0$, and $y < 0$, $z < 0$ apply above the body along left-running Mach lines, or below along right running Mach lines respectively.

B. BOUNDARY CONDITIONS

1. Flow Tangency Condition

The boundary condition comes from the tangential flow condition, (III-18)

$$v(x) = [1 + ik(x - x_0)] \quad (\text{IV-8})$$

2. Upstream Condition

The final linearized equation is a hyperbolic differential equation with boundary condition

$$\psi(x, y) = 0 \quad (\text{IV-9})$$

when

$$x - x_\ell < |my|$$

for the solution shown in equations (IV-5) or

$$\Psi(x, z) = 0$$

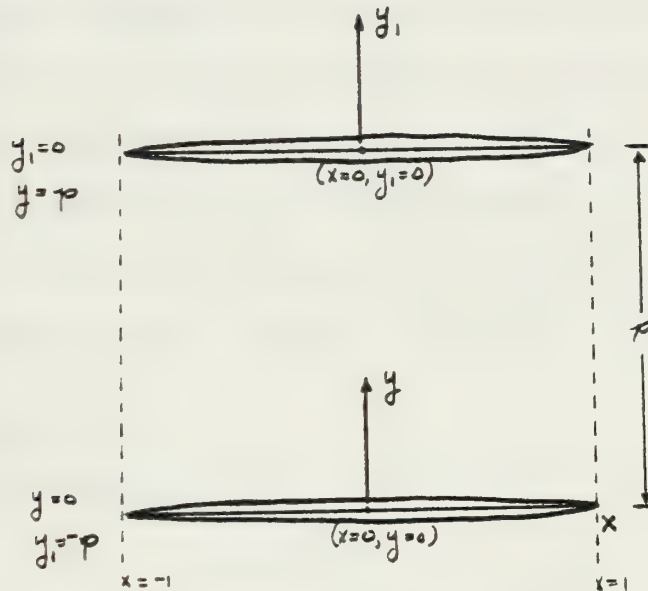
when

$$x - x_{\ell} < |z|$$

for the solution shown in equations (IV-7).

V. PROBLEM FORMULATION

A. CO-ORDINATE SYSTEM



Assume the geometry shown above. Both blades are thin airfoils of semichord c . All measurements are non-dimensional, normalized to c . The (x, y) co-ordinate system has its origin at the center of the reference (lower) blade. The (x, y_1) system is centered at the middle of the adjacent (upper) blade. The origin of the (x, y_1) system is located at $(0, p)$ in the reference system. Generalizing this convention, the same symbols shall be used for the same quantities on both blades. Where discrimination is required, the quantity associated with the adjacent blade will be marked with

superscript ¹, the quantity associated with the reference blade will be either unsuperscripted or marked with a superscript ⁰.

Each blade is assumed to perform a small amplitude harmonic oscillation about its mid-chord point. Both blades are assumed to have identical reduced frequencies, k , and the motion of the adjacent blade lags that of the reference blade by a phase angle σ .

The blades are immersed in a uniform flow from the left at $M = 1$. The objective is to determine the oscillatory pressure distributions and aerodynamic forces generated by the blades' oscillations. Cascade solidity, $\tau = 2/p$.

B. BOUNDARY CONDITIONS

1. Upstream Condition

$$\psi = 0 \quad \text{whenever}$$

$$x + 1 < |my| \tag{V-1}$$

and

$$x + 1 < |my_1|, \text{ simultaneously}$$

2. Flow Tangency Condition

Along the reference blade

$$\lim_{y \rightarrow 0} \psi_y(x, y) = (1 + ikx) \tag{V-2a}$$

Along the adjacent blade

$$\lim_{y_1 \rightarrow 0} \psi_{y_1}(x, y_1) = (1 + ikx)e^{i\sigma} \quad (V-2b)$$

where

σ is the phase angle between the blades oscillations

C. BASIC SOLUTION TECHNIQUE

Assume that the unsteady potential, ψ , may be written as the sum of four components

$$\psi(x, y) = \phi^0(x, y) + \psi^0(x, y) + \phi^1(x, y_1) + \psi^1(x, y_1) \quad (V-3)$$

where:

- ϕ^0 = potential due to the reference blade alone, known from equation (IV-7)
- ϕ^1 = potential due to the adjacent blade alone, known from equation (IV-7)
- ψ^0 = interference potential required to satisfy tangential flow condition along reference blade, unknown
- ψ^1 = interference potential required to satisfy tangential flow condition along adjacent blade, unknown.

This total potential must satisfy the tangential flow condition at the plane of both the reference and adjacent blades. Thus

$$\begin{aligned} \phi_Y^0(x, y=0) + \phi_{Y_1}^1(x, y_1=-p) + \psi_Y^0(x, y=0) + \psi_{Y_1}^1(x, y_1=-p) \\ = (1 + ikx) \end{aligned} \quad (V-4a)$$

at the reference blade, and

$$\begin{aligned} \phi^0(x, y=p) + \phi_{y_1}^1(x, y_1=0) + \psi_Y^0(x, y=p) + \psi_{y_1}^1(x, y_1=0) \\ = (1 + ikx) e^{i\sigma} \end{aligned} \quad (V-4b)$$

at the adjacent blade.

But from the unsteady potential solution for a single oscillating blade one has

$$\phi_Y^0(x, y=0) = 1 + ikx \quad (V-5a)$$

and

$$\phi_{y_1}^1(x, y_1=0) = (1 + ikx) e^{i\sigma} \quad (V-5b)$$

Thus

$$\phi_{y_1}^1(x, y_1=-p) + \psi_{y_1}^1(x, y_1=-p) + \psi_Y^0(x, y=0) = 0 \quad (V-6a)$$

along the reference blade, and

$$\phi_Y^0(x, y=p) + \psi_Y^0(x, y=p) + \psi_{y_1}^1(x, y_1=0) = 0 \quad (V-6b)$$

along the adjacent blade.

From equation (IV-7)

$$\phi^0(x, y) = -\frac{1}{m} \int_{-1}^{x-my} v^0(s) J_0[\omega \sqrt{(x-s)^2 - (my)^2}] e^{i\lambda(s-x)} ds$$

$$y > 0 \quad (V-7a)$$

$$= \frac{1}{m} \int_{-1}^{x+my} v^0(s) J_0[\omega \sqrt{(x-s)^2 - (my)^2}] e^{i\lambda(s-x)} ds$$

$$y < 0 \quad (V-7b)$$

where

$$v^0(s) = 1 + iks$$

$$\lambda = k/m^2$$

$$m = (\gamma+1)w$$

$$\omega = \frac{k^2(1-m^2)}{m^4}$$

$$w = \text{mean value of } \tilde{\phi}_x \text{ in the channel}$$

$$\phi^1(x, y_1) = -\frac{1}{m} \int_{-1}^{x-my_1} v^1(s) J_0[\omega \sqrt{(x-s)^2 - (my_1)^2}] e^{i\lambda(s-x)} ds$$

$$y_1 > 0 \quad (IV-8a)$$

$$= \frac{1}{m} \int_{-1}^{x+my_1} v^1(s) J_0[\omega \sqrt{(x-s)^2 - (my_1)^2}] e^{i\lambda(s-x)} ds$$

$$y_1 < 0 \quad (IV-8b)$$

where

$$v^1(s) = (1 + iks)e^{i\sigma}$$

Henceforth attention will be restricted to the flow within the channel, $0 \leq y \leq p$, $-p \leq y_1 \leq 0$ leaving (IV-7a) and (IV-8b) as the governing equations of interest.

The two interference potentials are assumed to have forms identical to (IV-7a) and (IV-8b).

Set

$$\psi^0(x, y) = -\frac{1}{m} \int_{-1}^{x-my} u^0(s) J_0[\omega \sqrt{(x-s)^2 - (my)^2}] e^{i\lambda(s-x)} ds$$

(V-9a)

$y > 0$

$$\psi^1(x, y_1) = \frac{1}{m} \int_{-1}^{x+m} u^1(s) J_0[\omega \sqrt{(x-s)^2 - (my_1)^2}] e^{i\lambda(s-x)} ds$$

(V-9b)

$y_1 < 0$

where

$u^0(s)$ and $u^1(s)$ are unknown functions to be determined so as to satisfy equations (V-6)

Substitution of (V-7a), (V-8b) and (V-9) into (V-6) yields

$$u^0(x) + \psi_{y_1}^1(x, y_1 = -p) + \phi_{y_1}^0(x, y_1 = -p) = 0 \quad (V-10a)$$

$$u^1(x) + \psi_Y^0(x, y=p) + \phi_Y^0(x, y=p) = 0 \quad (V-10b)$$

Recalling Gorelov's transformation discussed in [4] and shown in equations (IV-6) above, set

$$\phi^0(x, z) = \phi^0(x, y) e^{i\lambda x}$$

$$\phi^1(x, z_1) = \phi^1(x, y_1) e^{i\lambda x}$$

$$\psi^0(x, z) = \psi^0(x, y) e^{i\lambda x}$$

$$\psi^1(x, z_1) = \psi^1(x, y_1) e^{i\lambda x}$$

where $z = my$

$$z_1 = my_1$$

Then

$$\frac{e^{i\lambda x}}{m} u^0(x) + \phi_{z_1}^1(x, z_1 = -x_*) + \psi_{z_1}^1(x, z_1 = -x_*) = 0 \quad (V-11a)$$

and

$$\frac{e^{i\lambda x}}{m} u^1(x) + \phi_z^0(x, z = x_*) + \psi_z^0(x, z = x_*) = 0 \quad (V-11b)$$

where

$$x_* = mp.$$

To employ the collocation method, assume that $u^1(x)$ and $u^0(x)$ can be approximated as the sum of a set of elementary functions f_j so that

$$u^0(x) \approx \sum_{j=0}^n \theta_j^0 f_j(x) \quad (V-12a)$$

$$u^1(x) \approx \sum_{j=0}^n \theta_j^1 f_j(x) \quad (V-12b)$$

where $f_j(x) = 0$ when $x \leq x_* - 1$

Note that here both u^0 and u^1 are expressed in terms of the same elementary functions, f_j .

Because of the slightly supersonic nature of the problem observe that $u^0(x) = 0$ and $u^1(x) = 0$ when $x \leq x_* - 1$.

Equations (V-12) may now be rewritten as

$$\begin{aligned} e^{i\lambda x} \sum_{j=0}^n \theta_j^0 f_j(x) + \frac{\partial}{\partial z_1} \int_{x_*-1}^{x+z_1} \sum_{j=0}^n \theta_j^1 f_j(s) J_0[\omega \sqrt{(x-s)^2 - z_1^2}] e^{i\lambda s} ds \\ = -e^{i\sigma} \frac{\partial}{\partial z_1} \int_{-1}^{x+z_1} (1+iks) J_0[\omega \sqrt{(x-s)^2 - z_1^2}] e^{i\lambda s} ds \end{aligned}$$

at $z_1 = -x_*$ (V-13a)

$$\begin{aligned}
e^{i\lambda x} \sum \theta_j^1 f_j(x) - \frac{\partial}{\partial z} \int_{x_*-1}^{x-z} \sum \theta_j^0 f_j(s) J_0[\omega \sqrt{(x-s)^2 - z^2}] e^{i\lambda s} ds \\
= \frac{\partial}{\partial z} \int_{-1}^{x+z} (1+iks) J_0[\omega \sqrt{(x-s)^2 - z^2}] e^{i\lambda s} ds
\end{aligned}$$

$$\text{at } z = x_* \quad (V-13b)$$

where

$$f_j(x) = 0 \text{ for } x \leq x_* - 1$$

This simplifies to

$$\begin{aligned}
e^{i\lambda x} \sum \theta_j^0 f_j(x) + \sum \theta_j^1 \left\{ \frac{\partial}{\partial z_1} \int_{x_*-1}^{x+z_1} f_j(s) J_0[\omega \sqrt{(x-s)^2 - z_1^2}] e^{i\lambda s} ds \right\} \\
= -e^{i\sigma} \frac{\partial}{\partial z_1} \int_{-1}^{x+z_1} (1+iks) J_0[\omega \sqrt{(x-s)^2 - z_1^2}] e^{i\lambda s} ds
\end{aligned}$$

$$\text{at } z_1 = -x_* \quad (V-14a)$$

and

$$\begin{aligned}
e^{i\lambda x} \sum \theta_j^1 f_j(x) - \sum \theta_j^0 \left\{ \frac{\partial}{\partial z} \int_{x_*-1}^{x-z} f_j(s) J_0[\omega \sqrt{(x-s)^2 - z^2}] e^{i\lambda s} ds \right\} \\
= \frac{\partial}{\partial z} \int_{-1}^{x-z} (1+iks) J_0[\omega \sqrt{(x-s)^2 - z^2}] e^{i\lambda s} ds \quad (V-14b)
\end{aligned}$$

at $z = x_*$ where

$$f_j(x) = 0 \text{ for } x \leq x_* - 1.$$

Performing the indicated differentiation yields

$$\begin{aligned} & e^{i\lambda x} \left\{ \theta_j^0 f_j(x) + \left[\theta_j^1 \left\{ \int_{x_*-1}^{x-x_*} f_j(s) \frac{J_1[\omega \sqrt{(x-s)^2 - x_*^2}] \omega x_* e^{i\lambda s} ds}{\sqrt{(x-s)^2 - x_*^2}} \right. \right. \right. \\ & \quad \left. \left. \left. + f_j(x-x_*) e^{i\lambda(x-x_*)} \right\} \right] \right\} \\ & = e^{i\sigma} \left\{ - \int_{-1}^{x-x_*} (1+iks) \frac{\omega x_* J_1[\omega \sqrt{(x-s)^2 - x_*^2}] e^{i\lambda s} ds}{\sqrt{(x-s)^2 - x_*^2}} \right. \\ & \quad \left. - [1+ik(x-x_*)] e^{i\lambda(x-x_*)} \right\} \end{aligned} \quad (V-15a)$$

and

$$\begin{aligned} & e^{i\lambda x} \left\{ \theta_j^1 f_j(x) + \left[\theta_j^0 \left\{ \int_{x_*-1}^{x-x_*} f_j(s) \frac{J_1[\omega \sqrt{(x-s)^2 - x_*^2}] \omega x_* e^{i\lambda s} ds}{\sqrt{(x-s)^2 - x_*^2}} \right. \right. \right. \\ & \quad \left. \left. \left. + f_j(x-x_*) e^{i\lambda(x-x_*)} \right\} \right] \right\} \\ & = - \int_{-1}^{x-x_*} (1+iks) \frac{\omega x_* J_1[\omega \sqrt{(x-s)^2 - x_*^2}] e^{i\lambda s} ds}{\sqrt{(x-s)^2 - x_*^2}} - [1+ik(x-x_*)] e^{i\lambda(x-x_*)} \end{aligned} \quad (V-15b)$$

Gorelov's formulation, equations [2.6, 2.7, 2.8, and 2.9] of [4], can be obtained directly from equations (V-15) by substituting

$$f_j(x) = \cos j\eta - \cos j\eta_*$$

$$v^0 = \sum_{j=0}^n \mu_j^0 \cos j\eta$$

$$v^1 = \sum_{j=0}^n \mu_j^1 \cos j\eta$$

where:

$$\eta = \cos(-x)$$

$$\eta_* = \cos(1-x_*)$$

In comparing the two systems care must be taken to note the differing symbols and coordinate systems. The corresponding quantities are:

Here	in[4]
θ_j^0, θ_j^1	$v_{0\sigma}, v_{i\sigma}$
μ_j^0, μ_j^1	$\theta_{0\sigma}, \theta_{i\sigma}$
$\eta = \cos^{-1}(-x)$	$\eta = \cos^{-1}(1-x)$
$\eta_* = \cos^{-1}(1-x_*)$	$\eta_* = \cos^{-1}(1-x_*)$
z, z_1	y, y_1
j	σ
σ	ψ

Here $-1 \leq x \leq 1$; in [4] $0 \leq x \leq 2$. This transformation accounts for the differing definitions of η . Making the substitutions results in the system

n

$$\sum_{j=0}^n \{ \theta_j^0 [\cos j\eta - (1-\delta_{01}) \cos j\eta_*] \\ + \theta_j^1 \int_{x_*-1}^{x-x_*} \frac{\partial}{\partial z_1} J_0 [\omega \sqrt{(x-s)^2 - z_1^2}] [\cos j\hat{\eta} - (1-\delta_{10}) \cos j\eta_*] e^{i\lambda s} ds$$

$$+ \theta_i^1 [\cos j\bar{\eta} - (1-\delta_{10}) \cos j\eta_*] e^{i\lambda(x-x_*)} \}$$

$$= - \sum_{j=0}^n \{ -\mu_j^1 \int_{-1}^{x-x_*} \frac{\partial}{\partial z_1} J_0 [\sqrt{(x-s)^2 - z_1^2}] \cos j\hat{\eta} e^{i\lambda s} ds \\ - \mu_j^1 \cos j\bar{\eta} e^{i\lambda(x-x_*)} \}$$

$$(V-16a)$$

$$\text{at } z_1 = -x_*$$

$$x > x_* - 1$$

and

$$\sum \{ \theta_j^1 [\cos j\eta - (1-\delta_{oj}) \cos j\eta_*]$$

$$+ \theta_j^0 \int_{x_*-1}^{x-x_*} \frac{\partial}{\partial z} J_0 [\omega \sqrt{(x-s)^2 - z^2}] [\cos j\eta - (1-\delta_{oj}) \cos j\eta_*] e^{i\lambda s} ds$$

$$+ \theta_j^0 [\cos j\bar{\eta} - (1-\delta_{oj}) \cos j\eta_*] e^{i\lambda(x-x_*)} \}$$

$$= \sum_{j=0}^n \{ \mu_j^0 \int_{-1}^{x-x_*} \frac{\partial}{\partial z} J_0 [\omega \sqrt{(x-s)^2 - z^2}] \cos j\hat{\eta} e^{i\lambda s} ds$$

$$- \mu_j^0 \cos j\bar{\eta} e^{i\lambda(x-x_*)} \} \quad (V-16b)$$

where:

$$\eta = \arccos(-x)$$

$$\eta_* = \arccos(1-x_*)$$

$$\bar{\eta} = \arccos(-s)$$

$$\hat{\eta} = \arccos(-x+x_*)$$

$$\delta_{\bar{o}j} = \text{Dirac } \delta \text{ function} = \begin{matrix} 1 & \text{for } j = 0 \\ 0 & \text{for } j \neq 0 \end{matrix}$$

Given the change in coordinates and notation, this system is equivalent to that shown in [4].

This was the system programmed for computer solution. Because the function, $f_j(x)$, is unaffected by the differentiation with respect to y (or z) the exact form used need not be specified, so that the system shown in (V-15) may be programmed with f undetermined. A subroutine may be written to return the function desired and the remaining program left perfectly general. In the program developed with this thesis both the Gorelov functions shown above and the Legendre polynomials were employed. All the integrals may now be evaluated at $n+1$ points, x_i , on both blades in $mp-1 < x_i < 1$ and the resulting linear system solved for θ_j^0 and θ_j^1 , $j=1,2,\dots,n+1$.

The interference potentials may be constructed by taking

$$\psi^0(x,z) = \frac{-1}{m} \int_1^{x-z} [\theta_j^0 f_j(x)] J_0[\omega \sqrt{(x-s)^2 - z^2}] e^{i\lambda s} ds$$

$$z > 0 \quad (V-17a)$$

$$\psi^1(x,z_1) = \frac{1}{m} \int_{-1}^{x+z_1} [\theta_j^1 f_j(x)] J_0[\omega \sqrt{(x-s)^2 - z_1^2}] e^{i\lambda s} ds$$

$$z_1 < 0 \quad (V-17b)$$

where

$$f_j(x) = 0, \quad \text{for all } x \leq x_* - 1$$

Once the potentials have been calculated as outlined above, the surface pressure may be calculated using the relationship

$$C_p = -2(\psi_x + ik\psi) \quad (V-18a)$$

$$= -2[\Psi_x + i(k-\lambda)\Psi] e^{-i\lambda x} \quad (V-18b)$$

Because all the plates are assumed to be in steady oscillation with uniform phase shift, σ , between neighboring plates, then

$$v^1(x) = v^0(x)e^{i\sigma}, \quad u^1(x) = u^0(x)e^{i\sigma}, \quad \Psi(x,y) = -\Psi(x,-y)$$

in this case

$$C_{\ell\alpha} = 2 \int_{-1}^1 [\Psi_x(x,+0) + i(k-\lambda)\Psi(x,+0)] e^{-i\lambda x} dx \quad (V-19)$$

where

$$\begin{aligned} \Psi(x_1,+0) &= \frac{-1}{m} \int_{-1}^x [(1+iks)+u^0(s)] J_0[\omega(x-s)] e^{i\lambda(s)} ds \\ &+ \frac{e^{i\sigma}}{m} \int_{-1}^{x-x_*} [(1+iks)+u^0(s)] J_0[\omega\sqrt{(x-s)^2-(x_*)^2}] e^{i\lambda(s)} ds \end{aligned}$$

where $x_* = mp$. (V-20)

Results from this approach, in the form of values of $C_{\ell\alpha}$ for $k = 0.1$, at various values of w are presented in the results for

approximations based both on Gorelov's formulation, and on the Legendre polynomials.

D. COLLOCATION SOLUTION OF THE POTENTIAL EQUATION EXPANDED FOR SMALL k ,

In order to provide a partially independent check of the results of the main program, the Gorelov function representation of the collocation solution was expanded for small k , and solved at two collocation points, $n = 2$. The resulting potentials, and partial derivatives with respect to x and y were then used to replace the corresponding numerical routines in the main program. The output resulting from the approximations were compared with the purely numerical results obtained from the computer program.

1. Solution For The Unknown Potential Coefficients

The basic system of linear equations used to determine the unknown coefficients is

$$\frac{1}{m} e^{i\lambda x} u^0 + \phi_{z_1}^1 + \psi_{z_1}^1 = 0, \quad z = 0, \quad z_1 = -x_* = -mp \quad (V-21a)$$

$$\frac{1}{m} e^{i\lambda x} u^1 + \phi_z^0 + \psi_z^0 = 0, \quad z_1 = 0, \quad z = x_* = mp \quad (V-21b)$$

where

$$0 = u^0(x) = u^1(x) \quad \text{when} \quad x \leq -1+x_*$$

otherwise

$$u^0 = \sum_{j=1}^n \theta_j^0 (\cos j\eta - \cos j\eta_*) + \theta_0^0$$

$$u^1 = \sum_{j=1}^n \theta_j^1 (\cos j\eta - \cos j\eta_*) + \theta_0^1$$

where

$$\eta = \arccos(-x)$$

$$\eta_* = \arccos(1 - x_*)$$

Thus, for $n = 2$, the system becomes

$$e^{i\lambda x} \{ \theta_0^0 + \theta_1^0 (\cos \eta - \cos \eta_*) \}$$

$$+ \frac{\partial}{\partial z_1} \int_{-1}^{x+z_1} v^1(s) J_0[\omega \sqrt{(x-s)^2 - z_1^2}] e^{i\lambda s} ds$$

$$+ \frac{\partial}{\partial z_1} \int_{x_*-1}^{x+z_1} u^1(s) J_0[\omega \sqrt{(x-s)^2 - z_1^2}] e^{i\lambda s} ds = 0$$

$$z_1 = -x_*$$

(V-22a)

along the reference blade and

$$e^{i\lambda x} \{ \theta_0^1 + \theta_1^1 (\cos \eta - \cos \eta_*) \} - \frac{\partial}{\partial z} \int_{-1}^{x-z} v^0(s) J_0 [\omega \sqrt{(x-s)^2 - z^2}] e^{i\lambda s} ds$$

$$- \frac{\partial}{\partial z} \int_{x_*-1}^{x-z} u^0(s) J_0 [\omega \sqrt{(x-s)^2 - z^2}] e^{i\lambda s} ds \quad (V-22b)$$

$z = x_*$

along the adjacent blade

where $u^0(s) = \theta_0^0 + \theta_1^0 (\cos \eta - \cos \eta_*)$

$$u^1(s) = \theta_0^1 + \theta_1^1 (\cos \eta - \cos \eta_*)$$

$$v^0(s) = 1 + iks$$

$$v^1(s) = (1 + iks) e^{i\sigma}$$

For k sufficiently small, this system may be further simplified by the following approximations

$$J_0 [\omega \sqrt{(x-s)^2 - z^2}] \approx 1 - O(\omega^2) \approx 1 \quad (V-23a)$$

$$J_0 [\omega \sqrt{(x-s)^2 - z^2}] \approx 1 - O(\omega^2) \approx 1 \quad (V-23b)$$

$$e^{i\lambda x} \approx 1 + i\lambda x - O(\lambda^2 x^2) \approx 1 + i\lambda x$$

$$e^{i\lambda s} \approx 1 + i\lambda s - O(\lambda^2 s^2) \approx 1 + i\lambda s$$

where $O(\omega^2)$ means "of the order of magnitude of ω^2 "

$$-1 \leq x \leq 1, \quad -1 \leq s \leq x - x_*$$

$$\lambda = k/m^2, \quad \omega^2 = \frac{k^2(1-m^2)}{m^4}$$

The interference source distributions may be replaced by

$$u^0(s) = \theta_0^0 + \theta_1^0(-s + x_* - 1)$$

$$u^1(s) = \theta_0^1 + \theta_1^1(-s + x_* - 1) .$$

If higher order terms are neglected, the result is a system linear in k and λ

$$\begin{aligned} (1+i\lambda x) [\theta_0^0 + \theta_1^0(-x-1+x_*)] + \frac{\partial e^{i\sigma}}{\partial z_1} \int_{-1}^{x+z_1} (1+iks)(1+i\lambda s) ds \\ + \frac{\partial}{\partial z_1} \int_{x_*-1}^{x+z_1} [\theta_0^1 + \theta_1^1(-s-1+x_*)] (1+i\lambda s) ds = 0 \\ z_1 = -x_* \end{aligned} \quad (V-24a)$$

$$\begin{aligned} (1+i\lambda x) [\theta_0^1 + \theta_1^1(-x-1+x_*)] - \frac{\partial}{\partial z} \int_{-1}^{x-z} (1-iks)(1+i\lambda s) ds \\ - \frac{\partial}{\partial z} \int_{x_*-1}^{x-z} [\theta_0^0 + \theta_1^0(-s-1+x_*)] (1+i\lambda s) ds = 0 \\ z = x_* \end{aligned} \quad (V-24b)$$

Product terms containing $(k\lambda) = \frac{k^2}{m^2}$ may be neglected as of higher order in k , yielding

$$(1+i\lambda x) [\theta_0^0 + \theta_1^0(-x-1+x_*)] + \frac{\partial e^{i\sigma}}{\partial z_1} \int_{-1}^{x+z_1} [1+i(\lambda+k)s] ds$$

$$+ \frac{\partial}{\partial z_1} \int_{x_*-1}^{x+z_1} \{ [\theta_0^1 + \theta_1^1(-s-1+x_*)] \}$$

$$+ i\lambda s [\theta_0^1 + \theta_1^1(-s-1+x_*)] \} ds = 0$$

$$z_1 = -x_* \quad (V-25a)$$

$$(1+i\lambda x) [\theta_0^1 + \theta_1^1(-x-1+x_*)] - \frac{\partial}{\partial z} \int_{-1}^{x-z} [1+i(\lambda+k)s] ds$$

$$- \frac{\partial}{\partial z} \int_{x_*-1}^{x-z} \{ [\theta_0^0 + \theta_1^0(-s-1+x_*)] + i\lambda s [\theta_0^1 + \theta_1^1(-s-1+x_*)] \} ds = 0$$

$$z = x_* \quad (V-25b)$$

Evaluating the indicated derivatives yields

$$(1+i\lambda x) [\theta_0^0 + \theta_1^0(-x-1+x_*)] + [1+i(k+\lambda)(x-x_*)] e^{i\sigma}$$

$$+ [\theta_0^1 + \theta_1^1(2x_*-x-1)] + i\lambda(x-x_*) [\theta_0^1 + \theta_1^1(2x_*-x-1)] \} = 0$$

$$(V-26a)$$

$$(1+i\lambda x) [\theta_0^1 + \theta_1^1 (-x-1+x_*)] + [1+i(\lambda+k)(x-x_*)]$$

$$+ \{ [\theta_0^0 + \theta_1^0 (2x_* - x - 1)] + i\lambda(x-x_*) [\theta_0^0 + \theta_1^0 (2x_* - x - 1)] \} = 0$$

(V-26b)

Thus:

$$\theta_0^0 (1+i\lambda x) + \theta_1^0 [(-x-1+x_*) + i\lambda x (-x-1+x_*)]$$

$$+ \theta_0^1 [1 + i\lambda(x-x_*)] + \theta_1^1 [(2x_* - x - 1) + i\lambda(x-x_*)(2x_* - x - 1)]$$

$$= -e^{i\sigma} [1 + i(k+\lambda)(x-x_*)] \quad (V-27a)$$

$$\theta_0^1 [1 + i\lambda x] + \theta_1^1 [(-x+1-x_*) + i\lambda x (-x+1-x_*)]$$

$$+ \theta_0^0 [1+i\lambda x(x-x_*)] + \theta_1^0 [(2x_* - x - 1) + i\lambda(x-x_*)(2x_* - x - 1)]$$

$$= -[1 + i(\lambda+k)(x-x_*)] \quad (V-27b)$$

This system may be solved at two points, x_1 and x_2 , for

θ_0^0 , θ_1^0 , θ_0^1 , and θ_1^1 .

2. Calculation Of The Potential

The potential is given by

$$\psi(x, y) = \psi(x, z) e^{-i\lambda x} \quad (V-28)$$

where

$$\begin{aligned} \Psi(x, z) = & -\frac{1}{m} \int_{-1}^x [v^0(s) + u^0(s)] J_0[\omega \sqrt{(x-s)^2}] e^{i\lambda s} ds \\ & + \frac{1}{m} \int_{-1}^{x-x_*} [v^1(s) + u^1(s)] J_1[\omega \sqrt{(x-s)^2 - x_*^2}] e^{i\lambda s} ds. \end{aligned}$$

$$u^0(s) = u^1(s) = 0 \quad \text{for all } s \leq x_* - 1.$$

Thus

$$\begin{aligned} \Psi(x, z) = & -\frac{1}{m} \int_{-1}^x v^0(s) J_0[\omega \sqrt{(x-s)^2}] e^{i\lambda s} ds \\ & - \frac{1}{m} \int_{x_*-1}^x u^0(s) J_0[\omega \sqrt{(x-s)^2}] e^{i\lambda s} ds \\ & + \frac{1}{m} \int_{-1}^{x-x_*} v^1(s) J_0[\omega \sqrt{(x-s)^2 - x_*^2}] e^{i\lambda s} ds \\ & + \frac{1}{m} \int_{x_*-1}^{x-x_*} u^1(s) J_0[\omega \sqrt{(x-s)^2 - x_*^2}] e^{i\lambda s} ds \quad (V-29) \end{aligned}$$

Making the same small frequency approximations as in the previous section yields

$$\begin{aligned}
 \Psi(x, z) = & -\frac{1}{m} \int_{-1}^x v^0(s) (1+i\lambda s) ds \\
 & -\frac{1}{m} \int_{x_*-1}^x u^0(s) (1+i\lambda s) ds \\
 & + \frac{1}{m} \int_{-1}^{x-x_*} v^1(s) (1+i\lambda s) ds \\
 & + \frac{1}{m} \int_{x_*-1}^{x-x_*} u^1(s) (1+i\lambda s) ds \quad (V-30)
 \end{aligned}$$

From the general formulation

$$\Psi = \Phi^1 + \Phi^0 + \Psi^1 + \Psi^0 \quad (V-31)$$

Thus, along the reference blade

$$\begin{aligned}
 -m\Phi^0(x, z=0) &= \int_{-1}^x v^0(s) (1+i\lambda s) ds = \int_{-1}^x (1+iks) (1+i\lambda s) ds \\
 &= \int_{-1}^x [1+i(k+\lambda)s] ds = \left[s + i(k+\lambda)\frac{s^2}{2} \right]_{-1}^x \quad (V-32) \\
 &= x + i\frac{(k+\lambda)}{2}x^2 + 1 - i\left(\frac{k+\lambda}{2}\right)
 \end{aligned}$$

$$\phi^0(x, z=0) = -\frac{1}{m}[(1+x) + i(\frac{k+\lambda}{2})(x^2-1)] \quad (V-33)$$

By inspection

$$\phi^1(x, z_1=-x_*) = \frac{e^{i\sigma}}{m}\{(1+x-x_*) + i(\frac{k+\lambda}{2})[(x-x_*)^2-1]\} \quad (V-34)$$

$$-m\psi^0(x, z=0) = \int_{x_*-1}^x u^0(s)(1+i\lambda s)ds \quad (V-35)$$

$$= \int_{x_*-1}^x [\theta_0^0 + \theta_1^0(-s+x_*-1)](1+i\lambda s)ds$$

$$= \int_{x_*-1}^x \theta_0^0(1+i\lambda s) + \theta_1^0(-s+x_*-1)(1+i\lambda s)ds$$

$$= \theta_0^0[(x+1-x_*) + \frac{i\lambda}{2}(x^2-x_*^2+2x_*-1)]$$

$$+ \int_{x_*-1}^x \theta_1^0[-x+x_*-1+i\lambda(-s^2+sx_*-s)]ds$$

$$= \theta_0^0[(x+1-x_*) + \frac{i\lambda}{2}(x^2-x_*^2+2x_*-1)]$$

$$+ \theta_1^0\{(-\frac{s^2}{2}+sx_*-s) + i\lambda(-\frac{s^3}{3} + \frac{s^2x_*}{2} - \frac{s^2}{2})\} \Big|_{x_*-1}^x \quad (V-36)$$

$$\begin{aligned}
&= \theta_0^0 \left[(x+1-x_*) + \frac{i\lambda}{2} (x^2 - x_*^2 + 2x_* - 1) \right] \\
&+ \theta_1^0 \left\{ \left(-\frac{x^2}{2} + x_* - x \right) - \left[-\frac{(x_*-1)^2}{2} + x_* (x_*-1) - (x_*-1) \right] \right. \\
&+ i\lambda \left[\left(-\frac{x^3}{3} + \frac{x^2 x_*}{2} - \frac{x^2}{2} \right) + \frac{(x_*-1)^3}{3} - x_* \frac{(x_*-1)^2}{2} + \frac{(x_*-1)^2}{2} \right] \left. \right\}
\end{aligned}$$

$$\begin{aligned}
&= \theta_0^0 \left[(x+1-x_*) + \frac{i\lambda}{2} (x^2 - x_*^2 + 2x_* - 1) \right] \\
&+ \theta_1^0 \left\{ -\frac{x^2}{2} + x(x_*-1) + \frac{(x_*-1)^2}{2} - (x_*-1)^2 \right. \\
&+ i\lambda \left[-\frac{x^3}{3} + x^2 \frac{(x_*-1)}{2} + \frac{x_*^3}{3} - x_*^2 + x_* - \frac{1}{3} \right. \\
&\quad \left. \left. - \frac{x_*^3}{2} + x_*^2 - \frac{x_*}{2} + \frac{x_*^2}{2} - x_* + \frac{1}{2} \right] \right\}
\end{aligned}$$

$$\begin{aligned}
&= \theta_0^0 \left[(x+1-x_*) + \frac{i\lambda}{2} (x^2 - x_*^2 + 2x_* - 1) \right] \\
&+ \theta_1^0 \left\{ -\frac{x^2}{2} + x(x_*-1) - \frac{(x_*-1)^2}{2} \right. \\
&+ i\lambda \left[-\frac{x^3}{3} + \frac{x^2 (x_*-1)}{2} - \frac{x_*^3}{6} + \frac{x_*^2}{2} - \frac{x_*}{2} + \frac{1}{6} \right] \left. \right\}
\end{aligned}$$

$$\begin{aligned}
\Psi^0 &= -\frac{1}{m}\{\theta_0^0[(x+1-x_*) + \frac{i\lambda}{2}(x^2 + x_*^2 + 2x_*-1)] \\
&\quad + \theta_1^0\{-\frac{x^2}{2} + x(x_*-1) - \frac{(x_*-1)^2}{2} \\
&\quad + i\lambda[-\frac{x^3}{3} + \frac{x^2(x_*-1)}{2} - \frac{(x_*-1)^3}{6}]\}\} \quad (V-37)
\end{aligned}$$

$m\Psi^1(x, z_1=-x_*)$ may be evaluated by substituting $x-x_*$ for x in the expression for $m\Psi^0$ and exchanging θ_0^1 and θ_1^1 for θ_0^0 and θ_1^0

$$\begin{aligned}
m\Psi^1(x, z_1=-x_*) &= \theta_0^1\{[(x-x_*)+1-x_*] + \frac{i\lambda}{2}[(x-x_*)^2 - x_*^2 + 2x_* - 1]\} \\
&\quad + \theta_1^1\{-\frac{(x-x_*)^2}{2} + (x-x_*)(x_*-1) - \frac{(x_*-1)^2}{2} \\
&\quad + i\lambda[-\frac{(x-x_*)^3}{3} + \frac{x^2(x_*-1)}{2} - \frac{(x_*-1)^3}{6}]\} \quad (V-38)
\end{aligned}$$

$$\begin{aligned}
&= \theta_0^1(x+1-2x_*) + \frac{i\lambda}{2}[(x^2-2xx_*+x_*^2) - x_*^2 + 2x_* - 1] \\
&\quad + \theta_1^1\{-\frac{(x^2-2xx_*+x_*^2)}{2} + (x - \frac{3}{2}x_* + \frac{1}{2})(x_*-1) \\
&\quad + i\lambda[-\frac{x^3}{3} + x^2x_* - xx_*^2 + \frac{x_*^3}{3} + (x^2-2xx_*+x_*^2)\frac{(x_*-1)}{2} \\
&\quad \quad - \frac{(x_*-1)^3}{6}]\}
\end{aligned}$$

$$\begin{aligned}
&= \theta_0^1 \{ (x+1-2x_*) + \frac{i\lambda}{2} [x^2 - 2xx_* + 2x_* - 1] \\
&\quad + \theta_1^1 \{ -\frac{x^2}{2} + xx_* - \frac{x_*^2}{2} + x(x-1) + \frac{1}{2}(1-3x_*)(x_*-1) \\
&\quad + i\lambda [-\frac{x^3}{3} + x^2 \frac{(3x_*-1)}{2} + x(-\frac{x_*^2-2xx_*^2-1)}{2} \\
&\quad \quad + \frac{x_*^3}{3} + \frac{x_*^3}{2} - \frac{1}{2}] \}
\end{aligned}$$

$$\psi^1(x, z_1 = x_*)$$

$$\begin{aligned}
&= \frac{1}{m} \{ \theta_0^1 \{ (x+1-2x_*) + \frac{i\lambda}{2} [x^2 - 2xx_* + 2x_* - 1] \\
&\quad + \theta_1^1 \{ [-\frac{x^2}{2} + x(2x_*-1) - xx_*^2 + 2x_* - \frac{1}{2}] \\
&\quad + i\lambda [-\frac{x^3}{3} + x^2 \frac{(3x_*-1)}{2} + x \frac{(-3x_*^2-1)}{2} + 5\frac{x_*^3}{6} - \frac{1}{2}] \} \}
\end{aligned}$$

(V-39)

A comparison of the results for the full program and the approximation is given below for $k = 0.01$, $w = 0.05$, $\sigma = \rho$, $n = 2$, yielding $\omega^2 \approx 6.1 \times 10^{-3}$, $\lambda \approx 0.083$

	Full program	Approx
$x = .1285$		
ϕ	-5.363, .327i	-5.379, .549i
ϕ_x	-8.634, .361i	-8.66, .379i
$x = .5643$		
ϕ	-9.754, .809i	-9.804, 1.000i
ϕ_x	-12.234, .692i	-12.272, .7608i
C_{ℓ_α}	+31.658, -3.1032i	C_{ℓ_α} 31.7019, -3.2165i

VI. RESULTS

The collocation method was used to solve the partial differential equation resulting from the Gorelov approximation of transonic potential flow in an unstaggered cascade. The system was solved using both the spanning functions proposed by Gorelov in [4], resulting in the equations (V-16); and the Legendre polynomials, resulting in equations (V-15) with f_j replaced by the Legendre polynomial, P_j . The resulting values of $C_{l\alpha}$ for $k = .1$, $\tau = 1$, $\sigma = 1$ and seven collocation points on each blade are presented in figures VI-2, VI-3, and VI-4.

Figure VI -1 presents a diagram which is useful in commenting on the other results. This shows the location of the collocation points and first three interference reflections as a function of w expressed as a percentage of that portion of the chord subject to reflection. The collocation points are equally spaced throughout this interval, 12.5% from the leading edge of the interference zone, 12.5% between each pair of points and 12.5% from the blade trailing edge. The independent variable, w , is plotted vertically so that the dependent variable, percent of chord subject to interference, may be more conveniently visualized along the blade. (The curves are not precisely linear, but are very nearly so in the range shown.)

Figure VI -2 shows the $C_{l\alpha}$ calculated with $k = 0.0$ in comparison with the results obtained from Ackeret theory.

Agreement is good where there is no reflection and the portion of the blade subject to interference is affected by a constant interference potential, $w \geq 0.11$. Throughout the rest of the curve the results calculated here oscillate above and below the theoretical values. This appears to be due to the discrete nature of the approximation used in the collocation method. Rarely is the fraction of the chord subject to interference reflection equal to the fraction of the collocation points which feel it. Where the collocation point fraction lags, as near $w = 0.6$, the collocation results are lower than those due to Ackeret theory. When the collocation point fraction leads, as it does for $w \leq .04$ and briefly for $w \approx .08$, the the collocation results are higher than those due to Ackeret theory. The fault appears to be an intrinsic feature of the small number of points sampled. This results in a set of coefficients similar to those which would be obtained from a generalized Fourier series based on the integration of the Taylor series expansion about each point. This obviously cannot be a good approximation when both the function and its derivative are discontinuous at the reflections.

Figures VI -3 and VI -4 show the results of using Legendre polynomials and Gorelov's functions as spanning functions. The results for both formulations are identical. Gorelov's results are presented for comparison. Agreement is good for $w > 0.05$ except for an anomalous point, marked A .

It is believed that this anomaly is due to the location of the first reflection just ahead of the last collocation point (cf. "A" on Figure VI -1). This will yield a very small contribution from the reflection potential to the linear system from which the collocation points are determined. The resulting system will have a large dynamic range and may be ill-conditioned.

The discrepancy between these results, and those in [3] for $w < 0.5$ is still unexplained, as is the outlying value for $w = 0.5$.

The discontinuities in the imaginary results are believed to be due primarily to the reflection/collocation interaction explained above.

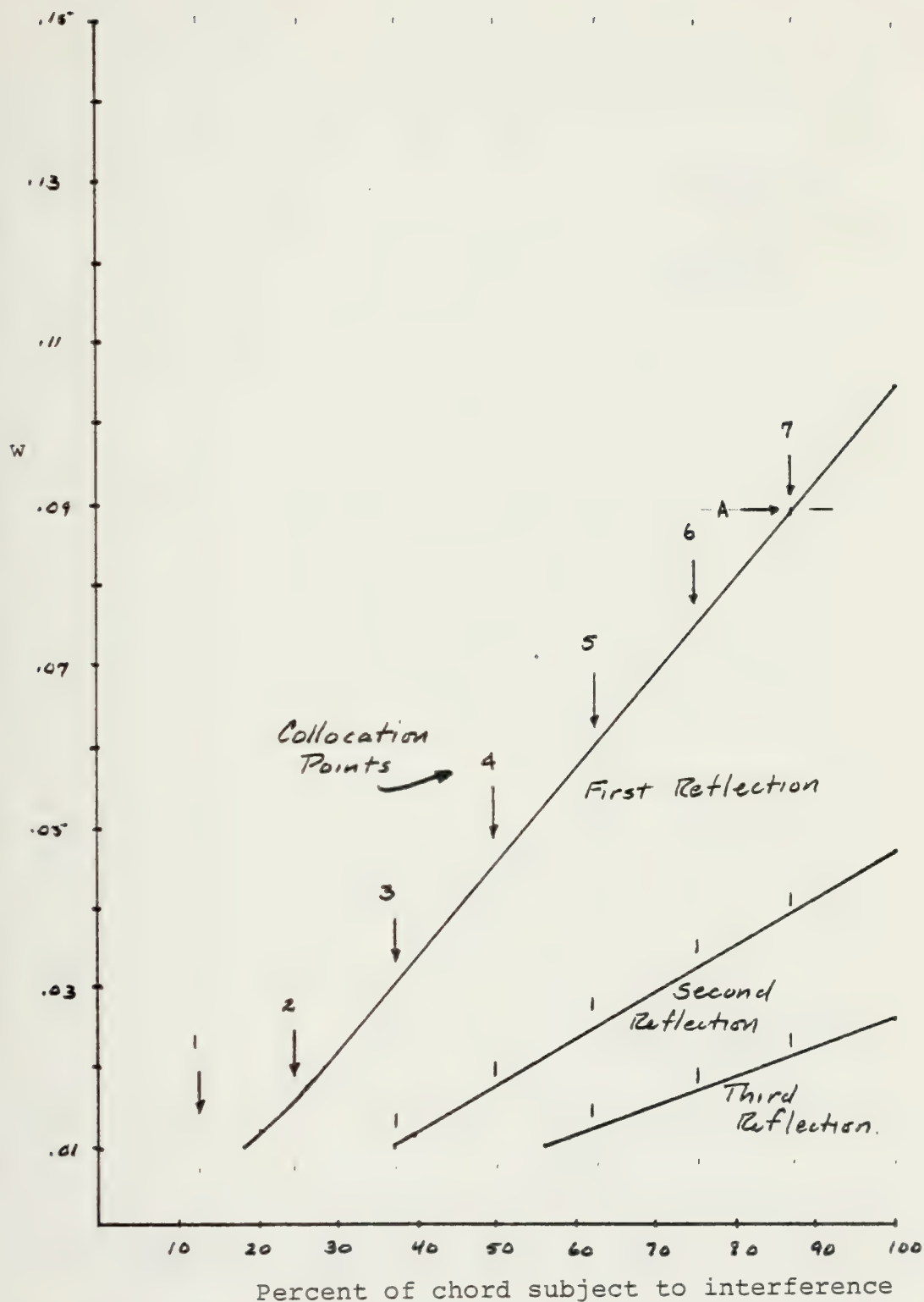


FIGURE VI-1. Location of Reflections and Collocation Points Shown as Percent of Chord Subject to Interference

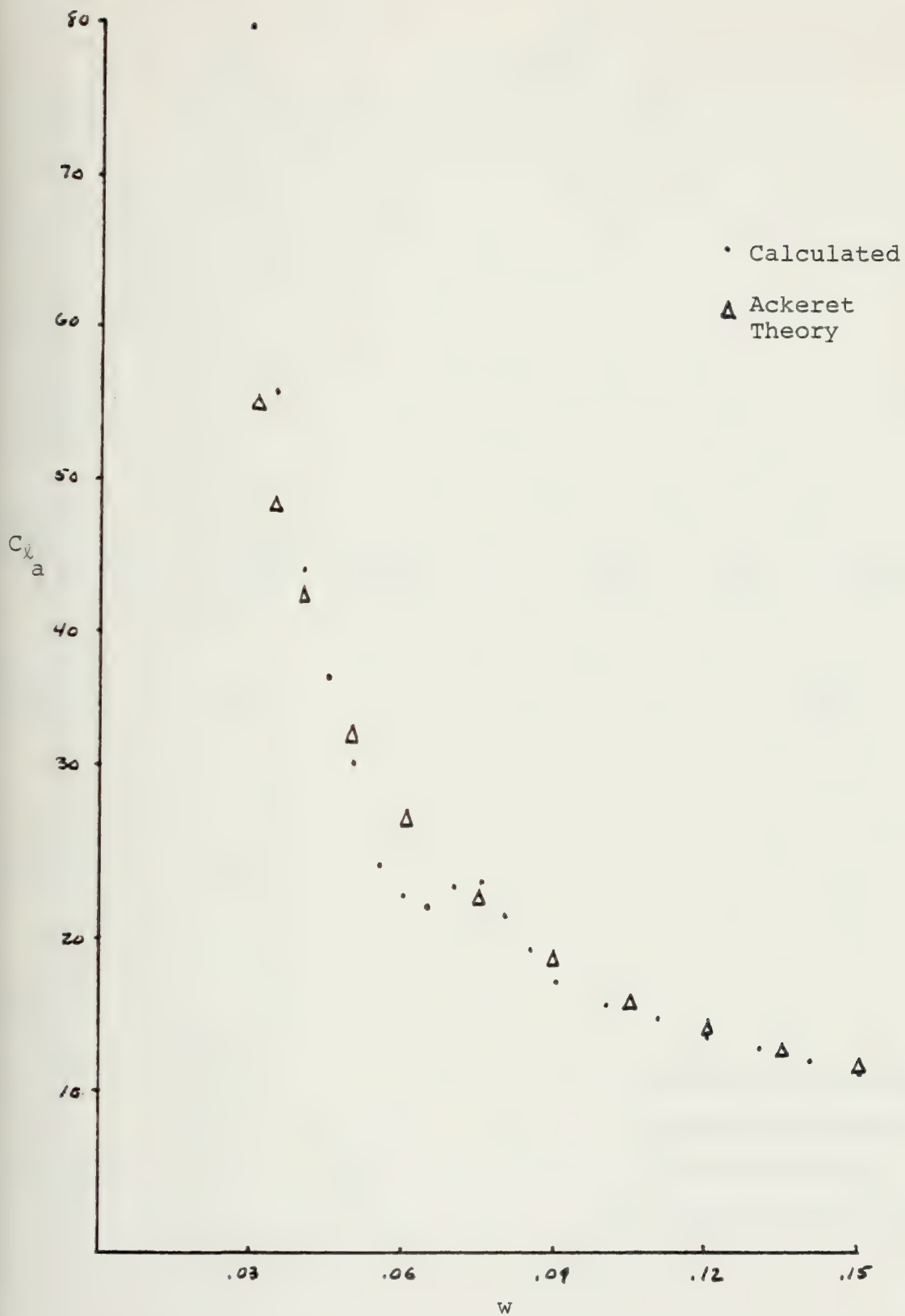


FIGURE VI-2. Comparison of Cl_α -vs- w to that Obtained from Ackeret Theory

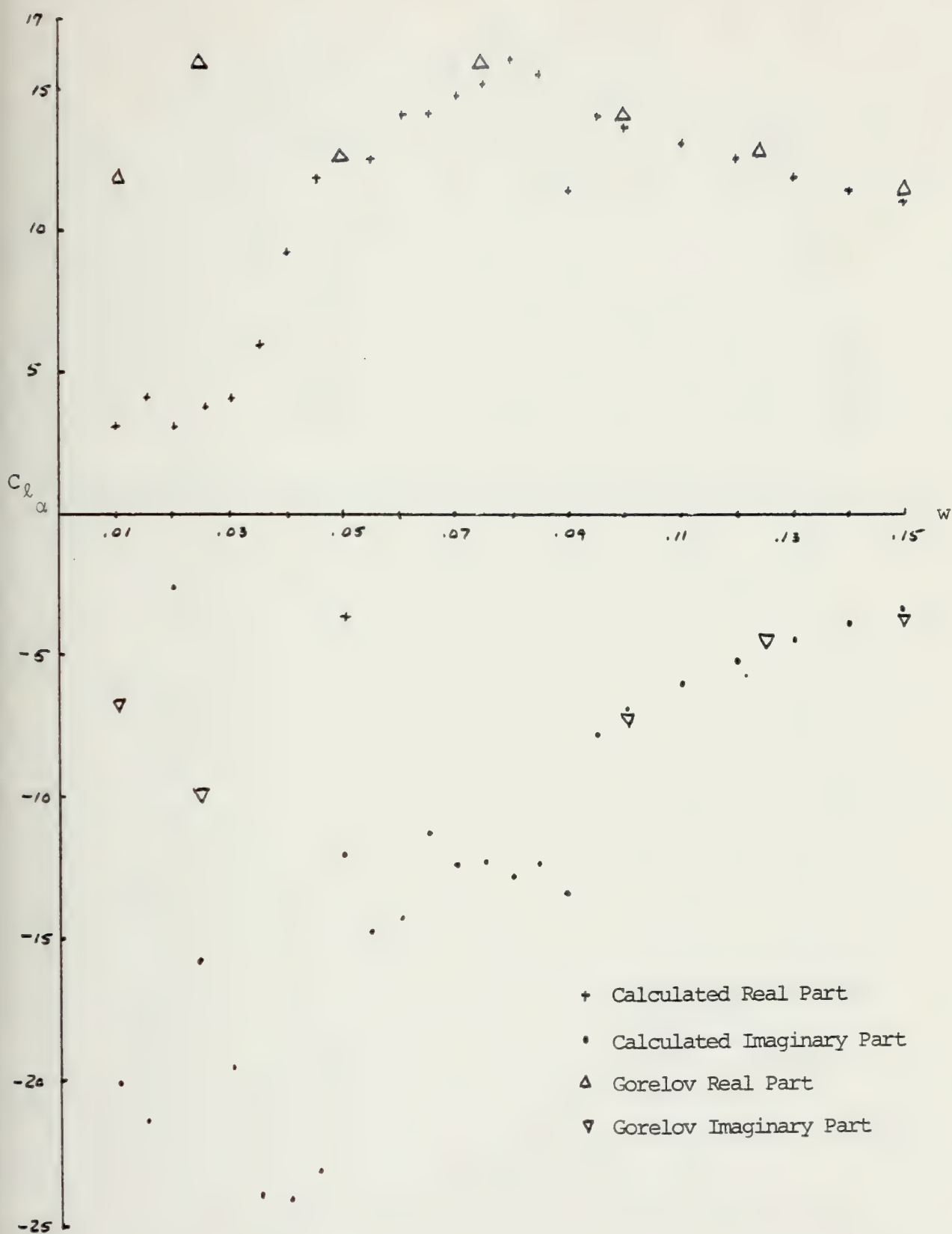


FIGURE VI-3. Plot of Cl_α -vs- w , Legendre Polynomials
 $k = 0.1$, $\tau = 1.0$, $\sigma = \pi$, $n = 7$
 compared with Gorelov's results

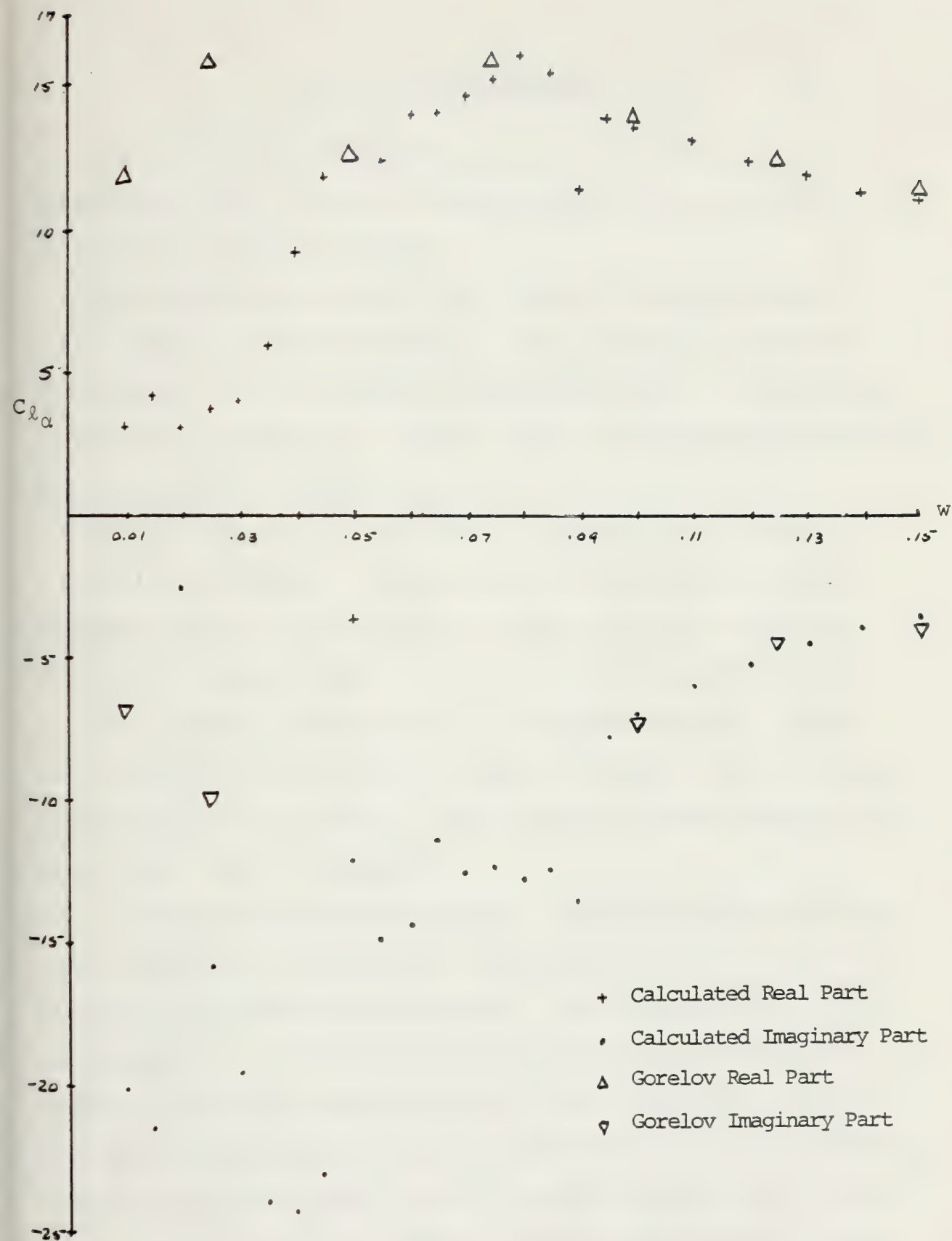


FIGURE VI-4. Plot of $C_{l\alpha}$ -vs- w , Gorelov Spanning Function $k = 0.1$, $\tau = 1.0$, $\sigma = \pi$, $n = 7$ compared with Gorelov's results

VII . RECOMMENDATIONS

There are two recommendations to be made about the techniques used in the collocation method, and a new area in which it might be employed.

The program developed in the course of writing this thesis employs adaptive Simpson's integration to calculate the elements of a completely determined system. This system is solved to provide the coefficients of the spanning functions. Two improvements may be made:

1. The Simpson's integration scheme may be replaced by a Gaussian integrator. Experience has shown that several thousand function evaluations are required by the Simpson's integration routine when C_{α} is to be evaluated for small w . This entails large amounts of computer time and leads to increased accumulations of numerical error. Use of Gaussian integration would probably improve both of these characteristics with little loss of accuracy.

2. The present program treats a completely determined system of dimension $2n+1$ by $2n+1$, and then solves that system to find the collocation coefficients. This procedure has worked satisfactorily in this thesis, but may not work as well at higher frequencies where the final linear system of equations may be ill-conditioned. As an alternative, it is recommended that the boundary conditions be applied at more than n points, say m points, where m is twice or three times as many points, and that the least squares technique be used to determine the

the collocation coefficients which give the minimum square error over-all. This may be thought of as "sampling more data" in order to get more information about the unknown function. The present program could be easily modified in this regard by replacing the spanning function matrix, Q1ZINT, by a new matrix of the form

$$Q1ZINT' = X^T X$$

where X is the new m by n+1 ($m > n+1$) matrix, and replacing the present right-hand-side vector, Q1COF with

$$Q1COF' = X^T Y$$

where Y is the new m by 1 right-hand-side vector. An alternative would be to employ a prepackaged statistical linear regression routine after either modifying the routine to accept complex data, or transforming the present system into a larger system of real numbers only.

The new area in which the collocation method might be employed is the calculation of the potential flow about a staggered cascade. The method could be employed to calculate both the potential in the channel and above the upper blade. The program presented has been designed to enable the

calculation of flow within the channel of a staggered cascade. Unfortunately, there was not enough time to extend the study to this case.

APPENDIX A

PROGRAM DESCRIPTION

This section describes the computer program used to calculate the interference solution to the Gorelov linearization for unsteady transonic flow in a channel. The program written in IBM Fortran IV with the basic structure outlined by Stevens [5]. The basic points are:

1. Organization of the program into small subroutines, each of which performs a specific task.
2. Transmission data to and from subroutines via a formal parameter argument list. No common statements are used.

The end objective is code which is both easy to modify and maintain.

Each subroutine is designed with optional diagnostic printing of its input and output. This is controlled by the parameter IPT. The diagnostic output is printed (only) if $IPT > 0$. Each routine accepts IPT, sets $IOT = IPT - 1$, and then passes IOT as the print parameter to routines it calls. By this method, diagnostic output can be "cascaded" to any desired level. Large initial values of IPT should be avoided because of the spectacular amount of output which can be generated by the double integrals within Q1DCOF.

1. Main Program; including subroutines READ and ABSA.

The basic structure is given in Figure A-1. MAIN calls READ to read input data and then ABSA to calculate the

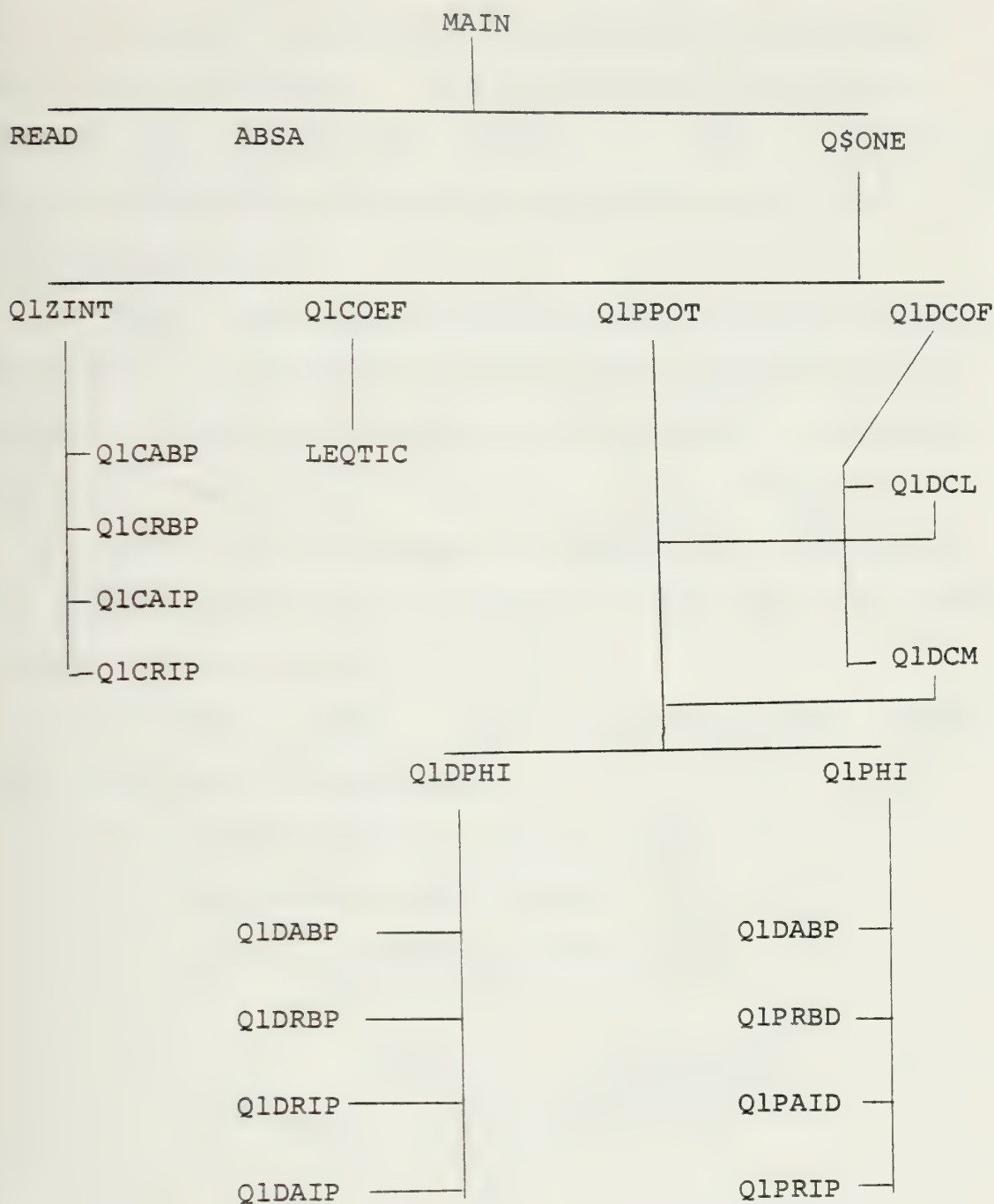


Figure A-1. Program Hierarchy

collocation points. The version of ABSA shown evenly spaces the collocation points across that portion of the blade subject to interference. ABSA may be easily replaced if different point spacing is desired, or if additional points are to be added for an overdetermined system and least squares approximation.

2. Q\$ONE This subroutine controls the actual potential calculation. It performs no calculation itself, but calls subordinate subroutines where the calculations are actually performed. The calling hierarchy is shown in Figure A-1.

3. Q1ZINT This subroutine calculates the linear equation system arising from the boundary conditions. Hierarchy is shown in Figure A-2.

The matrix output is carried through Q1INT. Q1ZINT calls the following subprograms

- a. Q1CRBP returns the value of ϕ_z^0
- b. Q1CABP returns the value of $\phi_{z_1}^1$
- c. Q1CRIP returns the values of ψ_z^0

$$\frac{\partial}{\partial z} \int_{-1+x_*}^{x-z} f_j(s) J_0[\omega \sqrt{(x-s)^2 - z^2}] ds$$

where f_i is one of the set of i elementary functions, $j=1,n$

- d. Q1CAIP returns the values of $\psi_{z_1}^1$

Q1ZINT

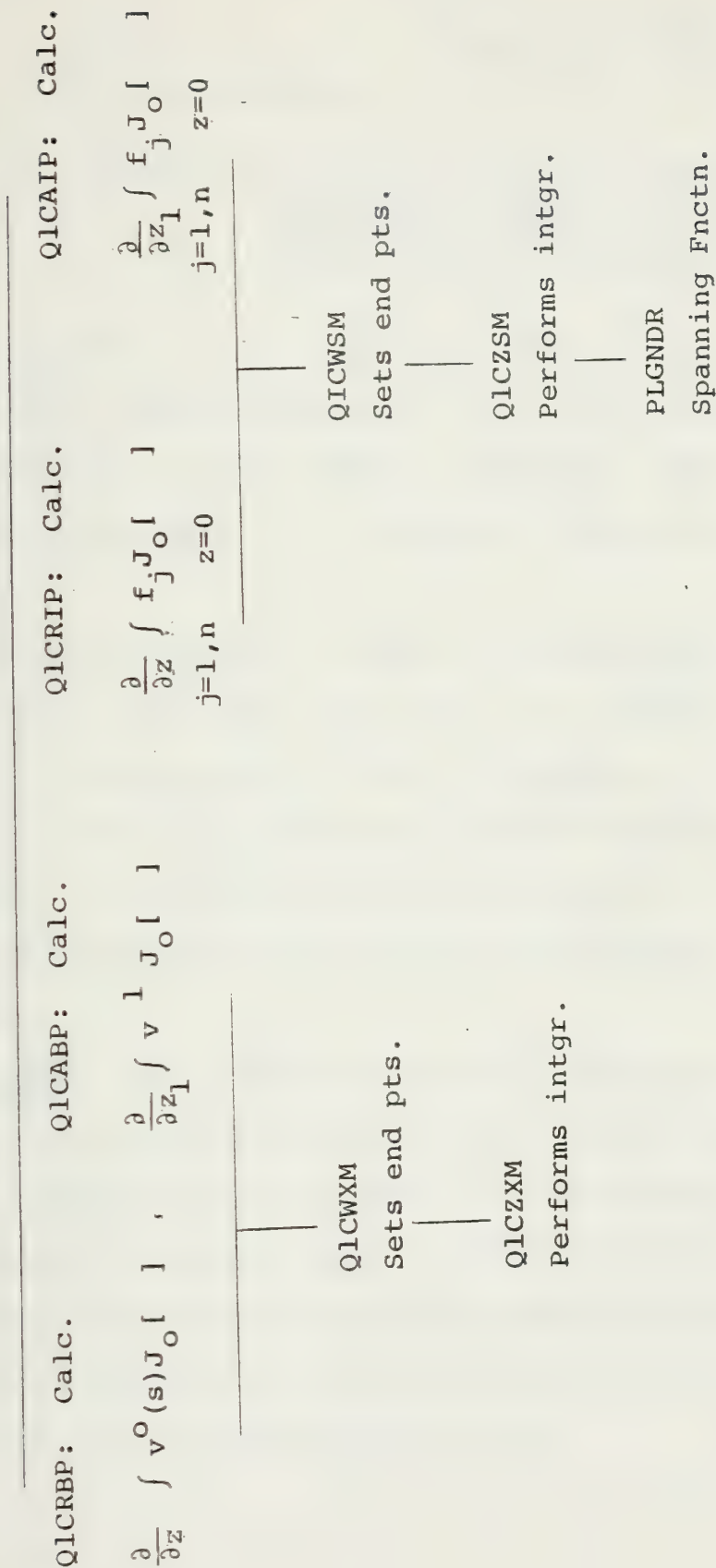


Figure A-2

$$\frac{\partial}{\partial z_1} \int_{-1+\text{OFFSET}+x_*}^{x+z_1} f_j(s) J_0[\omega \sqrt{(x-s)^2 - y_1^2}] ds$$

where f_j is one of the set of elementary functions. OFFSET is a parameter included to facilitate program conversion to a staggered cascade

e. PLGNDR is the subprogram which returns $f_j(x)$, the elementary spanning function. No other routine contains explicit reference to the spanning function. This facilitates easy replacement of the spanning functions should this be desired.

Q1CRBP and Q1CABP in turn call Q1CWXM and Q1CZXM. Q1CWXM computes end-points and then calls Q1CZXM, a complex integration routine based on SIMP by Shampine and Allen [6]. Q1CAIP and Q1CRIP call Q1CWSM and Q1CZSM to perform the integration. Q1ZINT passes the constant matrix to Q1COEF in the array Q1INT and the right-hand-side vector in the array Q1COF.

4. Q1COEF This subroutine employs the IMSL routine LEQT1C to solve the linear system received from Q1ZINT. The resulting coefficients are Q1ABCF for the adjacent blade and Q1RBCF for the reference blade. LEQT2C, the high precision complex IMSL routine may be directly substituted for LEQT1C. Q1COEF may be rewritten to employ the generalized inverse required for least squares approximation

$$A = (x^T x)^{-1} x^T y \quad \text{where } x = Q1INT$$

$$y = Q1COF$$

after first performing the multiplication necessary to replace Q1INT and Q1COF with the proper matrix products in the call to LEQT1C.

5. Q1PPOT This subroutine calculates the potential, ϕ , and ϕ_x , at each collocation point along the reference blade but only if Q1PPOT receives a value of $IPT > 0$, requiring $IPT \geq 2$ on input to the main program. If $IPT \leq 0$, then the subroutine is exited before any calculations are performed. This subroutine is most useful for debugging Q1ZINT and Q1COEF. Q1PPOT calls Q1PABP, Q1PRBP, Q1PRIP, Q1PAIP, Q1DABP, Q1DRBP, Q1DRIP, and Q1DAIP, all of which will be described in the next section.

6. Q1DCOF This subroutine calculates the dimensionless coefficients of lift and moment; C_{l_α} , C_{m_α} . Its internal hierarchy is shown in Figure A-3.

- a. Q1DCL calculates the nondimensional complex coefficient C_{l_α}
- b. Q1DCM calculates the nondimensional complex coefficient C_{m_α} .
- c. Q1PRBP and Q1PABP calculate the potentials due to the reference and adjacent blades respectively. Q1PWXM and Q1PZXM are called to perform the actual integration.

- d. Q1PRIP and Q1PAIP calculate the interference potentials along the reference and adjacent blades. Q1PZSM is called to perform the integration.
- e. Q1DRBP, Q1DABP, Q1DRIP, and Q1DAIP correspond exactly to subroutines above except that the value returned is the partial derivative of the potential with respect to X. Q1DWXM, Q1DZSM, and Q1DXSM perform the co-reponding integrals.

6. Program Listing The program listing shown below incorporates the Legendre functions as spanning functions. Listings for a subroutine employing Gorelov's spanning function and the linear approximation program follow.


```

C      IMPLICIT REAL*8(A-T,C,P,R-Y), COMPLEX*16(Q,Z)
C      DIMENSION X(13), Q2PT(13), Q2CP(13)
C      BLCK ONE READ AND EDIT DATA
C      CALL ERRSET(208,256,-1,1)
C      WRITE(6,91C)
C      CALL READ(CK,DR,CW,RHO,OFFSET,SIGMA,N,NF,IPT)
C      IF(IPT.GT.0)WRITE(6,950)DK,DR,DW,SIGMA,OFFSET,RHG,N,NF
C      FCRMAT(1,1,5X,1,CCRCLCV CASCADCE PROGRAM)
C      FCRMAT(1,0,1,5X,1,INPUT VALUES TRANSMITTED TO MAIN PROGRAM:1,/,
C      1,1,10X,DK,1,13X,DR,1,13X,DW,1,10X,SIGMA,10X,CFFSET,/,
C      1,1,10X,5(E12.5,3X),/,1,5X,NFCN,
C      2,0,1,10X,RHC,1,12X,N,1,5X,12,4X,12)
C      3,1,1,10X,1(E12.2,3X),1,12,4X,12)
C      CALL ABSA(N,CFFSET,X,RHO,DW)
C
C      BLCK TWO CALLS CALCULATION ROUTINES FOR ZCNE ONE AND ZCNE TWO
C      CALL Q$CNE(CK,DR,CW,RHO,OFFSET,SIGMA,N,NF,IPT,X)
C      GO TO 1
C      END
C      SUBROUTINE READ (CK,CR,DW,RHO,OFFSET,SIGMA,N,NF,IPT)
C      IMPLICIT REAL*8(A-T,C,P,R-Y)
C      READ(5,905) DK,DR,CW,RHP,OFFSET,SIGMA,N,NF,IPT
C      IF (DK.LT.0.000) STOP
C      IF (IPT.GT.0) WRITE (6,910)
C      GAMMA = 1.4EO
C      CR = RHP * [SQRT((GAMMA + 1.000) * DW)
C      RFC = DR
C      IF(N.LE.13) GO TO 1
C      N = 13
C      WRITE(6,935)
C      1 IF(OFFSET.LE.1.900) GO TO 2
C      1 WRITE(6,940)
C      OFFSET = CFFSET-1.000
C      GO TO 1
C      2 CCNTINUE
C      IF (IPT.GT.0) WRITE (6,925) DK,DR,DW,SIGMA,CFFSET,RFC,N,NF,IPT
C      RETURN
C      FCRMAT(6F10.4,3I2)
C      905 FCRMAT(1,1,10X,1,STAGGERED SUPERSYNIC CASCADCE PROGRAM)
C      910 FCRMAT(1,0,1,10X,DK,1,13X,DR,1,13X,DW,1,10X,SIGMA,10X,OFFSET,
C      925 1,/,1,10X,5(E12.5,3X),/,1,5X,NFCN,2X,IPT,
C      2,0,1,10X,RHC,1,12X,N,1,5X,12,4X,12)
C      3,1,1,10X,1(E12.2,3X),1,12,4X,12,4X,12)
C      935 FCRMAT(1,0,1,5X,1,ORIGINAL N TOO LARGE (.GT.1.9) - RESET TO N = 13)
C      940 FCRMAT(1,0,1,5X,1,OFFSET TOO LARGE (.GT.1.9), RESET AS CFFSET =
C      1,OFFSET - 1.000)
C      END

```


CLDR0245
 CLDR0255
 CLDR0265
 CLDR0275
 CLDR0285
 CLDR0295
 CLDR0305
 CLDR0315
 CLDR0325
 CLDR0335
 CLDR0345
 CLDR0355
 CLDR0365
 CLDR0375
 CLDR0385
 CLDR0395
 CLDR0405
 CLDR0415
 CLDR0425
 CLDR0435
 CLDR0445
 CLDR0455
 CLDR0465
 CLDR0475
 CLDR0485

```

SUBROUTINE ABSA (N,OFFSET,X,DR,CW
)
IMPLICIT REAL*8 (A-H,O,P,R-Y), COMPLEX*16(Q,Z)
DIMENSION X(13)
XINT = (2.000-DR)/DFLOAT(N+1)
XL = DR - 1.000 + 1.00-8
CC 10 I = 1,N
X(I) = XL + (XINT*DFLOAT(I))
IF(X(I).EQ.0.000) X(I) = .1C-14
10 CONTINUE
RETURN
END
SUBROUTINE C$ONE (CK,CR,CW,RHO,OFFSET,SIGMA,N,NF,IPT,X)
IMPLICIT REAL*8(A-H,O,P,R-Y), COMPLEX*16(Z,C)
DIMENSION Q1COF(26), CLINT(26,26), C1RBP(13), Q1ABP(13)
DIMENSION Q1RBCF(13), Q1ABCF(13)
DIMENSION X(13)
IF(IPT.GT.0)WRITE(6,990)DK,CR,RHO,OFFSET,SIGMA,N,NF,IPT
1 ICT = IPT - 1
2 ICFMAT(,0) SUBROUTINE Q$ONE ENTERED WITH:,,',
3 N,5X,DK,13X,CR,13X,RHC,12X,OFFSET,9X,SIGMA,10X,
NF,2X,IPT,2X,/,12,2X,13)
5X,5(E12,5,3X),12,2X,13)
CALL Q1ZINT(DK,DR,DW,RHO,OFFSET,SIGMA,N,X,ICT,C1CCF,CLINT)
CALL Q1ZCOEF(Q1COF,Q1INT,N,IOT,Q1AECF,Q1RBCF)
CALL Q1PPCT(DK,DR,CW,RHC,OFFSET,SIGMA,N,X,C1AECF,C1RBCF,ICT)
1 ICT = 0
CALL Q1CCCF(DK,DR,CW,RHO,OFFSET,SIGMA,N,Q1AECF,Q1RBCF,ICT)
RETURN
END
C
C
C
SUBROUTINE Q1ZINT(DK,DR,DW,RHO,OFFSET,SIGMA,N,X,
IPT,Q1COF,Q1INT)
IMPLICIT REAL*8(A-H,O,P,R-Y), COMPLEX*16(Z,Q)
DIMENSION Q1COF(26), Q1INT(26,26), Q1INTRP(13), Q1INTAP(13)
DIMENSION X(13)
1 ICFMAT(,0,10X,Q1ZINT ENTERED WITH:,,',/
2 N,10X,DK,13X,DR,13X,RHO,12X,OFFSET,9X,SIGMA,10X,
NF,2X,IPT,2X,/,12,2X,13)
3 N,1CX,5(E12,5,3X),12,2X,13)
CCCNST = CDEXP(DCMPLX(0.000,SIGMA))
IF(IPT.GT.0) WRITE(6,990) DK,DR,RHO,OFFSET,SIGMA,N,IPT
1 ICT = IPT - 1
2 GAMMA = 1.4C0
3 CLAMDA = 1.0C0/(GAMMA+1.000)*CW
CC 90 I = 1,N
IN = I + N
990
  
```


QLDRO485
 QLDRO490
 QLDRO495
 QLDRO500
 QLDRO505
 QLDRO510
 QLDRO515
 QLDRO520
 QLDRO525
 QLDRO530
 QLDRO535
 QLDRO540
 QLDRO545
 QLDRO550
 QLDRO555
 QLDRO560
 QLDRO565
 QLDRO570
 QLDRO575
 QLDRO580
 QLDRO585
 QLDRO590
 QLDRO595
 QLDRO600
 QLDRO605
 QLDRO610
 QLDRO615
 QLDRO620
 QLDRO625
 QLDRO630
 QLDRO635
 QLDRO640
 QLDRO645
 QLDRO650
 QLDRO655
 QLDRO660
 QLDRO665
 QLDRO670
 QLDRO675
 QLDRO680
 QLDRO685
 QLDRO690
 QLDRO695
 QLDRO700
 QLDRO705
 QLDRO710
 QLDRO715
 QLDRO720

```

XSTN = X(I)
XD = XSTN - DR
CEXP = CDEXP(DCMPLX(0.0D0, DLAMCA*XSTN))
CALL QICABP(DK, DR, DW, RHO, OFFSET, SIGMA, XSTN, QINTRP, N, IOT)
CALL QICRIP(DK, DR, DW, RHO, OFFSET, XSTN, QINTRP, N, ICT)
DC 20 J = 1, N
JN = N + J
J1 = J - 1
QIINT(I, J) = PLGNDR(XSTN, DR, J1) * CEXP
QIINT(IN, J) = QINTRP(J)
QIINT(I, JN) = QINTAP(J)
QIINT(IN, JN) = PLGNDR(XSTN - OFFSET, CR, J1) * CEXP
CCNTINUE
CICOF(IN) = -QICREF(DK, DR, DW, RHC, XSTN, IOT)
CICOF(I) = -QICABP(DK, DR, DW, RHO, OFFSET, SIGMA, XSTN, IOT)
CCNTINUE
RETURN
END
COMPLEX FUNCTION QICRBP*16(DK, DR, CW, RHO, XSTN, IPT)
IMPLICIT REAL*8 (A-F, C, P, R-Y), COMPLEX*16 (Q, Z)
DIMENSION QINP(2)
IF (IPT.GT.0) WRITE(6,990) DK, DR, DW, RHO, XSTN, IPT
FORMAT(0, 10X, 'QICREP ENTERED WITH:', /,
, 10X, 'DK, 11X, 'DR, 11X, 'DW, 11X, 'RHO, 10X, 'XSTN, 9X, 'IPT,
, 10X, 5(E12.5, '), I3)
IF(XSTN.LE.RHO-1.0D0) GOTO 20
IF(XSTN.GT.2.0D0) GOTO 20
ICT = IPT - 1
QCK = DCMPLX(0.0D0, DK)
CALL QICWXM(DK, DR, DW, RHO, XSTN, QINP, IOT)
QICRBP = -QCK*QINP(2) - QINP(1)
IF(IPT.LE.0) RETURN
GOTO 30
20 QICRBP = DCMPLX(0.0D0, 0.0D0)
IF(IPT.LE.0) RETURN
30 WRITE(6,995) QICREP
995 FCRMAT(0, 10X, 'QICRBP = ', E14.7, ' ', E14.7)
RETURN
END
COMPLEX FUNCTION QICABP*16(DK, DR, DW, RHO, OFFSET, SIGMA, XSTN, IPT)
IMPLICIT REAL*8 (A-F, O, P, R-Y), COMPLEX*16 (C, Z)
DIMENSION QINP(2)
IF(IPT.GT.0) WRITE(6,990) DK, CR, CW, RHO, OFFSET, SIGMA, XSTN, IPT
FCRMAT(0, 10X, 'QICABP ENTERED WITH:', /,
, 10X, 'DK, 11X, 'CR, 11X, 'DW, 11X, 'RHO, 10X, 'OFFSET, 7X,
, 10X, 5(E12.5, '), I3)
IF(XSTN.LE.RHO-1.0D0) GOTO 20
IF(XSTN.GT.2.0D0) GOTO 20
XSTN = XSTN - OFFSET
IF(XSTN.LE.RHO-1.0D0) GOTO 20

```


QLDR0725
QLDR0730
QLDR0735
QLDR0740
QLDR0745
QLDR0750
QLDR0755
QLDR0760
QLDR0765
QLDR0770
QLDR0775
QLDR0780
QLDR0785
QLDR0790
QLDR0795
QLDR0800
QLDR0805
QLDR0810
QLDR0815
QLDR0820
QLDR0825
QLDR0830
QLDR0835
QLDR0840
QLDR0845
QLDR0850
QLDR0855
QLDR0860
QLDR0865
QLDR0870
QLDR0875
QLDR0880
QLDR0885
QLDR0890
QLDR0895
QLDR0900
QLDR0905
QLDR0910
QLDR0915
QLDR0920
QLDR0925
QLDR0930
QLDR0935
QLDR0940
QLDR0945
QLDR0950
QLDR0955
QLDR0960

```

IF(XASTN.GT.2.000) GOTO 20
ICT = IPT - 1
CCK = DCPLX(0.000,DK)
QCONST = CDEXP(DCPLX(0.000,SIGMA))
CALL QICWXM(CK,DR,DW,RHO,XASTN,CINF,IOT)
QICABP = -(CK*QINP(2) + QINP(1)) * QCONST
IF(IPT.LE.0) RETURN
GC TO 30
CICABP = DCPLX(0.000,0.000)
IF (IPT.LE.0) RETURN
WRITE(6,995) QICABP
FCRMT(0,10X,'QICABP = ',E14.7,',',E14.7)
END
ROUTINE QICWXM(DK,DR,DW,RHO,XSTN,QINP,IPT)
N IS THE MAXIMUM DEGREE OF THE TERM U IN THE INTEGRALS
IMPLICIT REAL * 8 (A-H,O,P,R-Y), COMPLEX * 16 (Z,Q)
DIMENSION QINP(2)
IF(IPT.GT.0) WRITE(6,990) DK,DR,RHC,XSTN,IPT
FORMAT(0,10X,'QICWXM ENTERED WITH ARGUMENTS:',/,
1,10X,DK,16X,ER,16X,RHC,15X,XSTN,14X,IPT,/,
2,10X,4(E13.6,5X),12,5X,I3)
ICT = IPT - 1
B = XSTN - RHO - 1.0D-10
A = -1.000
CALL QICZXM(DK,DR,CW,RHO,XSTN,A,B,1,QANS,IOT)
QINP(1) = QANS
CALL QICZXM(CK,DR,CW,RHO,XSTN,A,B,2,QANS,IOT)
QINP(2) = QANS
IF(IPT.LE.0) RETURN
WRITE(6,995)
FCRMT(0,10X,'J',10X,'QICWXM RESULTS:',/,
1,10X,I,10X,I=1,2
DC 1001 I = 1,2
FORMAT(0,10X,I3,5X,E12.6,',',E12.6)
1001 WRITE(6,996) I, CINF(I)
END
ROUTINE QICZXM(CK,DR,DW,RHC,XSTN,A,B,J,QANS,IPT)
IMPLICIT REAL*8(A-E,G,H,M,O,P,R-Y), COMPLEX*16(F,S,Z)
DIMENSION FV(5),LORR(30),F1F(30),F2T(30),F3T(30),CAT(30)
1 ARESTT(30),QUEST(30),EPST(30),QPSUM(30)
1 F(X) = (X**J)*CDEXP(QEXP*(X))
1 *(OMEGA*RHO/(DSQRT((XSTN-X)*(XSTN-X)-YY)))*
2 MMBSJ1(OMEGA*DSQRT((XSTN-X)*(XSTN-X)-YY),IER))
GAMMA = 1.400

```

CC

QLDR0965
QLDR0970
QLDR0975
QLDR0980
QLDR0985
QLDR0990
QLDR0995
QLDR1000
QLDR1005
QLDR1010
QLDR1015
QLDR1020
QLDR1025
QLDR1030
QLDR1035
QLDR1040
QLDR1045
QLDR1050
QLDR1055
QLDR1060
QLDR1065
QLDR1070
QLDR1075
QLDR1080
QLDR1085
QLDR1090
QLDR1095
QLDR1100
QLDR1105
QLDR1110
QLDR1115
QLDR1120
QLDR1125
QLDR1130
QLDR1135
QLDR1140
QLDR1145
QLDR1150
QLDR1155
QLDR1160
QLDR1165
QLDR1170
QLDR1175
QLDR1180
QLDR1185
QLDR1190
QLDR1195
QLDR1200

```

DM2 = (GAMMA + 1.00) * DW
CMEGA = DSQRT(DK*EK*(1.000-DM2)/(CM2*DM2))
YY = RHO * RHC
CLAMDA = DK/DM2
QEXP = DCMLPX(0.00C, CLAMDA)
J1 = J-1
ACC = 1.00-6
U = 9.00-13
IF (IPT.GT.C) WRITE(6,990)DK,DR,DH,RHO,XSTN,A,B,J,IPT
FCRMAT(0.15X,'QICZM ENTERED WITH ARGUMENTS: ',/,
1.15X,'DK',13X,'DW',10X,'RHC',12X,'XSTN',11X,
2.14X,'B',14X,'J',2X,'IPT',/,
3.15X,'E14.7',12,2X,13)
EF0URU = 4.0*U
IFLAG = 1
EPS = ACC
CERROR = DCMLPX(0.000,0.000)
LVL = 1
LCRR(LVL) = 1
QPSUM(LVL) = 0.0
ALPHA = A
CA = B - A
AREST = 0.0
FV(1) = F(ALPHA)
FV(3) = F(ALPHA + 0.5*DA)
FV(5) = F(ALPHA + DA)
KCOUNT = 3
WT = DA/6.0
CEST = WT*(FV(1) + 4.0*FV(3) + FV(5))
CX = 0.5*DA
FV(2) = F(ALPHA + 0.5*DX)
FV(4) = F(ALPHA + 1.5*DX)
KCOUNT = KOUNT + 2
WT = DX/6.0
QUESTL = WT*(FV(1) + 4.0*FV(2) + FV(3))
QUESTR = WT*(FV(3) + 4.0*FV(4) + FV(5))
QSUM = QUESTL + QUESTR
ARESTL = WT*(CDABS(FV(1)) + CDABS(4.0*FV(2))) + CDABS(FV(3)))
ARESTR = WT*(CDABS(FV(3)) + CDABS(4.0*FV(4))) + CDABS(FV(5)))
AFEA = ARESTL + ARESTR - AREST
QDIFF = CEST - QSUM
IF(CDABS(QDIFF).LE.EPS*DABS(AREA))GO TO 2
IF(DABS(DX).LE.EFCURU*DABS(ALPHA))GO TO 5
IF(LVL.GE.30)GO TO 5
IF(KCOUNT.GE.2000)GO TO 6
LVL = LVL + 1
LCRR(LVL) = 0

```

1


```

11  FIT(LVL) = FV(3)
12  F2T(LVL) = FV(4)
13  F3T(LVL) = FV(5)
14  DA = DX
15  DAT(LVL) = DX
16  AREST = ARESTL
17  ARESTT(LVL) = ARESTR
18  CEST = QESTL
19  CESTT(LVL) = QESTR
20  EPST = EPS/1.4
21  EPST(LVL) = EPS
22  FV(5) = FV(3)
23  FV(3) = FV(2)
24  GC TO 1
25  CERROR = QERROR + QDIFF/15.C
26  IF (LORR(LVL).EQ.0) GO TO 4
27  CSUM = QPSUM(LVL) + CSUM
28  LVL = LVL - 1
29  IF (LVL.GT.1) GO TO 3
30  QANS = CSUM - (B*J1) * CDEXP(QEXF*B)
31  IF (IPT.GT.0) GO TO 11
32  IF (IER.EQ.129) GO TO 11
33  IF (IFLAG.EQ.1) RETURN
34  WRITE(6,990) DK,DR,DW,RHO,XSTN,A,B,J,IPT
35  FCFORMAT(1,15X,'RESULTS: QANS =',E14.7,/,
36  1,15X,'IFLAG',2X,'IER',5X,'QERRCR',/,
37  2,15X,'I3,2X,'I3,5X,'E14.7',/,
38  RETURN
39  QPSUM(LVL) = QSUM
40  LCRR(LVL) = 1
41  ALPHA = ALPHA + DA
42  DA = DAT(LVL)
43  FV(1) = F1T(LVL)
44  FV(3) = F2T(LVL)
45  FV(5) = F3T(LVL)
46  AREST = ARESTT(LVL)
47  QEST = QESTT(LVL)
48  EPS = EPST(LVL)
49  GC TO 1 2
50  IF LAG = 2 3
51  GC TO 2
52  IF LAG = 3
53  GC TO 3
54  END
55  SUBROUTINE C1CAIP(DK,CR,DW,RHO,CFST,SGMA,XSTN,QCAIP,N,IPT)
56  IMPLICIT REAL*8 (A-H,C,P,R-Y), COMPLEX * 16 (C-Z)

```



```

DIMENSION CIMP(13), QCAIP(13)
IF (IPT.GT.0) WRITE(6,990) CK,DR,CW,RHO,OFST,SGMA,XSTN,IPT
950 FCRMAT(0,10X,'Q1CAIP ENTERED WITH: ',/,
,10X,CK,11X,DR,11X,DW,11X,RHC,1CX,'CFST',9X,'SGMA',9XQ
1 XSTN,9X,IPT,/,',10X,7(E12.5,','),I3)
2 XSTN=XSTN-OFST
IF(XSTN.LE.DR+RHO-1.0D0) GOTO 20
IF(XASTN.GT.2.0D0) GOTO 20
ICT=IPT-1
QCCNST=DCPLX(0.0D0,SGMA))
CALL Q1CWSM(DK,DR,DW,RHO,XASTN,N,CINP,IOT)
DC10 I=1,N
QCAIP(I)=CINP(I)
CONTINUE
10 IF(IPT.LE.0) RETURN
GOTO 30
20 ZERC=DCPLX(0.0D0,0.0D0)
DC25 I=1,N
QCAIP(I)=ZERO
CONTINUE
25 IF(IPT.LE.0) RETURN
30 WRITE(6,995)
995 FCRMAT(0,10X,'Q1CAIP RESULTS: J QCAIP(J)')
CC40 J=1,N
WRITE(6,996) J,QCAIP(J)
996 FCRMAT(0,26X,I2,3X,E14.7,',',E14.7)
CONTINUE
40 RETURN
ENCL
SUBROUTINE Q1CRIP (CK,DR,DW,RHC,CFESET,XSTN,CCRIP,N,IPT)
IMPLICIT REAL*8 (A-H,O,P,R-Y), COMPLEX * 16 (C,Z)
DIMENSION CCRIP(13)
IF (IPT.GT.0) WRITE(6,990) CK,DR,CW,RHO,XSTN,IPT
950 FCRMAT(0,10X,'Q1CRIP ENTERED WITH: ',/,
,10X,CK,11X,DR,11X,DW,11X,RHC,1CX,
1 XSTN,9X,IPT,/,',10X,5(E12.5,','),I3)
2 XSTN=XSTN-OFST
IF(XSTN.LE.DR+RHC+CFESET-1.0D0) GOTO 20
IF(XSTN.GT.2.0D0) GOTO 20
ICT=IPT-1
CALL Q1CWSM(DK,DR,DW,RHO,XSTN,N,CCRIP,IOT)
IF(IPT.LE.0) RETURN
GOTO 30
20 ZERC=DCPLX(0.0D0,C.0D0)
CC25 I=1,N
CCRIP(I)=ZERO
990 FCRMAT(0,10X,'Q1CRIP RESULTS: J QCRIP(J)')

```

QLCR1445
 QLCR1450
 QLCR1455
 QLCR1460
 QLCR1465
 QLCR1470
 QLCR1475
 QLCR1480
 QLCR1485
 QLCR1490
 QLCR1495
 QLCR1500
 QLCR1505
 QLCR1510
 QLCR1515
 QLCR1520
 QLCR1525
 QLCR1530
 QLCR1535
 QLCR1540
 QLCR1545
 QLCR1550
 QLCR1555
 QLCR1560
 QLCR1565
 QLCR1570
 QLCR1575
 QLCR1580
 QLCR1585
 QLCR1590
 QLCR1595
 QLCR1600
 QLCR1605
 QLCR1610
 QLCR1615
 QLCR1620
 QLCR1625
 QLCR1630
 QLCR1635
 QLCR1640
 QLCR1645
 QLCR1650
 QLCR1655
 QLCR1660
 QLCR1665
 QLCR1670
 QLCR1675
 QLCR1680


```

DC 40 J = 1,N
WRITE (6,556) J,QCRIP(J)
FCRMAT(, ,26X,12,3X,E14.7, ,',E14.7)
CCCONTINUE
RETURN
ENC
SUBROUTINE QICWSM (CK,DR,DW,RHO, XSTN,N,QINF,IFT)
N IS THE MAXIMUM DEGREE OF THE TERM U IN THE INTEGRALS
IMPLICIT REAL * 8 (A-H,O,P,R-Y), COMPLEX * 16 (Z,Q)
DIMENSION QINP(13)
IF (IPT.GT.0) WRITE (6,990) DK,DR,RHC,XSTN,N,IPT
FCRMAT(, ,10X, ,QICWSM ENTERED WITH ARGUMENTS: ,/,
1, , ,10X, ,DK, ,16X, ,DR, ,16X, ,RHO, ,15X, ,XSTN, ,14X, ,N, ,6X, ,IPT, ,/,
2, , ,10X, ,4(E13.6,5X),12,5X,13)
IOT = IPT - 1
B = XSTN - RHO - 1.0D-10
A = DR - 1.0D0
DC 30 J = 1,N
CALL QICZSM(DK,DR,DW,RHO,XSTN,A,B,J,QANS,IOT)
QINF(J) = QANS
CONTINUE
30 IF (IPT.LE.0) RETURN
WRITE (6,995)
FCRMAT(, ,10X, ,QICWSM RESULTS: ,/,
1, , ,11X, ,J, ,6X, ,QINP(1),
DC 1001 I = 1,N
FORMAT(, ,10X,13,5X,E12.6, ,',E12.6)
1001 WRITE (6,996) I, QINP(I)
RETURN
END
SUBROUTINE QICZSM(DK,DR,DW,RHO,XSTN,A,B,J,QANS,IPT)
IMPLICIT REAL*8 (A - E,G,H,M,O,P,R - Y), COMPLEX*16 (F,Q,Z)
DIMENSION FV(5),LCRR(20),F1T(30),F2T(30),F3T(30),DAT(30),
1 ARESTT(30),QESTT(30),EPST(30),QPSUM(30)
F(X) = PLGNCR(X,DR,J1)*CDEXP(QEXP*(X))
1 CMEGA = RHO / (SQRT((XSTN-X)*(XSTN-X)-YY)) *
2 MBSJ1(OMEGA*DSQRT((XSTN-X)*(XSTN-X)-YY),IER)
GAMMA = 1.4D0
CM2 = (GAMMA + 1.0D0) * DW
CMEGA = CSQRT(DK*DK*(1.0D0-DM2) / (CM2*CM2))
YY = RHC * RHO
LLAMDA = DK/DM2
QEXP = DCMPLX(0.0D0, [LAMDA])
J1 = J-1
ACC = 1.0D-6
U = 9.0D-13

```

QLCR1685
 QLCR1690
 QLCR1695
 QLCR1700
 QLCR1705
 QLCR1710
 QLCR1715
 QLCR1720
 QLCR1725
 QLCR1730
 QLCR1735
 QLCR1740
 QLCR1745
 QLCR1750
 QLCR1755
 QLCR1760
 QLCR1765
 QLCR1770
 QLCR1775
 QLCR1780
 QLCR1785
 QLCR1790
 QLCR1795
 QLCR1800
 QLCR1805
 QLCR1810
 QLCR1815
 QLCR1820
 QLCR1825
 QLCR1830
 QLCR1835
 QLCR1840
 QLCR1845
 QLCR1850
 QLCR1855
 QLCR1860
 QLCR1865
 QLCR1870
 QLCR1875
 QLCR1880
 QLCR1885
 QLCR1890
 QLCR1895
 QLCR1900
 QLCR1905
 QLCR1910
 QLCR1915
 QLCR1920


```

950 IF (IPT.GT.0) WRITE(6,990)DK,DR,CW,RHO,XSTN,A,B,J,IPT
FCMAT(0.,15X,0.01CZSM ENTERED WITH ARGUMENTS:.,/,
1.,15X,DK,13X,CR,13X,CW,10X,RHO,12X,XSTN,11X,
2.,A,14X,B,14X,J,12X,IPT,/,
3.,15X,7(E14.7,.,),12,2X,I3),
EFCURU = 4.C*U
IFLAG = 1
EPS = ACC
CERROR = DCMLPX(0.000,0.000)
LVL = 1
LCRR(LVL) = 1
CFSUM(LVL) = 0.0
ALPHA = A
DA = B - A
AREA = C.0
AREST = 0.0
FV(1) = F(ALPHA)
FV(3) = F(ALPHA + 0.5*DA)
FV(5) = F(ALPHA + DA)
KCUNT = 3
WT = DA/6.0
QEST = WT*(FV(1) + 4.0*FV(3) + FV(5))
DX = 0.5*DA
FV(2) = F(ALPHA + 0.5*DX)
FV(4) = F(ALPHA + 1.5*DX)
KCUNT = KCUNT + 2
WT = DX/6.0
QESTL = WT*(FV(1) + 4.0*FV(2) + FV(3))
QESTR = WT*(FV(3) + 4.0*FV(4) + FV(5))
QSUM = QESTL + QESTR
ARESTL = WT*(CDABS(FV(1)) + CDABS(FV(3)) + CDABS(FV(5)))
ARESTR = WT*(CDABS(FV(2)) + CDABS(FV(4)) + CDABS(FV(5)))
AREA = ARESTL + ARESTR - AREST
QDIFF = QEST - QSUM
IF(CDABS(QDIFF).LE.EPS*CDABS(AREA))GO TO 2
IF(DABS(CX).LE.EFOUR*CDABS(ALPHA))GO TO 5
IF(LVL.GE.3)GO TO 5
IF(KCUNT.GE.2000)GC TC 6
LVL = LVL + 1
LCRR(LVL) = 0
F1T(LVL) = FV(3)
F2T(LVL) = FV(4)
F3T(LVL) = FV(5)
DA = DX
DAT(LVL) = CX
AREST = ARESTL
AFESTT(LVL) = ARESTR
QEST = QESTL

```

QLCR1920
 QLCR1930
 QLCR1935
 QLCR1940
 QLCR1945
 QLCR1950
 QLCR1955
 QLCR1960
 QLCR1965
 QLCR1970
 QLCR1975
 QLCR1980
 QLCR1985
 QLCR1990
 QLCR1995
 QLCR2000
 QLCR2005
 QLCR2010
 QLCR2015
 QLCR2020
 QLCR2025
 QLCR2030
 QLCR2035
 QLCR2040
 QLCR2045
 QLCR2050
 QLCR2055
 QLCR2060
 QLCR2065
 QLCR2070
 QLCR2075
 QLCR2080
 QLCR2085
 QLCR2090
 QLCR2095
 QLCR2100
 QLCR2105
 QLCR2110
 QLCR2115
 QLCR2120
 QLCR2125
 QLCR2130
 QLCR2135
 QLCR2140
 QLCR2145
 QLCR2150
 QLCR2155
 QLCR2160

Q1DR2405
 Q1DR2410
 Q1DR2415
 Q1DR2420
 Q1DR2425
 Q1DR2430
 Q1DR2435
 Q1DR2440
 Q1DR2445
 Q1DR2450
 Q1DR2455
 Q1DR2460
 Q1DR2465
 Q1DR2470
 Q1DR2475
 Q1DR2480
 Q1DR2485
 Q1DR2490
 Q1DR2495
 Q1DR2500
 Q1DR2505
 Q1DR2510
 Q1DR2515
 Q1DR2520
 Q1DR2525
 Q1DR2530
 Q1DR2535
 Q1DR2540
 Q1DR2545
 Q1DR2550
 Q1DR2555
 Q1DR2560
 Q1DR2565
 Q1DR2570
 Q1DR2575
 Q1DR2580
 Q1DR2585
 Q1DR2590
 Q1DR2595
 Q1DR2600
 Q1DR2605
 Q1DR2610
 Q1DR2615
 Q1DR2620
 Q1DR2625
 Q1DR2630
 Q1DR2635
 Q1DR2640

```

103 PLGNDR = (5.000*X2 - 3.000)*X/2.000
    RETURN
104 PLGNDR = ((3.501*X2 - 3.001)*X2 + 3.000)/8.000
    RETURN
105 PLGNDR = ((6.301*X2 - 7.001) *X2 +1.501)*X/8.000
    RETURN
106 PLGNDR=((231.000*X2-315.000)*X2+105.000)*X2-5.000)/16.000
    RETURN
107 PLGNDR=((429.000*X2-693.000)*X2+315.000)*X2-35.000)
    *X/16.000
    RETURN
108 PLGNDR=((6435.000*X2-1202.000)*X2+6930.000)*X2
    -1260.000)*X2+35.000)/128.000
    RETURN
109 PLGNDR=((12155.000*X2-25740.000)*X2+18018.000)
    *X2-4620.000)*X2+315.000)/128.000
    RETURN
110 PLGNDR=((146185.000*X2-109395.000)*X2+5000.000)
    *X2-30030.000)*X2-3465.000)*X2-63.000)/256.000
    RETURN
111 PLGNDR=((188175.000*X2-230945.000)*X2+218750.000)*X2
    -90090.000)*X2+15015.000)*X2-693.000)*X2/256.000
    RETURN
112 PLGNDR=((1676039.000*X2-15395535.000)*X2+2078505)*X2
    -1021020.000)*X2+225225.000)*X2-18018.000)*X2-231.000)/1024.000
    RETURN
    ENCL
    SUBROUTINE C1COEF( Q1COF,Q1INT,N,IPT,Q1ABCF,Q1RBCF)
    IMPLICIT REAL * 8(A-H,O,P,R-Y), COMPLEX * 16(Z,C)
    DIMENSION Q1COF(26), Q1INT(26,26), ZHA(300)
    DIMENSION C1ABCF(13), Q1RBCF(13)
    IF (IPT.GE.C) WRITE (6,90)
    N = 1
    N2 = 2*N
    IF (IPT.LE.0) GO TO 5
    WRITE (6,98) N,N2
    FCFORMAT(10,5X,'Q1COEF ENTERED WITH ',I2,' DEG PWR SERIES (',I2,
    58 1', SQUARE MATRIX)')
    DC 2 I = 1,N2
    WRITE(6,92) I, I, Q1COF(I)
    FCFORMAT(10,10X,'Q1COEF EQUATION SYSTEM, ROW ',I2,/,
    92 1',10X,'Q1COF(',I2,') = ',E14.7,' ',E14.7)
    DC 2 J = 1,N
    J2=J+N
    WRITE(6,91) I,J,Q1INT(I,J),I,J2,Q1INT(I,J2)
    91 2 FCFORMAT(10,15X,2('Q1INT(',I2,',' ,I2,') = ',E14.7,' ',E14.7,10X))
    2 CCNTINUE
  
```



```

QDRIPP = QIDRIP(DK,DR,DR,DW,RHO,OFFSET,XSTN,CIRBCF,N,IOT)
QCAIPO = QIDAIP(DK,DR,DR,DW,RZERO,OFFSET,SIGMA,XSTN,QIABCF,N,IOT)
QDAIPP = QIDAIP(DK,DR,DR,DW,RHO,OFFSET,SIGMA,XSTN,QIABCF,N,IOT)
QCR = QDRO + QDAF + QDRIP + QCAIPP
QDA = QDRP + QDAO + QCRIPP + QDAIFC
QCPHI(I) = QIDCFI(DK,DR,DR,DW,RHO,OFFSET,SIGMA,N,CIRBCF,
1 XSTN,IOT)
WRITE(6,55) I,XSTN
WRITE(6,91) QCR,QCRO,QDAP,QDRIP,QCAIPP
WRITE(6,92) CCA,QCRP,QDAO,QCRIPP,CCAIPQ
QCR(I) = QCR
QCAAI(I) = CCA
95 FCRMAT(1,0,XSTN,CFHI(I),QRR(I),QAA(I)
1, BL TCTAL D(POT)/CX',14X,'REF EL D(POT)/CX',8X,
2, 'ADJ BL D(POT)/DX',8X,'REF EL INT D(POT)/DX',5X,
3, 'ADJ BL INT D(POT)/DX')
10 CONTINUE
WRITE(6,994)
994 FCRMAT(1,1,10X,'SLMMARY LISTING',/,
1,0,10X,'XSTN',7X,'SINGLE BLADE TOTAL POTENTIAL',6X,
1, 'REF BLADE POTENTIAL',15X,'ADJ BLADE POTENTIAL')
DC 20 I = 1,N
XSTN = X(I)
WRITE(6,94) XSTN,CFHI(I),QRR(I),QAA(I)
94 FCRMAT(1,1,10X,F6.4,3(5X,E14.7,1,1,E14.7))
20 CONTINUE
WRITE(6,996)
996 FFORMAT(1,0,10X,'XSTN',7X,'SINGLE ELADE TOTAL D(PCT)/DX',6X,
1, 'REF ELADE D(POT)/DX',15X,'ADJ ELADE D(POT)/DX')
DC 50 I = 1,N
XSTN = X(I)
WRITE(6,94) XSTN,CCPHI(I),QDRR(I),CDAI(I)
50 CONTINUE
RETURN
END
SUBROUTINE QIDCOF(CK,CR,DR,DW,RHC,CFFSET,SIGMA,N,QIABCF,QIRBCF,IFT)
IMPLICIT REAL*8 (A-H,O,P,R-Y), COMPLEX*16 (Q,Z)
DIMENSION QIABCF(13), QIRBCF(13)
IF(IPT.GT.0)WRITE(6,990) DK,DR,CK,RHO,OFFSET,SIGMA,N,IPT
990 FCFMAT(1,1,10X,'QIDCOF - CALCULATION OF COMPLEX DIMENSIONLESS AERO-
1 DYNAMIC COEFFICIENTS',/,
3,0,10X,'CK',13X,'DR',13X,'DW',12X,'OFFSET',5X,'SIGMA',
4,10X,'N',3X,'IPT',/,
56(E12.5,3X),12,2X,I3)
ICT = IPT - 3
995 QICCL = QICCL(DK,DR,DW,RHO,OFFSET,SIGMA,N,QIABCF,CIRBCF,IOT)
QICCM = QIDCCM(DK,DR,DW,RHC,CFFSET,SIGMA,N,QIABCF,QIRBCF,IOT)
GAMMA = 1.4CO

```


QLCR4565
QLCR4570
QLDR4575
QLDR4580
QLDR4585
QLDR4590
QLDR4595
QLDR4600
QLDR4605
QLDR4610
QLDR4615
QLDR4620
QLDR4625
QLDR4630
QLDR4635
QLDR4640
QLDR4645
QLDR4650
QLDR4655
QLDR4660
QLDR4665
QLDR4670
QLDR4675
QLDR4680
QLDR4685
QLDR4690
QLDR4695
QLDR4700
QLDR4705
QLDR4710
QLDR4715
QLDR4720
QLDR4725
QLDR4730
QLDR4735
QLDR4740
QLDR4745
QLDR4750
QLDR4755
QLDR4760
QLDR4765
QLDR4770
QLDR4775
QLDR4780
QLDR4785
QLDR4790
QLDR4795
QLDR4800

```

IF(IPT.LE.0) RETURN
GOTO 60
C1PRBP = CCMPLX(C.C00,0.000)
IF(IPT.LE.0) RETURN
WRITE(6,995) Q1PRBP
995 FORMAT('0',10X,'Q1PRBP = ',E14.7,' ',E14.7)
RETURN
END
COMPLEX FUNCTION C1PABP*16(DK,DR,CW,RHO,OFFSET,SIGMA,XSTN,IPT)
IMPLICIT REAL*8 (A-F,C,P,R-Y), COMPLEX*16 (Q,Z)
DIMENSION QINP(2)
IF(IPT.GT.0) WRITE(6,990) DK,DR,DW,RHO,OFFSET,SIGMA,XSTN,IPT
990 FCMAT(C,C,10X,'Q1PABP ENTERED WITH:',/,/,
1,10X,'C',1X,'DR',11X,'DW',11X,'RHO',10X,'CFFSET',7),
2,SIGMA,8),XSTN,9X,'IPT',/,/,10X,7(E12.5,' ',1),1)
XASTN=XSTN-OFFSET
IF(XASTN.LE.RHO-1.000) GOTO 20
ICT=IPT-1
CCK = CCMPLX(0.000,DK)
QCONST = CDEXP(CCMPLX(0.000,SIGMA))
CALL Q1PWXM(DK,DR,CW,RHC,XASTN,CINP,IOT)
Q1PABP = (QCK*QINF(2) + QINP(1)) * QCONST
IF(IPT.LE.0) RETURN
GOTO 60
Q1PABP = CCMPLX(C.C00,0.000)
IF(IPT.LE.0) RETURN
990 WRITE(6,995) Q1PABP
995 FCMAT('0',10X,'Q1PABP = ',E14.7,' ',E14.7)
RETURN
END
SUBROUTINE C1PWXM(CK,DR,DW,RHO,XSTN,CINP,IPT)
N IS THE MAXIMUM DEGREE OF THE TERM U IN THE INTEGRALS
IMPLICIT REAL*8 (A-F,C,P,R-Y), COMPLEX*16 (Z,Q)
DIMENSION QINP(2)
IF(IPT.GT.0) WRITE(6,990) DK,DR,RHC,XSTN,IPT
990 FCMAT(C,C,10X,'Q1PWXM ENTERED WITH ARGUMENTS:',/,IPT,/,
1,10X,'DK',16X,'DR',16X,'RHO',15X,'XSTN',14X,'IPT',/,
2,10X,4(E13.6,5X),12,5X,13)
ICT=IPT-1
B = XSTN - RHO - 1.00-10
A = -1.000
CALL C1PZXM(CK,DR,CW,RHO,XSTN,A,E,1,QANS,IOT)
QINP(1) = QANS
CALL C1PZXM(CK,DR,CW,RHC,XSTN,A,E,2,QANS,IOT)
QINP(2) = QANS
IF(IPT.LE.0) RETURN

```

C
C
C


```

995 WRITE (6,995)
   FFORMAT(, , , 1CX, , QIPWXM RESULTS: , / ,
1 , , 11X, , J, , 6X, , QINP(1), )
   CC 1001 I = 1, 2
996 FFORMAT(, , , 10X, 13, 5X, E12.6, , , E12.6)
1001 WRITE (6,996) I, QINP(I)
   RETURN
END
SUBROUTINE QIPZXM (DK, DR, DW, RHO, XSTN, A, B, J, CANS, IPT)
IMPLICIT REAL*8 (A - E, G, I, M, O, P, R - Y), COMPLEX*16 (F, Q, Z)
DIMENSION FV(5), LCFR(60), F1T(60), F2T(60), F3T(60), DAT(60),
1 ARESTT(60), QESTT(60), EPST(60), CFSUM(60)
1 F(X) = (X**JI)*CDEXP(QEXP*(X))
1 * MMESJO((OMEGA * DSQRT( (XSTN-X)*(XSTN-X) - YY)), IER)
   GAMMA = 1.4D0
   YY = RHO * RPO
   DM2 = (GAMMA + 1.0D0) * DW
   CMEGA = CSQRT (DK*CK*(1.0D0-DM2)/(CM2*DM2))
   CLAMDA = CK/DM2
   CM = DSQRT(CM2)
   C2EXP = CDEXP (DCMPLX (0.0D0, -DLAMDA*XSTN))
   QEXP = DCMPLX(0.0D0, CLAMDA)
   J1 = J - 1
   U = 9.0D-13
   ACC = 1.0D-6
   IF (IPT.GT.0) WRITE(6,990)DK,DR,DW,RHO,XSTN,A,B,J,IPT
990 FFORMAT(, , , 15X, , CIPZXM ENTERED , IF ARGUMENTS: , / ,
1 , , 15X, , DK, , 13X, , CR, , 13X, , DW , , 10X, , RHO, , 12X, , XSTN, , 11X,
2 , A, , 14X, , B, , 14X, , J, , 2X, , IPT, , / ,
3 , , 15X, , 7(E14.7, , ,), 12, 2X, 13,
   EFFCURU = 4.0*U
   IFLAG = 1
   IFS = ACC
   CERROR = CCMPLX(0.0D0, 0.0D0)
   LVL = 1
   LCRRR(LVL) = 1
   CFSLM(LVL) = 0.0
   ALPHA = A
   C1 = B - A
   AREA = 0.0
   AREST = 0.0
   FV(1) = F(ALPHA)
   FV(3) = F(ALPHA + 0.5*DA)
   FV(5) = F(ALPHA + CA)
   KCUNT = 3
   WT = DA/6.0
   CEST = WT*(FV(1) + 4.0*FV(3) + FV(5))
   CX = 0.5*DA
1

```

CLDR4805
 QLDR4810
 QLDR4815
 QLDR4820
 QLDR4825
 QLDR4830
 QLDR4835
 QLDR4840
 QLDR4845
 QLDR4850
 QLDR4855
 QLDR4860
 QLDR4865
 QLDR4870
 QLDR4875
 QLDR4880
 QLDR4885
 QLDR4890
 QLDR4895
 QLDR4900
 QLDR4905
 QLDR4910
 QLDR4915
 QLDR4920
 QLDR4925
 QLDR4930
 QLDR4935
 QLDR4940
 QLDR4945
 QLDR4950
 QLDR4955
 QLDR4960
 QLDR4965
 QLDR4970
 QLDR4975
 QLDR4980
 QLDR4985
 QLDR4990
 QLDR4995
 QLDR5000
 QLDR5005
 QLDR5010
 QLDR5015
 QLDR5020
 QLDR5025
 QLDR5030
 QLDR5035
 QLDR5040


```

FV(2) = F(ALPHA + 0.5*DX)
FV(4) = F(ALPHA + 1.5*DX)
KCUNT = KCUNT + 2
WT = DX/6.0
QUESTL = WT*(FV(1) + 4.0*FV(2) + FV(3))
QUESTR = WT*(FV(3) + 4.0*FV(4) + FV(5))
QSUM = QUESTL + QUESTR
ARESTL = WT*(CDABS(FV(1)) + CDAES(4.0*FV(2)) + CDABS(FV(3)))
ARESTR = WT*(CDABS(FV(3)) + CDABS(4.0*FV(4)) + CDABS(FV(5)))
AREA = AREAL + ((ARESTL + ARESTR) - AREST)
CDIFF = QUEST - QSUM
IF(CDABS(CDIFF).LE.EPS*CDABS(AREA))GO TO 2
IF(CDABS(CX).LE.EFOUR*CDABS(ALPHA))GO TO 5
IF(LVL.GE.60)GO TO 5
IF(KOUNT.CE.2000)GC TC 6
LVL = LVL + 1
LCRR(LVL) = 0
F1T(LVL) = FV(3)
F2T(LVL) = FV(4)
F3T(LVL) = FV(5)
DA = DX
AREST(LVL) = CX
ARESTL = ARESTL
ARESTR(LVL) = ARESTR
QUESTL = QUESTL
QUESTR = QUESTR
EPS = EPS/1.4
EFST(LVL) = EPS
FV(5) = FV(3)
FV(3) = FV(2)
GC TO 1
QERROR = QERROR + CDIFF/15.0
IF(LORR(LVL).EQ.0)GC TC 4
QSUM = QSUM(LVL) + QSUM
LVL = LVL - 1
GO TO 3
QANS = QSUM * Q2EXP / DM
IF(IPT.GT.0) GO TO 11
IF(IFLAG.EC.1) RETURN
WRITE(6,990) DK,DR,DW,RHO,XSTN,A,E,J,IPT
WRITE(6,995) QANS,IFLAG,IER,QERR,F,E14.7,/,E14.7,/,
FCRMTAT(,1,15X,IFLAG,2X,IER,5X,QERR(F,/,
1,15X,I3,2X,I3,5X,E14.7,/,E14.7)
2 RETURN
QFSUM(LVL) = QSUM
LCRR(LVL) = 1
ALPHA = ALPHA + DA

```

2 3

11 995

4


```

1000  DA = CAT(LVL)
1001  FV(1) = F1T(LVL)
1002  FV(3) = F2T(LVL)
1003  FV(5) = F3T(LVL)
1004  AREST = ARESTT(LVL)
1005  AREST = ARESTT(LVL)
1006  EPS = EPST(LVL)
1007  GC TO 1
1008  IF LAG = 2
1009  GC TO 2
1010  IF LAG = 3
1011  GC TO 2
1012  GC TO 2
1013  END
1014  COMPLEX FUNCTION CIPAI*16(CK,DR,DW,RHO,OFST,SGMA,XSTN,CICF,N,IPT)
1015  IMPLICIT REAL*8 (A-H,C,P,R-Y), COMPLEX*16 (Q,Z)
1016  DIMENSION QICF(13)
1017  IF (IPT.GT.0) WRITE(6,990) DK,DF,DW,RHO,OFST,SGMA,XSTN,IPT
1018  FCRMAT(0,10X,'QIPAI ENTERED WITH:',/,
1019  1,10X,'CK,11X,'DR,11X,'DW,11X,'RHO,10X,'OFST',5X,'SGMA',5X,
1020  2,'XSTN,9X,'IPT,/,',10X,7(E12.5,','),13)
1021  XSTN = XSTN - OFST
1022  IF (XSTN.LE.DR+RHO-1.000) GO TO 20
1023  A = CR-1.000
1024  B = XSTN - RHO - 1.00E-8
1025  ICT = IPT - 1
1026  CCONST = CDEXP( CCPLX(0.000,SGMA))
1027  CALL QIPZSM(CK,DR,DW,RHO,XSTN,A,E,N,QANS,QICF,ICT)
1028  CIPAI = QANS
1029  IF (IPT.LE.0) RETURN
1030  GC TO 60
1031  CIPAI = CCPLX(0.000,0.000)
1032  IF (IPT.LE.0) RETURN
1033  WRITE(6,995) QIPAI
1034  FORMAT(0,10X,'CIPAI = ',E14.7,',',E14.7)
1035  RETURN
1036  END
1037  COMPLEX FUNCTION CIPRI*16(CK,DR,DW,RHO,OFFSE1,XSTN,CICF,N,IPT)
1038  IMPLICIT REAL*8 (A-H,C,P,R-Y), COMPLEX*16 (Q,Z)
1039  DIMENSION CICF(13)
1040  IF (IPT.GT.0) WRITE(6,990) DK,DR,DW,RHO,XSTN,IPT
1041  FCRMAT(0,10X,'QIPRI ENTERED WITH:',/,
1042  1,10X,'CK,11X,'DR,11X,'DW,11X,'RHO,10X,'
1043  2,'XSTN,9X,'IPT,/,',10X,5(E12.5,','),13)
1044  IF (XSTN.LE.DR+RHO+OFFSET-1.000) GC TO 20
1045  A = CR-1.000
1046  B = XSTN - RHO - 1.00E-8
1047  ICT = IPT - 1
1048  CALL QIPZSM(DK,DR,DW,RHO,XSTN,A,B,N,QANS,CICF,ICT)

```



```

Q1DR55230
Q1DR55235
Q1DR55305
Q1DR55345
Q1DR55405
Q1DR55455
Q1DR55505
Q1DR55555
Q1DR55605
Q1DR55655
Q1DR55705
Q1DR55755
Q1DR55805
Q1DR55855
Q1DR55905
Q1DR55955
Q1DR56005
Q1DR56055
Q1DR56105
Q1DR56155
Q1DR56205
Q1DR56255
Q1DR56305
Q1DR56355
Q1DR56405
Q1DR56455
Q1DR56505
Q1DR56555
Q1DR56605
Q1DR56655
Q1DR56705
Q1DR56755
Q1DR56805
Q1DR56855
Q1DR56905
Q1DR56955
Q1DR57005
Q1DR57055
Q1DR57105
Q1DR57155
Q1DR57205
Q1DR57255
Q1DR57305
Q1DR57355
Q1DR57405
Q1DR57455
Q1DR57505
Q1DR57555
Q1DR57605

```

```

Q1FRIP = -QANS
IF(IPT.LE.0) RETURN
GOTO 60
Q1PRIP = DCMPLX(0.0D0,0.0D0)
2C IF (IPT.LE.0) RETURN
60 WRITE(6,995) Q1PRIP
995 FCRMAT(0,10X,'Q1FRIP = ',E14.7,' ',E14.7)
RETURN
END
SUBROUTINE Q1PZSM (DK,DR,DW,RHO,XSTN,A,B,J,CANS,Q1CF,IPT)
IMPLICIT REAL*8(A - E,G,H,M,O,P,R - Y), COMPLEX*16(F,Q,Z)
DIMENSION Q1CF(13),LORR(60),FIT(60),F2T(60),F3T(60),CAT(60),
DIMENSION FV(5),EPST(60),QPSUM(60)
1 ARESTT(60),QUEST(60),CDEXF(CEXP*(X))
1 F(X) = CLGNCR(X,CR,J,Q1CF)*CDEXF(CEXP*(X))
1 * MMBSJC((OMEGA * DSQRT( (XSTN-X)*(XSTN-X) - YY)), IER)
GAMMA = 1.4D0
YY = RHC * RHO
DM2 = (GAMMA + 1.0D0) * DW
OMEGA = DSQRT (DK*DK*(1.0D0-DM2)/(DM2*DM2))
DLAMDA = DK/DM2
DN = DSQRT(DM2)
Q2EXP = CDEXP (DCMFLX (0.0D0, -DLAMDA*XSTN))
QEXP = DCMPLX(0.0D0,CLAMDA)
U = 9.0D-13
ACC = 1.0D-6
IF (IPT.GT.0) WRITE(6,990)DK,DR,DW,RHO,XSTN,A,B,J,IPT
FCFMAT(0,15X,'Q1PZSM ENTERED WITH ARGUMENTS:',/,
1,15X,DK,13X,DR,13X,DW,10X,'RHC',12X,'XSTN',11X,
2,A,14X,B,14X,J,12X,IPT,/,
3,15X,7(E14.7,/,12,2X,I3)
EFOURU = 4.0*U
IFLAG = 1
EPS = ACC
CERROR = DCMPLX(0.0DC,0.0D0)
LVL = 1
LCRR(LVL) = 1
QFSLM(LVL) = 0.0
ALPHA = A
DA = B - A
AREA = 0.0
AFEST = 0.0
FV(1) = F(ALPHA)
FV(3) = F(ALPHA) + C.5*DA)
FV(5) = F(ALPHA) + CA)
KCUNT = 3
WT = DA/6.0
QEST = WT*(FV(1) + 4.0*FV(3) + FV(5))

```


Q1DR6005
Q1DR6010
Q1DR6015
Q1DR6020
Q1DR6025
Q1DR6030
Q1DR6035
Q1DR6040
Q1DR6045
Q1DR6050
Q1DR6055
Q1DR6060
Q1DR6065
Q1DR6070
Q1DR6075
Q1DR6080
Q1DR6085
Q1DR6090
Q1DR6095
Q1DR6100
Q1DR6105
Q1DR6110
Q1DR6115
Q1DR6120
Q1DR6125
Q1DR6130
Q1DR6135
Q1DR6140
Q1DR6145
Q1DR6150
Q1DR6155
Q1DR6160
Q1DR6165
Q1DR6170
Q1DR6175
Q1DR6180
Q1DR6185
Q1DR6190
Q1DR6195
Q1DR6200
Q1DR6205
Q1DR6210
Q1DR6215
Q1DR6220
Q1DR6225
Q1DR6230
Q1DR6235
Q1DR6240

```

ALPHA = ALPHA + CA
CA = CAT(LVL)
FV(1) = F1T(LVL)
FV(3) = F2T(LVL)
FV(5) = F3T(LVL)
AREST = ARESTT(LVL)
CEST = CESTT(LVL)
EPST = EPST(LVL)
GC TO 1 2
IFLAG 2 = 3
IFLAG 2 = 2
IFLAG 2 = 2
GC TO 2
END
CCOMPLEX FUNCTION QLGNDR*16(X,DR,N,C1CF)
IMPLICIT REAL*8 (A-F,C,P,R-Y), COMPLEX*16(Q,Z)
DIMENSION Q1CF(13)
ZSUM = 0.0DC
DO 10 J1 = 1,N
  J = J1 - 1
  ZSUM = ZSUM + PLGNDR(X,DR,J) * C1CF(J1)
10 CONTINUE
ZSUM = ZSUM
GLGNDR = ZSUM
RETURN
END
CCOMPLEX FUNCTION Q1DPHI*16 (DK,DR,DW,RHO,CFFSET,SIGMA,
N,Q1ABCF,Q1RBCF,XSTN,IPT)
IMPLICIT REAL*8(A-H,O,P,R-Y), COMPLEX*16(C,Z)
DIMENSION Q1ABCF(13),Q1RBCF(13)
IF (IPT.GT.0) WRITE(6,990) CK,DR,DW,RHO,OFF(E1,SIGMA,XSTN,IPT)
FCMAT(0,10X,DK,11X,DR,11X,DW,11X,RHO,10X,OFFSET,7X,SIGMA,
1 8X,XSTN,9X,IPT,/,10X,7(E12.5,.,,),I3)
2 RZERC = 0.0DC
ICT = IPT - 1
C1DPHI = C1CREP(CK,DR,DW,RZERO,XSTN,IOT)
1 + Q1DIRP(DK,DR,DW,RZERO,OFFSET,XSTN,Q1RBCF,N,IOT)
2 + C1DAEP(CK,DR,DW,RZERO,OFFSET,SIGMA,XSTN,IOT)
3 + C1DAIP(CK,DR,DW,RZERO,OFFSET,SIGMA,XSTN,C1AECF,N,IOT)
IF (IPT.LE.0) RETURN
WRITE(6,995) Q1CPHI
FCMAT(0,10X,10X,Q1DPHI = ,E14.7,.,,E14.7)
995 RETURN
END
CCOMPLEX FUNCTION Q1DREP*16(DK,DR,DW,RHO,XSTN,IPT)
IMPLICIT REAL*8 (A-F,C,P,R-Y), COMPLEX*16(Q,Z)
DIMENSION C1NP(2)
IF (IPT.GT.C) WRITE(6,990) DK,DR,DW,RHO,XSTN,IPT

```



```

55C FCRMAT('0',10X,'QICRBP ENTERED WITH:',/,
1      ' ',10X,'DK',11X,'CR',11X,'DW',11X,'RHO',10X,'XSTN',9X,'IPT',
2      ' ',10X,5(E12.5,' '),13)
      IF(XSTN.LE.RHO-1.0CO) GOTO 20
      ICT = IPT - 1
      QDK = DCMLPX(0.0D0,DK)
      CALL Q1DWXM(DK,DR,DW,RHO,XSTN,QINP,IOT)
      QIDRBP = -QCK*QINP(2) - QINP(1)
      IF(IPT.LE.0) RETURN
      GOTO 60
20  QIDRBP = DCMLPX(0.0D0,0.0CO)
      IF(IPT.LE.0) RETURN
      WRITE(6,995) QIDREP
995  FCRMAT('0',10X,'QIDRBP = ',E14.7,' ',E14.7)
      RETURN
      ENCL
      COMPLEX FUNCTION Q1DABP*16(DK,DR,DW,RHO,CFFSET,SIGMA,XSTN,IPT)
      IMPLICIT REAL*8 (A-F,O,P,R-Y), COMPLEX*16 (Q,Z)
      DIMENSION QINP(2)
      IF(IPT.GT.0) WRITE(6,990) DK,CR,CW,RHO,OFFSET,SIGMA,XSTN,IPT
55C FCRMAT('0',10X,'QICABP ENTERED WITH:',/,
1      ' ',10X,'DK',11X,'DR',11X,'CW',11X,'RHO',10X,'OFFSET',7X,
2      ' ',SIGMA,8X,'XSTN',5X,'IPT',/, ' ',10X,7(E12.5,' '),13)
      XSTN = XSTN - OFFSET
      IF(XSTN.LE.RHO-1.0CO) GOTO 20
      ICT = IPT - 1
      QCK = DCMLPX(0.0D0,DK)
      CCCNST = CDEXP(-CCMLPX(0.0D0,SIGMA))
      CALL Q1DWXM(DK,DR,CW,RHO,XSTN,QINP,IOT)
      Q1DABP = (QCK*QINP(2) + QINP(1)) * QCONST
      IF(IPT.LE.0) RETURN
      GOTO 60
60  QIDABP = DCMLPX(0.0D0,0.0CO)
      IF(IPT.LE.0) RETURN
      WRITE(6,995) Q1DAEP
995  FCRMAT('0',10X,'Q1DABP = ',E14.7,' ',E14.7)
      RETURN
      END
      SUBROUTINE Q1DWXM(CK,DR,DW,RHO,XSTN,QINP,IPT)
      N IS THE MAXIMUM DEGREE OF THE TERM U IN THE INTEGRALS
      IMPLICIT REAL*8 (A-F,C,P,R-Y), COMPLEX*16 (Z,C)
      DIMENSION QINP(2)
      IF(IPT.GT.0) WRITE(6,990) DK,CR,RFC,XSTN,IPT
55C FCRMAT('0',10X,'Q1DWXM ENTERED WITH ARGUMENTS:',/,
1      ' ',10X,'DK',16X,'DR',16X,'RHO',16X,'XSTN',14X,'IPT',/,
2      ' ',10X,4(E13.6,5X),12,5X,13)

```

Q1DR6245
 Q1DR6250
 Q1DR6255
 Q1DR6260
 Q1DR6265
 Q1DR6270
 Q1DR6275
 Q1DR6280
 Q1DR6285
 Q1DR6290
 Q1DR6295
 Q1DR6300
 Q1DR6305
 Q1DR6310
 Q1DR6315
 Q1DR6320
 Q1DR6325
 Q1DR6330
 Q1DR6335
 Q1DR6340
 Q1DR6345
 Q1DR6350
 Q1DR6355
 Q1DR6360
 Q1DR6365
 Q1DR6370
 Q1DR6375
 Q1DR6380
 Q1DR6385
 Q1DR6390
 Q1DR6395
 Q1DR6400
 Q1DR6405
 Q1DR6410
 Q1DR6415
 Q1DR6420
 Q1DR6425
 Q1DR6430
 Q1DR6435
 Q1DR6440
 Q1DR6445
 Q1DR6450
 Q1DR6455
 Q1DR6460
 Q1DR6465
 Q1DR6470
 Q1DR6475
 Q1DR6480

QLDR6725
QLDR6730
QLDR6735
QLDR6740
QLDR6745
QLDR6750
QLDR6755
QLDR6760
QLDR6765
QLDR6770
QLDR6775
QLDR6780
QLDR6785
QLDR6790
QLDR6795
QLDR6800
QLDR6805
QLDR6810
QLDR6815
QLDR6820
QLDR6825
QLDR6830
QLDR6835
QLDR6840
QLDR6845
QLDR6850
QLDR6855
QLDR6860
QLDR6865
QLDR6870
QLDR6875
QLDR6880
QLDR6885
QLDR6890
QLDR6895
QLDR6900
QLDR6905
QLDR6910
QLDR6915
QLDR6920
QLDR6925
QLDR6930
QLDR6935
QLDR6940
QLDR6945
QLDR6950
QLDR6955
QLDR6960

```

AFEAT = 0.0
AREST = 0.0
FV(1) = F(ALPHA)
FV(3) = F(ALPHA + 0.5*DA)
FV(5) = F(ALPHA + CA)
KOUNT = 3
WT = DA/6.0
QEST = WT*(FV(1) + 4.0*FV(3) + FV(5))
DX = 0.5*DA
FV(2) = F(ALPHA + 0.5*CX)
FV(4) = F(ALPHA + 1.5*DX)
KOUNT = KOUNT + 2
WT = DX/6.0
CESTL = WT*(FV(1) + 4.0*FV(2) + FV(3))
CESTR = WT*(FV(3) + 4.0*FV(4) + FV(5))
CSUM = QESTL + QESTR
ARESTL = WT*(CDABS(FV(1)) + CDABS(4.0*FV(2)) + CCABS(FV(3)))
ARESTR = WT*(CDABS(FV(3)) + CDABS(4.0*FV(4)) + CDABS(FV(5)))
AFEAT = AREAT + ((ARESTL + ARESTR) - AREST)
QDIFF = QEST - CSUM
IF(CDABS(QDIFF).LE.EPS*DABS(AREA))GO TO 2
IF(CDABS(DX).LE.EFCURU*DABS(ALPHA))GO TO 3
IF(LVL.GE.6)GO TC 5
IF(KOUNT.GE.4000)GC TC 6
LVL = LVL + 1
LCRR(LVL) = 0
FIT(LVL) = FV(3)
FET(LVL) = FV(4)
FET(LVL) = FV(5)
CA = DX
CAT(LVL) = CX
ARESTL = ARESTL
ARESTT(LVL) = ARESTR
QESTT(LVL) = QESTR
EPS = EPS/1.4
EFST(LVL) = EPS
FV(5) = FV(3)
FV(3) = FV(2)
GC TO 1
GERRCR = QERROR + QDIFF/15.0
IF(CORR(LVL).EQ.0)GC TO 4
QSUM = CPSUM(LVL) + QSUM
LVL = LVL - 1
IF(LVL.GT.1)GO TO 3
CANS = ((B*J1)*CDEXP(QEXP*B) - CSUM)*Q2EXP/DM
IF(IPT.GT.0)GO TC 11
IF(IFLAG.EQ.1) RETURN

```

1

2 3


```

11 WRITE(6,990) DK,DF,DW,RHO,XSTN,A,B,J,IPT
995 WRITE(6,995) QANS,IFLAG,IER,CERRCR
1. FORMAT(1,15X,'C10ZXM RESULTS: QANS = ',E14.7,',',E14.7,',',
1. ,15X,'IFLAG,2X,IER,5X,'QERRCR',/,
2. ,15X,13,2X,13,5X,E14.7,',',E14.7)
RETURN
4 QFSUM(LVL) = QSUM
LCRR(LVL) = 1
ALPHA = ALPHA + CA
CA = DAT(LVL)
FV(1) = FIT(LVL)
FV(3) = F2T(LVL)
FV(5) = F3T(LVL)
AREST = ARESTT(LVL)
CEST = CESTT(LVL)
GC TO 1 2
IFLAG 2 3
GC TO 2 3
IFLAG 2 3
GC TO 2 3
END
COMPLEX FUNCTION Q1DAIP*16(DK,DR,DW,RHO,CFS1,SGMA,XSTN,Q1CF,N,IPT)
IMPLICIT REAL*8 (A-H,C,P,R-Y); COMPLEX * 16 (Q,Z)
DIMENSION Q1CF(13)
IF (IPT.GT.0) WRITE(6,990) DK,DR,14,RHO,OFST,SGMA,XSTN,IPT
IFCRMAT(0,10X,'Q1DAIP ENTERED WITH: ',/,
1. ,10X,'DK,11X,DR,11X,11X,RHO',10X,'CFS1',9X,'SGMA',9X,
2. ,XSTN,9X,'IPT',/,
3. ,XSTN,XSTN-OFST
4. ,XSTN,XSTN-OFST-1.0D0) GC TC 20
IF(XSTN.LE.CR+RHO-1.0D0) GC TC 20
ICT = IPT - 1
CCCNST = CCEXP(0.0D0,SGMA)
A = XASTN - RHO - 1.0D-8
CALL Q1DZSM(DK,DR,DW,RHO,XSTN,A,B,N,QANS,Q1CF,IOT)
Q1CAIP = QANS
IF(IPT.LE.0) RETURN
GC TO 60
2C Q1DAIP = DCMLX(0.0D0,0.0D0)
IF(IPT.LE.0) RETURN
60 WFITET(6,995) Q1DAIP
995 FCRMAT(0,10X,'Q1CAIP = ',E14.7,',',E14.7)
RETURN
END
COMPLEX FUNCTION Q1DRIP*16(DK,DR,DW,RHO,OFFSET,XSTN,Q1CF,N,IPT)
IMPLICIT REAL*8 (A-H,C,P,R-Y); COMPLEX * 16 (C,Z)
DIMENSION Q1CF(13)

```



```

9950 IF (IPT.GT.0) WRITE(6,990) CK,DR,CW,RHO,XSTN,IPT
FCRMAT(0,10X,'Q1DRIF ENTERED WITH: ',/,
1 '1CX,DK,11X,'DR',11X,'DW,11X,'RHO',10X,
2 'XSTN,9X,IPT,/,',10X,5(E12.5,1,1),13)
IF(XSTN.LE.CW+RHO+OFFSET-1.000) GO TO 20
ICT = IPT - 1
A = DR-1.000
B = XSTN - RHO - 1.0D-8
CALL Q1DZSM(CK,DR,CW,RHO,XSTN,A,B,N,QANS,Q1CF,IOT)
Q1CRIP = -QANS
IF(IPT.LE.0) RETURN
GOTO 60
20 Q1CRIP = DCMPLX(0.0DC,0.0D0)
IF (IPT.LE.0) RETURN
60 WRITE(6,995) Q1DRIP
995 WFCRNAT(0,10X,'Q1CRIP = ',E14.7,1,1,E14.7)
RETURN
ENCL
SUBROUTINE Q1DZSM (DK,DR,DW,RHO,XSTN,A,B,J,CANS,Q1CF,IPT)
IMPLICIT REAL*8(A - E,G,H,M,C,P,F - Y), COMPLEX*16(F,Q,Z)
DIMENSION Q1CF(13),LCRR(60),F1T(60),F2T(60),F3T(60),DAT(60),
D,AREST(60),QEST(60),EPST(60),CFSUM(60)
1 F(X) = CLGNCR(X,CK,DR,J,Q1CF)*CDEXF(CEXP*(X))
1 F(X) = NMBSJ1(OMEGA * DSQRT((XSTN-X)*(XSTN-X) - YY)), IER)
2 GAMMA = 1.400
YY = RHO * RHO
CM2 = (GAMMA + 1.000) * CW
OMEGA = DSQRT (DK*DK*(1.0D0-DM2) / (CM2*DM2))
CLAMDA = DK/DM2
CM = DSQRT(CM2)
CZEXP = CDEXP (DCMPLX (0.0D0, -[LAMDA*XSTN])
CZEXP = CCMPLX(0.0D0,CLAMDA)
U = 9.0D-13
ACC = 1.0D-6
IF (IPT.GT.0) WRITE(6,990)DK,DR,CW,RHO,XSTN,A,B,J,IPT
FCRMAT(0,15X,'Q1DZSM ENTERED WITH ARGUMENTS: ',/,
1 '1CX,DK,13X,'DR',13X,'DW',10X,'RHO',12X,'XSTN',11X,
2 'A,14X,'B,14X,'J,12X,'IPT,/,',
3 '1CX,15X,7(E14.7,1,1),12,2X,I3)
EFCURU = 4.C*U
IFLAG = 1
EPS = ACC
CERROR = DCMPLX(0.0D0,0.0D0)
LVL = 1
LCRR(LVL) = 0.0

```


QLCDR7445
 QLCDR7450
 QLCDR7455
 QLCDR7460
 QLCDR7465
 QLCDR7470
 QLCDR7475
 QLCDR7480
 QLCDR7485
 QLCDR7490
 QLCDR7495
 QLCDR7500
 QLCDR7505
 QLCDR7510
 QLCDR7515
 QLCDR7520
 QLCDR7525
 QLCDR7530
 QLCDR7535
 QLCDR7540
 QLCDR7545
 QLCDR7550
 QLCDR7555
 QLCDR7560
 QLCDR7565
 QLCDR7570
 QLCDR7575
 QLCDR7580
 QLCDR7585
 QLCDR7590
 QLCDR7595
 QLCDR7600
 QLCDR7605
 QLCDR7610
 QLCDR7615
 QLCDR7620
 QLCDR7625
 QLCDR7630
 QLCDR7635
 QLCDR7640
 QLCDR7645
 QLCDR7650
 QLCDR7655
 QLCDR7660
 QLCDR7665
 QLCDR7670
 QLCDR7675
 QLCDR7680

```

ALPHA = A
DA = B - A
AREAL = 0.0
AREST = 0.0
FV(1) = F(ALPHA)
FV(3) = F(ALPHA + 0.5*DA)
FV(5) = F(ALPHA + DA)
KOUNT = 3
WT = DA/6.0
CEST = WT*(FV(1) + 4.0*FV(3) + FV(5))
DX = 0.5*DA
FV(2) = F(ALPHA + 0.5*DX)
FV(4) = F(ALPHA + 1.5*DX)
KOUNT = KOUNT + 2
WT = DX/6.0
CESTL = WT*(FV(1) + 4.0*FV(2) + FV(3))
CESTR = WT*(FV(3) + 4.0*FV(4) + FV(5))
QSUM = CESTL + CESTR
ARESTL = WT*(CDABS(FV(1)) + CDABS(4.0*FV(2))) + CDABS(FV(3)))
ARESTR = WT*(CDABS(FV(3)) + CDABS(4.0*FV(4))) + CDABS(FV(5)))
AREAL = ARESTL + ARESTR
ALPHA = AREST - QSUM
QDIFF = CEST - QDIFF
LE.EPS*DAABS(AREAL) GO TO 2
IF(CDABS(DX).LE.EFCURU*DAABS(ALPHA)) GO TO 5
IF(LVL.GE.60) GO TO 5
IF(KOUNT.GE.4000) GO TO 6
LVL = LVL + 1
LCRR(LVL) = 0
F1T(LVL) = FV(3)
F2T(LVL) = FV(4)
F3T(LVL) = FV(5)
DA = DX
ARESTL = ARESTR
ARESTT(LVL) = ARESTR
CESTL = CEST
CESTT(LVL) = CESTR
EPS = EPS/1.4
CESTL = CESTL + EPS
FV(5) = FV(3)
FV(3) = FV(2)
FV(2) = FV(1)
LCRR(LVL) = LCRR(LVL) + 1
GERRCR = GERRCR + QDIFF/15.0
IF(LCRR(LVL).EQ.0) GO TO 4
QSUM = QSUM + QSUM
LVL = LVL + 1
IF(LVL.GT.1) GO TO 2
QANS = (QLGNDR(B,CR,J,QICF)*CDEXP*B) - QSUM)*Q2EXP/CM
  
```

1

2 3


```

11 IF (IPT.GT.0) GO TO 11
11 IF (IFLAG.EQ.1) RETURN
11 WRITE(6,995) DK,DR,DW,RHG,XSTN,A,E,J,IPT
995 WFORMAT(1,15X,'QANS,IFLAG,IER,GERRCRS = ',E14.7,/,
1,15X,'C1075M RESULTS: CANS = ',E14.7,/,
2,15X,'IFLAG,2X,IER,5X,'GERRCR,/,
2,15X,'13,2X,13,5X,E14.7,/,
2,15X,'13,2X,13,5X,E14.7)
RETURN
4 QSUM(LVL) = QSUM
QCRRL(LVL) = 1
ALPHA = ALPHA + DA
CFA = DAT(LVL)
FV(1) = FIT(LVL)
FV(2) = F2T(LVL)
FV(3) = F3T(LVL)
FV(5) = AREST(LVL)
AREST = QPST(LVL)
GEFEST = EPST(LVL)
5 GEFEST = EPST(LVL)
6 IFLAG TO 1 2
IFLAG TO 2 3
ICEND TO 2 2
ICEND

```


7. Gorelov Spanning Function.

The subprogram for function used by Gorelov is shown below.

```
REAL FUNCTION PLGNDR*8(X,DR,N)
IMPLICIT REAL*8(A-T,C-Z)
IF(N.EQ.0) GOTO 100
ETA = DARCCS(-X)
ETASTR = DARCCS(1.000 - DR)
FN = CFLCAT(N)
PLGNDR = DCCS(FN*ETA)-DCOS(FN*ETASTR)
RETURN
100 PLGNDR = 1.000
RETURN
END
```


8. Linear Expansion Program

The program based on the linear expansion for small k is shown below.


```

C      IMPLICIT REAL*8(A-H,O,P,R-Y), CCMFLEX*16(Q,Z)
C      DIMENSION X(13), Q2PT(13), Q2CP(13)
C      BLOCK ONE READ AND EDIT DATA
C      CALL ERRSET(208,256,-1,1)
C      WRITE(6,910)
C      CALL READ(OK,DR,DW,RHO,OFFSET,SIGMA,N,NF,IPT)
C      IF(IPT.GT.0)WRITE(6,950)OK,DR,DW,SIGMA,OFFSET,RHO,N,NF
C      FCRMAT(1,1,5X,'GORELOV CASCADE PROGRAM LINEARIZED FOR SMALL CK')
C      FCRMAT(1,0,EX,INPUT VALUES TRANSMITTED TO MAIN PROGRAM:,,/,
C      1,1,10X,OK,13X,CR,13X,DW,10X,SIGMA,10X,OFFSET,/,
C      1,1,10X,5(E12.5,3X),/,
C      2,0,10X,RHO,12X,N,5X,NFCN,
C      3,/,1,10X,1(E12.5,3X),12,4X,I2)
C      CALL ABSA(N,OFFSET,X,RHO,DW)

C      BLOCK TWO CALLS CALCULATION ROUTINES FOR ZONE ONE AND ZONE TWO
C      CALL Q$CNE(CK,DR,CW,RHO,OFFSET,SIGMA,N,NF,IFT,X)
C      GO TO 1
C      END
C      SLEROLINE READ (CK,DR,DW,RHO,CFFSET,SIGMA,N,NF,IPT)
C      IMPLICIT REAL*8(A-H,C,P,R-Y)
C      READ(5,905) OK,DR,DW,RHO,OFFSET,SIGMA,N,NF,IPT
C      IF (OK.LT.0.000) STOP
C      IF (IPT.GT.0) WRITE(6,910)
C      GAMMA = 1.4D0
C      DR = RHP * DSQRT((GAMMA + 1.0D0) * DW)
C      RHO = DR
C      IF(N.LE.2) GO TO 1
C      N = 2
C      WRITE(6,935)
C      1 IF(OFFSET.LE.1.9D0) GO TO 2
C      WRITE(6,940)
C      CFFSET = OFFSET-1.0D0
C      GO TO 1
C      CCNTINUE
C      2 IF (IPT.GT.0) WRITE(6,925) DK,DR,DW,SIGMA,CFFSET,RHO,N,NF,IFT
C      RETURN
C      905 FORMAT(6F10.4,3I2)
C      910 FORMAT(1,1,10X,GORELOV SLIGHTLY SUPERSONIC CASCADE PROGRAM',/,
C      1,15X,LINEARIZED FOR SMALL FREQUENCY, DK:'),
C      925 FCRMAT(1,0,10X,CK,13X,DR,13X,DW,10X,SIGMA,10X,OFFSET,
C      1,/,1,10X,5(E12.5,3X),/,
C      2,0,10X,RHO,12X,N,5X,NFCN,2X,IPT'
C      3,/,1,10X,1(E12.5,3X),12,4X,I2)
C      935 FCRMAT(1,0,5X,ORIGINAL N FCC LARGE (.GT.13) - RESET TO N = 13')
C      940 FCRMAT(1,0,5X,OFFSET TOO LARGE (.GT.1.9), RESET AS OFFSET =
C      10FFSET - 1.0D0)

```



```

IN = I + N
XSTN = X(I) - DR
XC = CEXP( DCMPLEX( C.000, DLAMDA*XSTN) )
CALL QICRIP( DK, DR, CW, RHO, OFFSET, SIGMA, XSTN, CINTAP, N, IOT )
CALL QICRIP( DK, DR, DW, RHC, OFFSET, XSTN, QINTRP, N, ICT )
DO 20 J = 1, N
  JN = J +
  J1 = J - 1
  Q1INT( I, 1 ) = DCMPLEX( 1.000, CLAMDA*XSTN )
  TEMP = (R-XSTN-1.000
  Q1INT( I, 2 ) = DCMPLEX( TEMP, CLAMCA*XSTN*TEMP )
  Q1INT( I, JN ) = QINTRP( J )
  Q1INT( I, JN ) = QINTAP( J )
  XASTN = XSTN - OFFSET
  Q1INT( I, 3 ) = DCMPLEX( 1.000, DLAMDA*XASTN )
  TEMP = DR-XASTN-1.000
  Q1INT( I, 4 ) = DCMPLEX( TEMP, DLAMDA*XASTN*TEMP )
  CCNTINUE
20  C1CCF( IN ) = -QICRBP( DK, DR, DW, RHO, XSTN, IOT )
   C1CCF( I ) = -QICABP( DK, DR, DW, RHC, CFFSET, SIGMA, XSTN, IOT )
   CCNTINUE
90  C1CCF( I ) = -QICABP( DK, DR, DW, RHC, CFFSET, SIGMA, XSTN, IOT )
   CCNTINUE
   ENCL
   CCMPLEX FUNCTION QICRBP*16( DK, DF, DW, RHO, XSTN, IPT )
   IMPLICIT REAL*8 ( A-H, O, P, R-Y ), COMPLEX*16 ( Q, Z )
   DIMENSION QINP( 2 )
   IF ( IPT.GT. 0 ) WRITE( 6, 990 ) DK, CR, CW, RHO, XSTN, IPT
   FCFORMAT( 0, 10X, 'QICRBP ENTERED WITH: ', /
   , 10X, DK, 11X, CR, 11X, DW, 11X, RHO, 10X, XSTN, 9X, IPT,
   /, XSTN, LE, RHO-1.000 ) GOTC 20
   IF( XSTN.GT. 2.000 ) GOTO 20
   ICT = IPT - 1
   QICK = DCMPLEX ( 0.000, DK )
   GAMMA = 1.400
   CLAMDA = DK/( ( GAMMA+1.000 ) * DW )
   C1CRBP = DCMPLEX( 1.000, ( DK+DLAMDA ) * ( XSTN-DR ) )
   IF ( IPT.LE.0 ) RETURN
   GCTO 30
20  C1CRBP = DCMPLEX( 0.000, 0.000 )
   IF ( IPT.LE.0 ) RETURN
30  WRITE( 6, 995 ) QICRBP
995  FORMAT( 'C', 10X, 'QICRBP = ', E14.7, ', ', E14.7 )
   RETURN
   ENCL
   CCMPLEX FUNCTION C1CABP*16( DK, DR, CW, RHO, OFFSET, SIGMA, XSTN, IPT )
   IMPLICIT REAL*8 ( A-H, O, P, R-Y ), COMPLEX*16 ( Q, Z )

```

```

LIN00570
LIN00580
LIN00590
LIN01000
LIN01010
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LIN01030
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LIN01210
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LIN01300
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LIN01320
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LIN01340
LIN01350
LIN01360
LIN01370
LIN01380
LIN01390
LIN01400
LIN01410
LIN01420
LIN01430
LIN01440

```



```

CJMENSICN QINP(2)
IF(IPT.GT.0) WRITE(6,990) DK,DR,DW,RHO,OFFSET,SIGMA,XSTN,IPT
FORMAT(0,'10X',Q1CABP ENTERED WITH:','/,
1  '10X',CK,11X,DR,11X,DW,11X,RHO,10X,OFFSET',7X,
2  'SIGMA',8X,XSTN,9X,IPT,/,',10X,7(E12.5,','),13)
XASTN = XSTN - OFFSET
IF(XASTN.LE.RHO-1.000) GOTO 20
IF(XASTN.GT.2.000) GOTO 20
IPT = IPT - 1
QCK = DCMLPX(0.000,DK)
CCCNST = CDEXP( DCMLPX(0.000,SIGMA))
GAMMA = 1.400
CLAMDA = CK/((GAMMA+1.000)*DW)
Q1CABP = DCMLPX(1.000,(DK+CLAMDA))*((XASTN-DR))*QCONST
IF(IPT.LE.0) RETURN
GC TO 30
20 Q1CABP = DCMLPX(C.000,0.000)
IF(IPT.LE.0) RETURN
30 WRITE(6,995) Q1CABP
995 FCFORMAT(0,'10X',Q1CABP = ',E14.7,',',E14.7)
RETURN
END
SUBROUTINE Q1CAIP(CK,DR,DW,RHO,CFST,SGMA,XSTN,QCAIP,N,IPT)
IMPLICIT REAL*8 (A-H,C,P,R-Y), COMPLEX * 16 (C,Z)
DIMENSION QINP(13),QCAIP(13)
IF(IPT.GT.0) WRITE(6,990) CK,DR,DW,RHO,CFST,SGMA,XSTN,IPT
FCFORMAT(0,'10X',Q1CAIP ENTERED WITH:','/,
1  '10X',CK,11X,DR,11X,DW,11X,RHO,10X,OFST',9X,SGMA',9X
2  'XSTN',9X,IPT,/,',10X,7(E12.5,','),13)
XASTN = XSTN - CFST
IF(XASTN.LE.DR+RHC-1.000) GO TO 20
IF(XASTN.GT.2.000) GOTO 20
IPT = IPT - 1
CCCNST = CDEXP( DCMLPX(0.000,SGMA))
GAMMA = 1.400
CLAMDA = CK/((GAMMA+1.000)*DW)
QCAIP(1) = DCMLPX(1.000,CLAMDA*(XASTN-RHC))
TEMP = DR+RHO-XASTN-1.000
QCAIP(2) = DCMLPX(TEMP,CLAMDA*(XASTN-RHO)*TEMP)
IF(IPT.LE.0) RETURN
GOTO 30
20 ZERO = DCMLPX(0.000,0.000)
DC25 I = 1,N
QCAIP(I) = ZERO
CCCONTINUE
25 IF(IPT.LE.0) RETURN
30 WRITE(6,995)
995 FCFORMAT(0,'10X',Q1CAIP RESULTS: J QCAIP(J)')

```



```

DC 40 J = 1,N
WRITE (6,556) J,CCAIIP(J)
FCRMAT(,26X,I2,3X,E14.7,,',E14.7)
CCNTINUE
RETURN
END
SUBROUTINE CCRIP (DK,DR,DH,RHO,CFFSET,XSTA,CCRIP,A,IPT)
IMPLICIT REAL*8 (A-H,C,P,R-Y), COMPLEX * 16 (Q,Z)
DIMENSION CCRIP(13)
IF (IPT.GT.0) WRITE(6,990) DK,DR,DH,RHO,XSTA,IPT
FCRMAT(,10X,CCRIP ENTERED WITH:,,',
,10X,DK,11X,DR,11X,DH,11X,RHO,1CX,
,XSTN,9X,IPT,/,10X,5(E12.5,,'),13)
IF(XSTN.LE.DR+RHO+CFFSET-1.0CO) GO TO 20
IF(XSTN.GT.2.0DO) GO TO 20
IPT = IPT - 1
GAMMA = 1.4CO
DLAMDA = DK/(GAMMA+1.0CO)*DW)
CCRIP(1) = CCPLX(1.0DO,DLAMDA*(XSTN-RHO))
TEMP = DR+RHO-XSTN-1.0DO
CCRIP(2) = CCPLX(TEMP,DLAMDA*(XSTN-RHO)*TEMP)
IF(IPT.LE.0) RETURN
GO TO 30
ZERO = CCPLX(0.0DO,0.0DO)
DC 25 I = 1,N
CCRIP(I) = ZERO
IF (IPT.LE.0) RETURN
WRITE (6,995)
FCRMAT(,10X,CCRIP RESULTS: J CCRIP(J),
DC 40 J = 1,N
WRITE (6,996) J,CCRIP(J)
FCRMAT(,26X,I2,3X,E14.7,,',E14.7)
CCNTINUE
RETURN
END
SUBROUTINE GLCOEF( Q1COF,Q1INT,A,IPT,Q1ABCF,Q1RBCF)
IMPLICIT REAL*8(A-H,C,P,R-Y), COMPLEX * 16 (Z,Q)
DIMENSION Q1COF(26), Q1INT(26,26), ZWA(300)
DIMENSION Q1ABCF(13),Q1RBCF(13)
IF (IPT.GE.0) WRITE (6,90)
IB = 26
M2 = 1
N2 = 2*N
IF (IPT.LE.0) GO TO 5
WRITE (6,98) N,N2
FCRMAT(,10X,Q1COEF ENTERED WITH ',I2,' DEG PWR SERIES (',I2,
1 DC 2 I = 1,N2

```



```

WRITE(6,92) I, I, Q1COF(I)
FORMAT(0,10X,Q1CCEF EQUATION SYSTEM, ROW,12,/,
1,10X,Q1COF(,I2,)) = ,E14.7,1,E14.7)
CC 2 J = 1,N
J2=J+N
WRITE(6,91) I,J,CLINT(I,J),I,J2,CLINT(I,J2)
WFCRMA(,15X,2(QINT(,I2,,I2,)) = ,E14.7,,E14.7,10X))
CCNTINUE
IA = 26
IJOB = 0
CALL LECTIC(Q1INT,N2,IA,Q1CCF,M,IE,IJOB,ZWA,IER)
IF(IER.EQ.0) GOTC 30
IF(IER.EQ.129) GC TO 10
WRITE(6,93)
FCRMA(,C,10X,Q1CCEF - ITERATIVE IMPROVEMENT FAILED, MATRIX TOO
1ILL-CONDITIONED. USE RESULTS WITH CAUTION.)
GC TO 30
WRITE(6,95)
FCRMA(,0,,10X,Q1COEF - MATRIX ALGORITHMICALLY SINGULAR. CCEFFIC
1IENTS SET TO ZERO.)
ZERC = DCMPLX(0.0D0,C.0D0)
CC 20 I = 1,N
I2 = I+N
Q1COF(I) = ZERO
Q1CCF(I2) = ZERO
CCNTINUE
CCNTINUE
CC 35 I = 1,N
IN = I + N
Q1ABCF(I) = Q1COF(IN)
Q1RBCF(I) = Q1COF(I)
CCNTINUE
IF(IPT.LE.0) RETURN (6,94)
IF(IPT.GE.C) WRITE (6,94)
CC 40 I = 1,N
IN1 = I-1
WRITE(6,99) I, IN1, Q1RBCF(I), Q1ABCF(I)
CCNTINUE
FCRMA(,0,5X,I2,10X,I2,10X,2(E14.7,5X,E14.7,10X))
FCRMA(,C,10X,SUBROUTINE Q1CCEF - COMPLEX POWER SERIES CCEFFIC
1IENTS)
FCRMA(,0,5X,INDEX,7X,DEG FCY,4X,REFERENCE BLADE TERMS Q1C
3CGF(INDEX),7X,ADJACENT BLADE TERMS Q1CCF(2*INDEX))
RETURN
SUBROUTINE CLPOT(DK,DR,DW,RHO,CFFSET,SIGMA,N,X,Q1ABCF,Q1RBCF,IPT)
IMPLICIT REAL*8(A-H,O,P,R-Y), COMPLEX*16(Z,Q)
DIMENSION Q1ABCF(13), Q1RBCF(13), CRR(13), CAA(13), CPHI(13)

```



```

WRITE (6,95) I,XSTN
WRITE (6,91) COR,QCRO,CDAP,QCRIPO,CCAIPP
WRITE (6,92) CCA,QCRP,CDAO,QDRIPP,CCAIP0
QCCR(I) = CCR
QCAA(I) = QCA
FCRMAT (,0 X-STATICN NUMBER ,I2, ,XSTN = ,F6.4,/,
1, ,BL TOTAL D(PCT)/DX, ,14X, ,REF EL D(POT)/DX, ,8X,
2, ,ADJ BL D(POT)/DX, ,8X, ,REF BL INT D(POT)/DX, ,5X,
3, ,ADJ BL INT D(POT)/DX, )
CCNTINUE
1C WRITE (6,954)
994 FCRMAT (,1, ,10X, ,SUMMARY LISTING,/,
1,C, ,10X, ,XSTN, ,7X, ,SINGLE BLADE TOTAL POTENTIAL, ,6X,
1, ,REF ELADE POTENTIAL, ,15X, ,ADJ ELADE POTENTIAL, )
DC 20 I = 1,N
XSTN = X(I)
WRITE (6,94) XSTN,CPhi(I),QRR(I),CAP(I)
94 FCRMAT (, ,10X,F6.4,3(5X,EL4.7, , ,E14.7) )
20 CCNTINUE
996 WRITE (6,996)
996 FCRMAT (,0, ,10X, ,XSTN, ,7X, ,SINGLE ELADE TOTAL C(POT)/CX, ,6X,
1, ,REF ELADE D(POT)/DX, ,15X, ,ADJ ELADE D(POT)/DX, )
DC 50 I = 1,N
XSTN = X(I)
WRITE (6,94) XSTN,CDPhi(I),QCRR(I),GDAA(I)
5C CCNTINUE
RETURN
END
SUBROUTINE C1DCOF (CK,DR,DW,RHO,CFFSET,SIGMA,N,C1ABCF,C1RBCF,IPT)
IMPLICIT REAL * 8 (A-F,O,P,R-Y), COMPLEX * 16 (Q,Z)
DIMENSION Q1ABCF(13),Q1RBCF(13)
IF (IPT.GT.0) WRITE (6,990) DK,DR,DW,RHO,OFFSET,SIGMA,N,IPT
C1(IPT,1,1,10X,Q1DCOF - CALCULATION OF COMPLEX DIMENSIONLESS AERC
FCRMAIC CCEFFICIENTS,/,
1, ,10X, ,DK, ,13X, ,DR, ,13X, ,Dw, ,13X, ,RHO, ,12X, ,OFFSET, ,5X, ,SIGMA,
3, ,0, ,10X, ,3X, ,IPT, ,/,
4, ,10X, ,N, ,3X, ,12,2X, ,13)
56 (E12.5,3X),12,2X, ,13)
ICIT = IPT - 3
999 CCCL = C1CCL(DK,DR,DW,RHO,OFFSET,SIGMA,N,Q1ABCF,Q1RBCF,IOT)
QCCM = Q1DCM(DK,DR,DW,RHO,OFFSET,SIGMA,N,C1ABCF,Q1RBCF,IOT)
GAMMA = 1.4CO
TAU = (2.0DO*DSQRT((GAMMA + 1.0CO)*DW))/DR
WRITE (6,90) DK,TAU,DW,N,SIGMA,CCCL,QDCM
FCRMAT (,0, ,5X, ,DK = ,F6.3, , ,TAU = ,F7.4, , ,Dh = ,F6.3, , ,N =
1,I2, , ,SIGMA = ,F6.3, , ,CL = ,F9.4, , ,F9.4, , ,CM = ,
2,F9.4, , ,F9.4)
RETURN
END

```



```

FV(5) = F3T(LVL)
AREST = ARESTT(LVL)
QEST = QESTT(LVL)
EFS = EPST(LVL)

```

```

5 GC TO 1 2
  IFLAG = 2 3
6 GC TO 2
  IFLAG = 2
GC TO 2

```

```

END
COMPLEX FUNCTION CIPHI*16 (DK,DR,CK,RHO,OFFSET,SIGMA,

```

```

1 N,Q1ABCF,Q1RBCF,XSTN,IPT)
  IMPLICIT REAL*8(A-H,O,P,R-Y),COMPLEX*16(C,Z)
  DIMENSION Q1ABCF(13),Q1RBCF(13)
  IF (IPT.GT.C) WRITE(6,990) CK,DR,CW,RHO,CFFSET,SIGMA,XSTN,IPT
  FCFORMAT(0,10X,Q1PHI ENTERED WITH:,,/
1 1,10X,DK,11X,DR,11X,DW,11X,RHO,10X,OFFSET,7X,SIGMA,
2 8X,1X,XSTN,9X,IPT,/,10X,7(E12.5,.,,13)
  RZERC = 0.000
  ICT = IPT - 1
  C1PHI = Q1PREP(CK,CR,DW,RZERO,XSTN,IOT)
1 + Q1PRIP(CK,DR,DW,RZERO,OFFSET,XSTN,Q1RBCF,N,IGT)
2 + Q1PAEP(CK,CR,CW,DR,OFFSET,SIGMA,XSTN,IOT)
3 + Q1PAIP(CK,DR,DW,DR,OFFSET,SIGMA,XSTN,C1ABCF,N,IOT)
  IF (IPT.LE.0) RETURN
  WRITE(6,995) Q1PHI
995 FCFORMAT(0,10X,C1PHI = ',E14.7,.',',E14.7)
  RETURN

```

```

COMPLEX FUNCTION Q1PRBP*16(DK,DF,CW,RHO,XSTN,IPT)
IMPLICIT REAL*8 (A-H,O,P,R-Y), COMPLEX*16 (Q,Z)
DIMENSION QINP(2)
IF (IPT.GT.0) WRITE(6,990) CK,CR,CW,RHO,XSTN,IPT
FCFORMAT(0,10X,C1PRBP ENTERED WITH:,,/
1 1,1CX,DK,11X,CR,11X,DW,11X,RHC,10X,XSTN,9X,IPT,
2 /,10X,5(E12.5,.,,13)
  IF (XSTN.LE.RHO-1.000) GOTO 20
  ICT = IPT - 1
  GAMMA = 1.400
  CM2 = (GAMMA + 1.000) * DW
  CLAMDA = CK/DM2
  PR = 1.0 + XSTN-RHO
  PIMAG = (CK+DLAMDA)*((XSTN-RHO)*(XSTN-RHO) - 1.000)
  Q1PRBP = - CCMPLEX(FR,PIMAG-CLAMCA*XSTN*PR)/CSCRT(CM2)
  IF (IPT.LE.0) RETURN
  GOTO 60

```

```

20 C1PREP = CCMPLEX(C-GDO,0.000)
  IF (IPT.LE.0) RETURN

```

```

LIN05770
LIN05780
LIN05790
LIN05800
LIN05810
LIN05820
LIN05830
LIN05840
LIN05850
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LIN05870
LIN05880
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LIN05900
LIN05910
LIN05920
LIN05930
LIN05940
LIN05950
LIN05960
LIN05970
LIN05980
LIN05990
LIN06000
LIN06010
LIN06020
LIN06030
LIN06040
LIN06050
LIN06060
LIN06070
LIN06080
LIN06090
LIN06100
LIN06110
LIN06120
LIN06130
LIN06140
LIN06150
LIN06160
LIN06170
LIN06180
LIN06190
LIN06200
LIN06210
LIN06220
LIN06230
LIN06240

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CCOMPLEX FUNCTION C1DREP*16(CK,DR,EH,RHO,XSTN,IPT)
IMPLICIT REAL*8 (A-H,C,P,R-Y), COMPLEX*16 (Q,Z)
DIMENSION C1NP(2)
IF (IPT.GT.0) WRITE(6,990) CK,DR,EH,RHO,XSTN,IPT
IF (IPT.CE.10X,'Q1DRBP ENTERED WITH: ',//
FCRMAT('C',10X,'11X',CR,'11X',DW,'11X',RHC,'1CX',XSTN,'9X',IPT,
1 1,10X,'11X',5(E12.5,'',13)
2 /,'XSTN,E,RHO-1.000) GOTO 20
IF(XSTN.LE.RHO-1.000) GOTO 20
ICT = IPT - 1
GAMMA = 1.400
GM2 = (GAMMA + 1.000) * DW
LLAMDA = DK/DM2
LM2 = (GAMMA + 1.000) * DW
PIMAG = (DK+DLAMDA)*(XSTN-RHO)
Q1DRBP = -DCMPLX(1.000,PIMAG-CLAMCA*XSTN)/DSQRT(DM2)
IF(IPT.LE.0) RETURN
GOTO 60
20 Q1DRBP = DCMPLX(0.000,0.000)
IF(IPT.LE.0) RETURN
60 WRITE(6,995) Q1DRBP
FCRMAT('C',10X,'Q1DRBP = ',E14.7,'',E14.7)
995 RETURN
END
CCOMPLEX FUNCTION C1DAEP*16(DK,DR,EH,RHO,OFFSET,SIGMA,XSTN,IPT)
IMPLICIT REAL*8 (A-H,C,P,R-Y), COMPLEX*16 (Q,Z)
DIMENSION C1NP(2)
IF (IPT.GT.0) WRITE(6,990) DK,DR,EH,RHO,OFFSET,SIGMA,XSTN,IPT
IF (IPT.CE.10X,'Q1DAEP ENTERED WITH: ',//
FCRMAT('C',10X,'11X',CR,'11X',DW,'11X',RHC,'10X','OFFSET',7X,
1 1,SIGMA,'8X',XSTN,'9X',IPT,'',10X,'7(E12.5,'',13)
2 XSTN=XSTN-CFFSET
IF(XSTN.LE.RHO-1.000) GOTO 20
ICT = IPT - 1
GCCNST = CDEXP( DCMPLX(0.000,SIGMA))
GAMMA = 1.400
GM2 = (GAMMA + 1.000) * DW
LLAMDA = DK/DM2
PIMAG = (DK+DLAMDA)*(XSTN-RHO)
Q1CABP = DCMPLX(1.000,PIMAG-DLAMCA*XSTN)*GCCNST/DSQRT(DM2)
IF(IPT.LE.0) RETURN
GOTO 60
20 Q1CABP = DCMPLX(0.000,0.000)
IF (IPT.LE.0) RETURN
60 WRITE(6,995) Q1CABP
FCRMAT('C',10X,'Q1DAEP = ',E14.7,'',E14.7)
995 RETURN
END
CCOMPLEX FUNCTION C1PAIP*16(CK,DR,EH,RHO,CFST,SGMA,XSTN,Q1CF,N,IPT)

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INPLICIT REAL*8 (A-H,C,P,R-Y), COMPLEX * 16 (G,Z)
DIMENSION QICF(13)
IF (IPT.GT.0) WRITE(6,990) CK,DR,CW,RHO,OFST,SGMA,XSTN,IPT
FCRMAT(0,10X,'QIPAI ENTERED WITH:',/,
1 1CX,DK,11X,DR,11X,DW,11X,RHO,10X,OFST,9X,SGMA,5X
2 XSTN,9X,IPT,/,',10X,7(E12.5,','),13)
XSTN = XSTN - OFST
IF(XSTN.LE.DR+RHC-1.000) GO TO 20
GCONST = CCEXP(DCMPLX(0.000,SGMA))
GAMMA = 1.400
DM2 = (GAMMA+1.000)*DW
DLAMDA = DK/DM2
XX=XSTN-RHO
Q1 = QICF(1)
Q2 = QICF(2)
PR1 = (XX+1.000-DR - XX/2.000)*XX - (CR-1.000)*(DR-1.000)/2.000
PR2 = ((CR-1.000-DR)*(2.000-DR)*DR-1.000)*DLAMDA/2.000
FIMAG1 = XX*XX*((DR-1.000)/2.000-XX)-((DR-1.000)*3)/6.000
PIMAG2 = PIMAG1*DLAMDA
PIMAG1 = PIMAG2*DLAMDA
PIMAG2 = PIMAG1*DLAMDA
QP1 = DCMPLX(PR1,FIMAG1-DLAMDA*PR1*XSTN)
QP2 = DCMPLX(PR2,PIMAG2-DLAMDA*PR2*XSTN)
QIPAI = (Q1*QP1 + Q2*QP2)/DSQRT(CM2)
IF(IPT.LE.0) RETURN
GCTC 30
QIPAI = DCMPLX(0.000,0.000)
IF(IPT.LE.0) RETURN
WRITE(6,995) QIPAI
FCRMAT(0,10X,'QIPAI = ',E14.7,',',E14.7)
RETURN
END
COMPLEX FUNCTION QIPRIP*16(CK,DR,CW,RHO,OFFSET,XSTN,CICF,N,IPT)
INPLICIT REAL*8 (A-H,O,P,R-Y), COMPLEX * 16 (Q,Z)
DIMENSION QICF(13)
IF (IPT.GT.0) WRITE(6,990) DK,DR,DH,RHO,XSTN,IPT
FCRMAT(0,10X,'QIPRIP ENTERED WITH:',/,
1 1CX,DK,11X,DR,11X,DH,11X,RHO,10X,
2 XSTN,9X,IPT,/,',10X,5(E12.5,','),13)
IF(XSTN.LE.CR+RHO+OFFSET-1.000) GO TO 20
GAMMA = 1.400
CM2 = (GAMMA+1.000)*DW
CLAMDA = DK/DM2
XX=XSTN-RHO
Q2 = QICF(2)
Q1 = QICF(1)
PF1 = XX+1.000-DR
PR2 = (CR-1.000)*XX - (DR-1.000)*(DR-1.000)/2.000

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FIMAG1 = (XX*XX+(2.0D0-DR)*DR-1.0D0)*DLAMDA/2.0D0
FIMAG2 = XX*XX*((DR-1.0D0)/2.0D0-XX)-((DR-1.0D0)**3)/6.0D0
FIMAG1=PIMAG1*DLAMCA
FIMAG2 = PIMAG2 * CLAMDA
CP1 = DCMPLEX(PR1,PIMAG1-DLAMDA*PR1*XSTN)
CP2 = DCMPLEX(PR2,PIMAG2-DLAMDA*PR2*XSTN)
Q1PRIP = -(Q1*QPI+C2*CP2)/DSQRT(LM2)
IF(IPT.LE.0) RETURN
GCTO 30
C1FRIPT = DCMPLEX(0.0D0,0.0D0)
IF (IPT.LE.0) RETURN
WRITE(6,995) Q1PRIP
FCRMTAT('C',10X,'Q1PRIP = ',E14.7,' ',E14.7)
RETURN
END
CCMPLEX FUNCTION Q1DAIP*16(DK,DR,CR,CW,RHO,OFST,SGMA,XSTN,CICF,A,IPT)
IMPLICIT REAL*8 (A-H,O,P,R-Y), COMPLEX * 16 (C,Z)
DIMENSION Q1CF(13)
IF (IPT.GT.0) WRITE(6,990) DK,DR,CR,CW,RHO,CFS1,SGMA,XSTN,IPT
FCRMTAT('O',10X,'Q1CAIP ENTERED WITH:',/,
,1CX,'CK',11X,'DR',11X,'DW',11X,'RHO',1CX,'OFST',9X,'SGMA',5X
,XSTN,9X,'IPT',/,/,10X,7(E12.5,/,),1E)
XSTN=XSTN-OFST
XSTN=XSTN-LE*DR+RHC-1.0D0 GC TO 20
IF(XSTN.LE.DR+RHC-1.0D0) GC TO 20
GCCNST = CODEXP( DCMPLEX(0.0D0,SGMA))
GAMMA = 1.4D0
CLAMDA = (GAMMA+1.0C0)*CW
CLAMDA = DK/DM2
XX = XSTN - RHO
Q1 = Q1CF(2)
Q1 = Q1CF(1)
PR1 = 1.0D0
PR2 = DR-1.0D0-XX
PIMAG1 = CLAMDA*XX
PIMAG2 = DLAMDA*(DR-1.0D0-XX)*XX*XX
CP1 = DCMPLEX(PR1,PIMAG1-DLAMDA*PR1*XSTN)
CP2 = DCMPLEX(PR2,PIMAG2-DLAMDA*PR2*XSTN)
Q1CAIP = (Q1*QPI + Q2*CP2)/DSQRT(LM2)
IF(IPT.LE.0) RETURN
GCTO 30
C1DAIP = CCMPLEX(0.0D0,0.0D0)
IF(IPT.LE.0) RETURN
WRITE(6,995) Q1DAIP
FCRMTAT('O',10X,'Q1DAIP = ',E14.7,' ',E14.7)
RETURN
END
CCMPLEX FUNCTION C1DRIP*16(DK,CR,CW,RHO,OFFSET,XSTN,CICF,N,IPT)
IMPLICIT REAL*8 (A-H,O,P,R-Y), COMPLEX * 16 (Q,Z)

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CIMENSICN QICF(13)
IF (IPT.GT.0) WRITE(6,990) DK,DR,CW,RHO,XSTN,IPT
FORMAT(0,'10X',QIDRIP,ENTERED WITH:','//FO',10X,
',10X',CK',11X',DR',11X',CW',11X',RFO',10X,
',XSTN',9X',IPT','//',10X,5(E12.5','),13)
2 IF(XSTN.LE.DR+RHO+OFFSET-1.0CO) GO TO 20
GAMMA = 1.4CO
CM2 = (GAMMA+1.000)*DW
CLAMDA = DK/CM2
Q1 = QICF(1)
Q2 = QICF(2)
XX = XSTN - RHO
PR1 = 1.0CO
PR2 = DR-1.000-XX)*XX*XX
PIMAG1 = CLAMDA*XX
PIMAG2 = CLAMDA*PR1*XSTN)
QP1 = CCMPLEX(FR1,PIMAG1-CLAMDA*PR1*XSTN)
QP2 = CCMPLEX(PR2,PIMAG2-CLAMDA*PR2*XSTN)
QICRIP = -(Q1*CP1+Q2*QP2)/DSQRT(DM2)
IF (IPT.LE.0) RETURN
GCTC 30
CLDRIP = CCMPLEX(C.000,0.000)
20 IF (IPT.LE.0) RETURN
30 WRITE(6,995) QIDRIP
995 FORMAT(0',10X',QIDRIP = ',E14.7','',E14.7)
RETURN
END

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LI N08420C
LI N08430C

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LIST OF REFERENCES

1. Elder, P. R., A Theoretical Analysis of Unsteady Transonic Cascade Flow, Thesis, Naval Postgraduate School, Monterey, California, 1972.
2. Schlein, P. B., A Study of Unsteady Transonic Interference Effects, Thesis, Naval Postgraduate School, Monterey, California, 1975.
3. Landahl, M., Unsteady Transonic Flow, Pergamon Press, 1961.
4. Gorelov, D. N., "Oscillations of a Plate Cascade in a Transonic Gas Flow," Mekhanika Zhidkosti i Gaza, v. 1, no. 1, pp. 69-74, 1966.
5. Garrick, I. E. and Rubinow, S. J., Flutter and Oscillating Air-Force Calculations for an Airfoil in Two-Dimensional Supersonic Flow, NACA Report 846, 1946.
6. Stevens, W. P.G. J. Meyers, and L. L. Constantine, "Structured Design," IBM Systems Journal, No. 2, pp. 115-139, 1974.
7. Shampine, L. F., and R. C. Allen, Numerical Computing: An Introduction, W. B. Saunders Company, 1973.
8. Ashley, H. and Landahl, M., Aerodynamics of Wings and Bodies, Addison-Wesley Publishing Company, Inc., 1965.
9. Miles, J. W., "The Compressible Flow Past an Oscillating Airfoil in a Wind Tunnel," Journal of Aeronautical Sciences, v. 23, no. 7, pp. 671-678, 1956.

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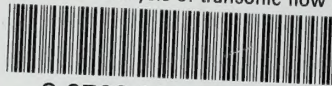
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